

Many-body localization in Quasiperiodic systems

Yi-Ting Tu (涂懿庭)

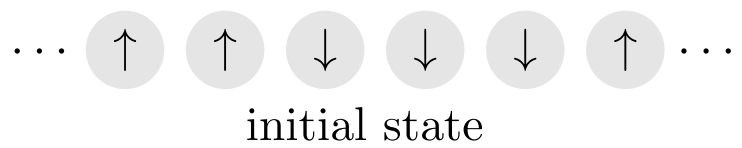
*Condensed Matter Theory Center and Joint Quantum Institute,
Department of Physics, University of Maryland*

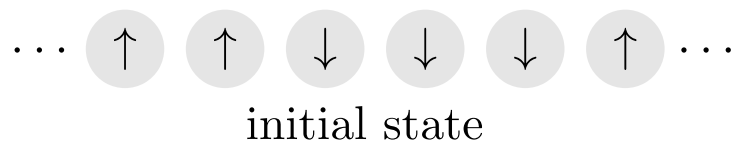
Doctoral defense: June 9, 2026

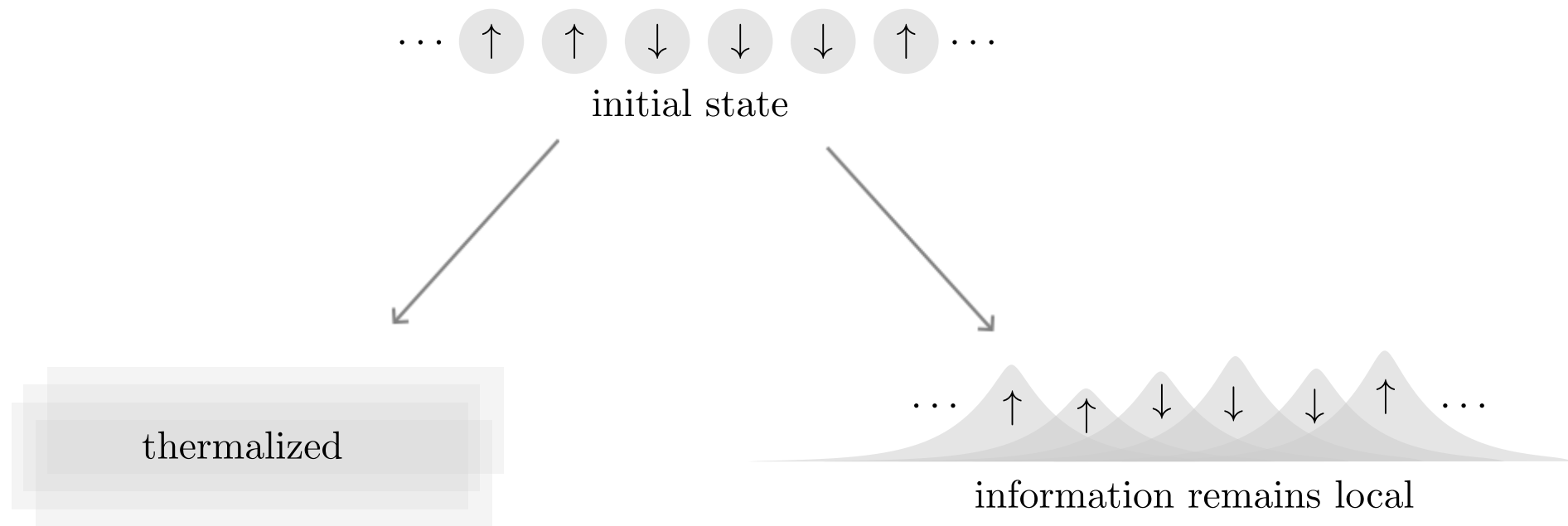


Outline

- Introduction
 - Many-body localization (MBL)
 - Random v.s. Quasiperiodic systems
- Non-ergodic extended (NEE) regime
 - Numerical results
 - Theoretical explanation
- Many-body critical (MBC) phases
 - Single-particle hierarchical mirror structure (HMS)
 - Effects of interaction on HMS
 - Asymptotically solvable model

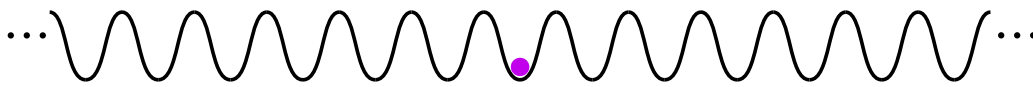




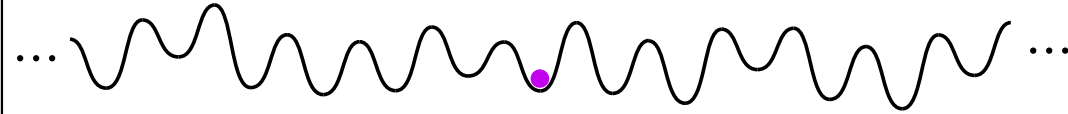


Single particle extended/localized systems

A local particle in a periodic potential:

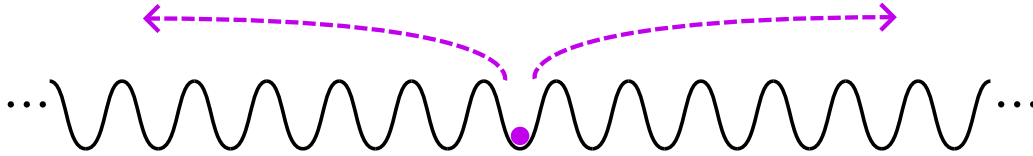


A local particle in a random potential:



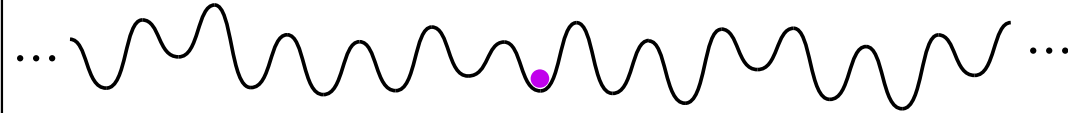
Single particle extended/localized systems

A local particle in a periodic potential:



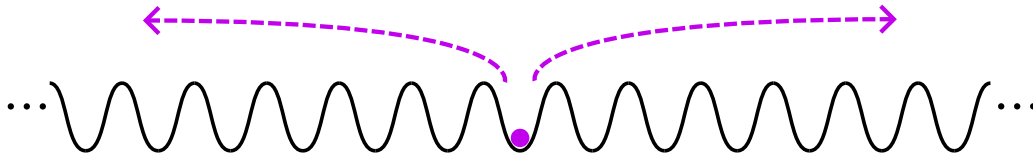
Travels further and further away
and never comes back

A local particle in a random potential:



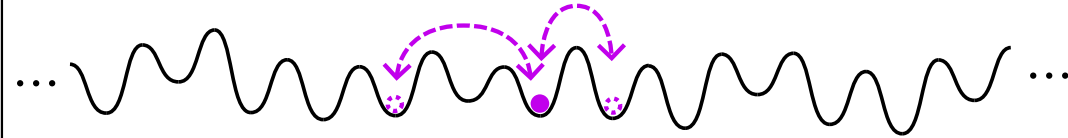
Single particle extended/localized systems

A local particle in a periodic potential:



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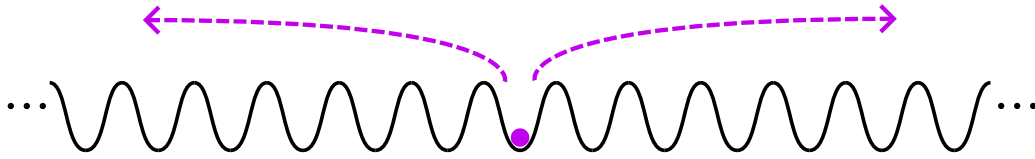
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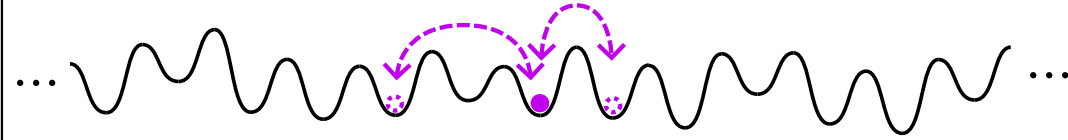
Comes back regularly
and does not go arbitrarily far away
[Anderson1958]

Many-body counterparts of extendedness/localization

Extended particle:

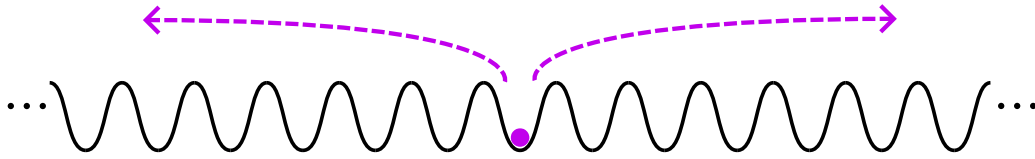


Localized particle:

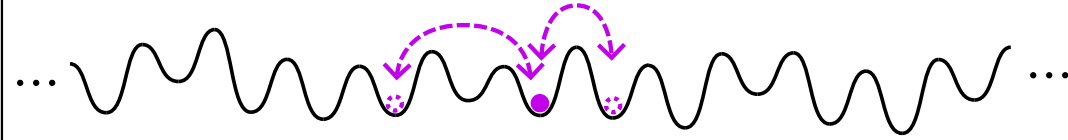


Many-body counterparts of extendedness/localization

Extended particle:

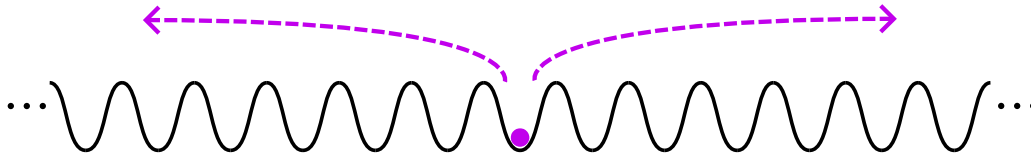


Localized particle:

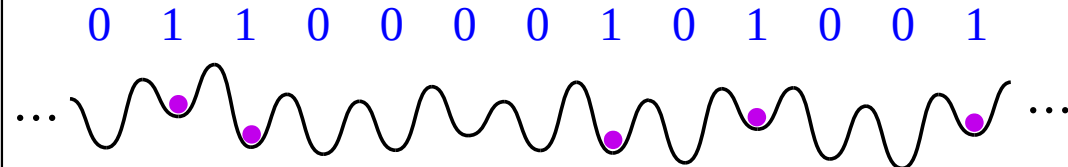
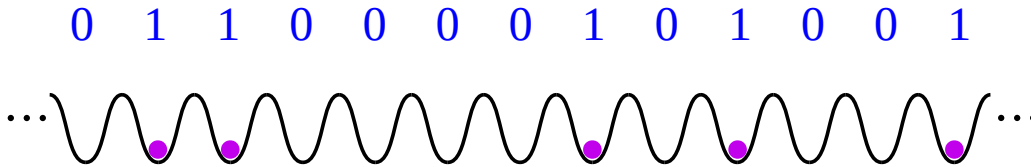
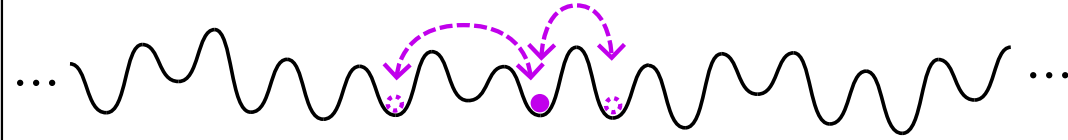


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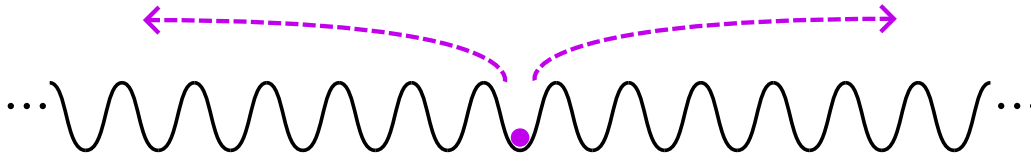


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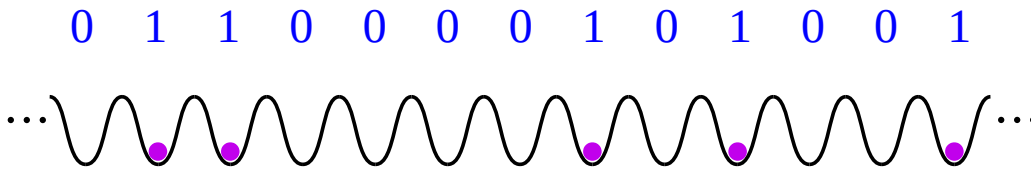
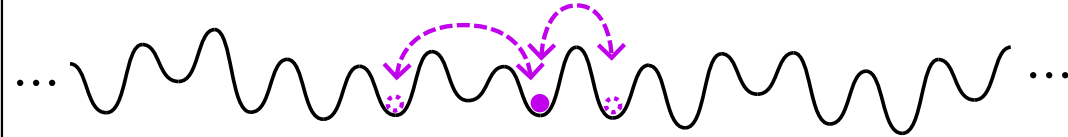


Many-body counterparts of extendedness/localization

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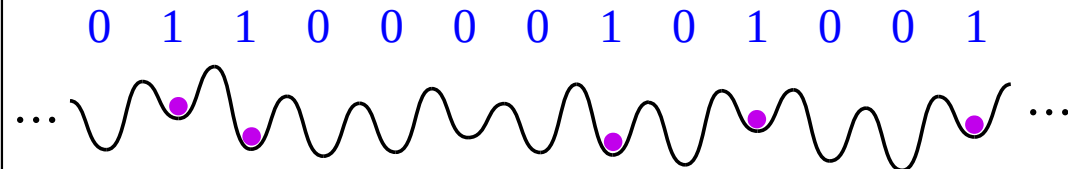


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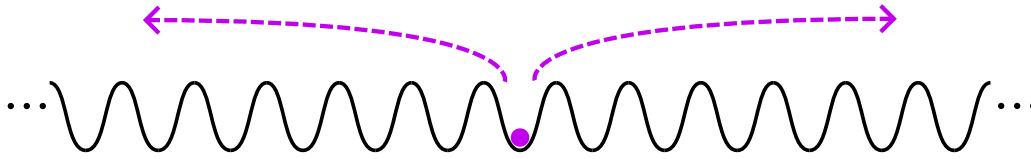
Infinite time

Information spreads out to infinite
Everything looks the same (thermal) locally

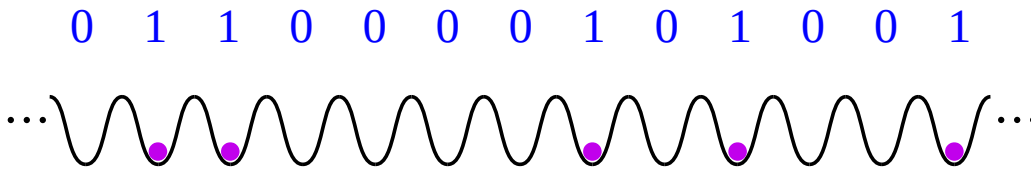
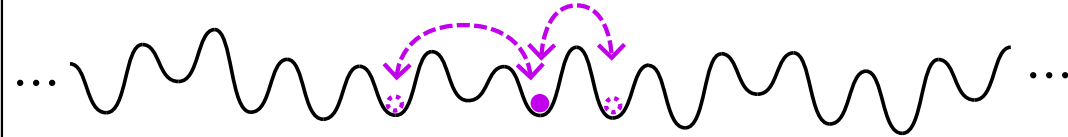


Many-body counterparts of extendedness/localization

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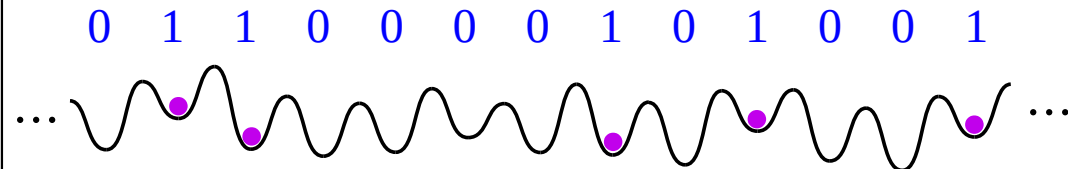


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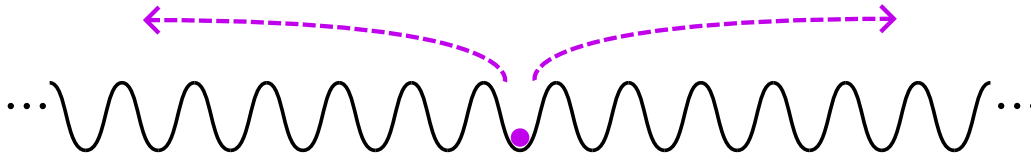


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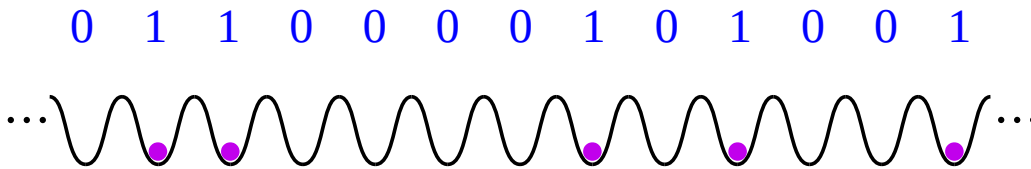
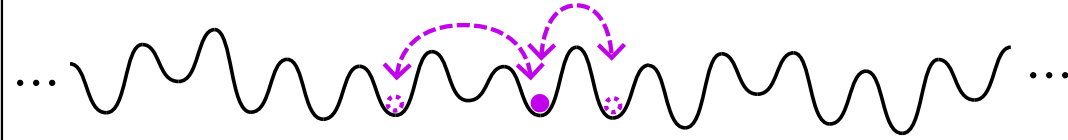
Information only spreads out a little
Memory remains in local observable expectations

Many-body counterparts of extendedness/localization

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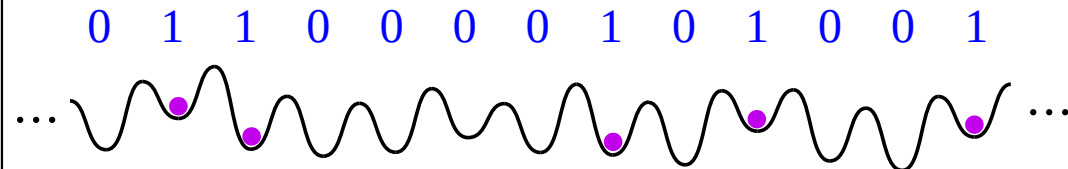
Localized particle:



Infinite time

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Everything looks the same (thermal) locally

ETH phase



Infinite time

Information only spreads out a little
Memory remains in local observable expectations

MBL phase

Many-body counterparts of extendedness/localization

ETH phase

MBL phase

Many-body counterparts of extendedness/localization

ETH phase

Satisfies the eigenstate thermalization hypothesis (ETH):

MBL phase

Many-body counterparts of extendedness/localization

ETH phase

Satisfies the eigenstate thermalization hypothesis (ETH):

$$\langle E_\alpha | O | E_\alpha \rangle \approx \langle O \rangle_{\text{microcanonical}}$$

O : local observable

$|E_\alpha\rangle$: energy eigenstate with E_α in the microcanonical shell

MBL phase

Many-body counterparts of extendedness/localization

ETH phase

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MBL phase

Has a complete set of local integrals of motion (LIOMs):

Many-body counterparts of extendedness/localization

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MBL phase

Has a complete set of local integrals of motion (LIOMs):

$$\{L_j\}, [H, L_j] = 0$$

L_j : localized around site j

Energy eigenstate uniquely labeled by eigenvalues of $\{L_j\}$

Many-body counterparts of extendedness/localization

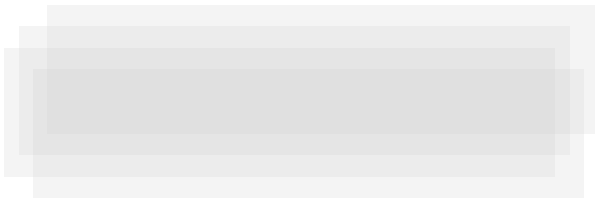
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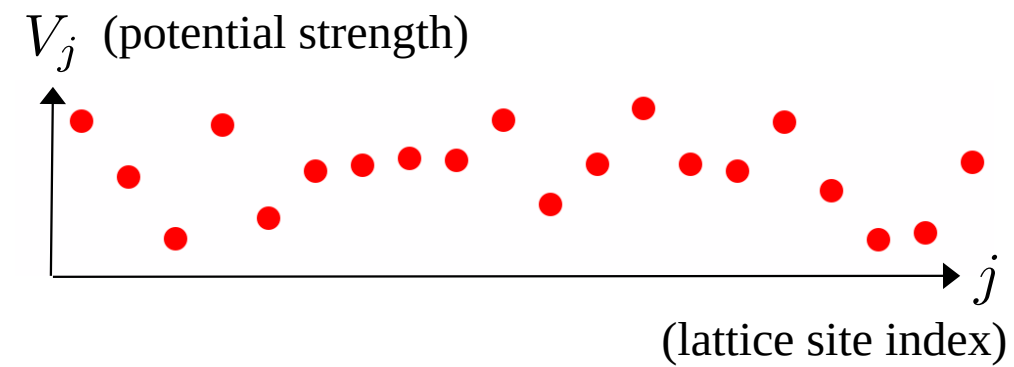
Energy eigenstate uniquely labeled by eigenvalues of $\{L_j\}$



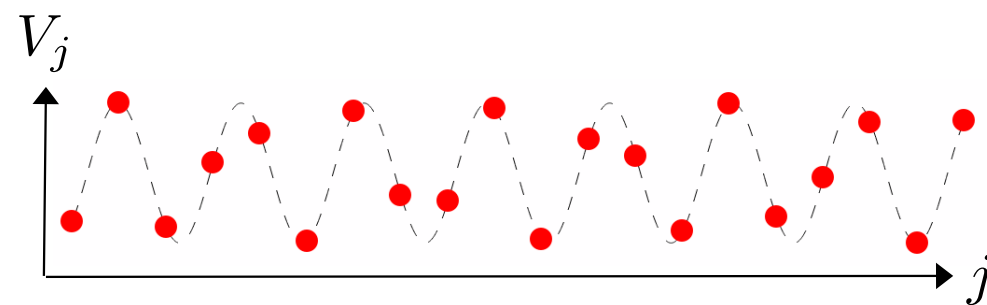
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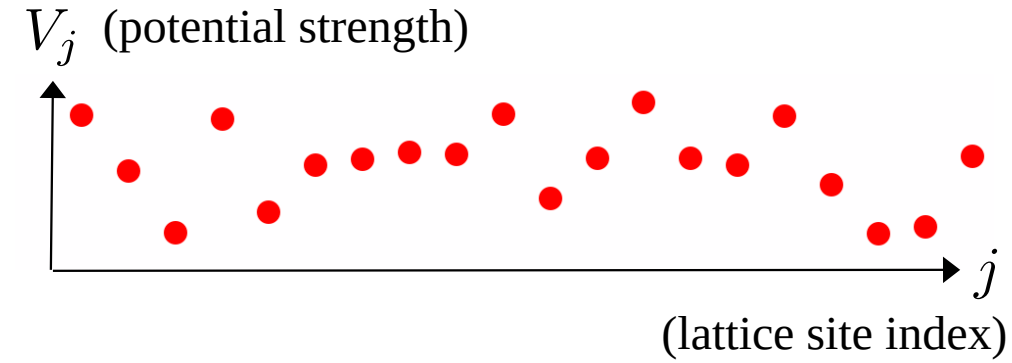
Random systems



Quasiperiodic systems

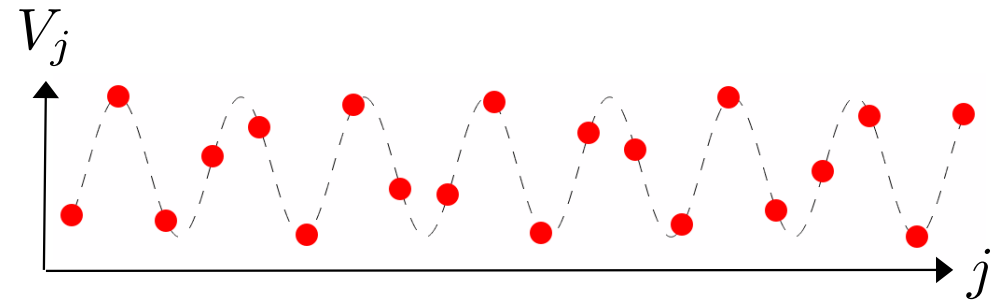


Random systems

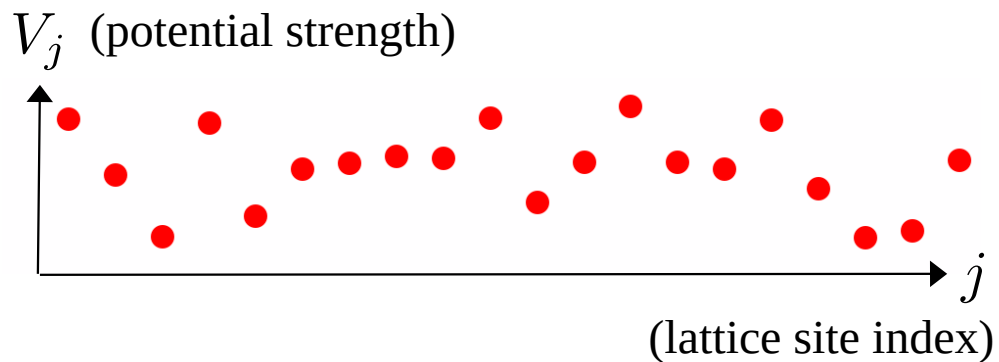


V_j : random numbers
indep. for each j

Quasiperiodic systems

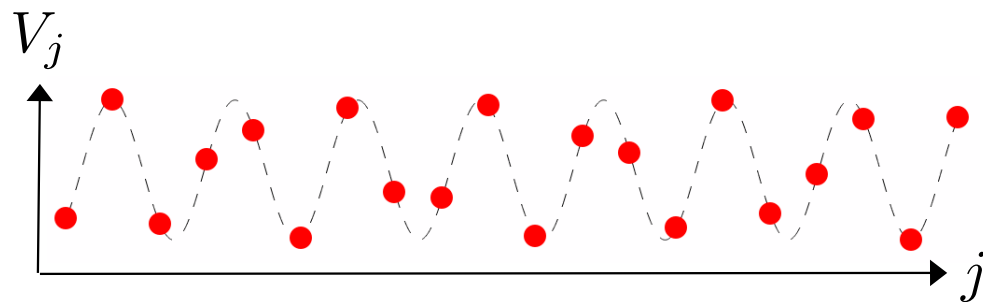


Random systems



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Quasiperiodic systems

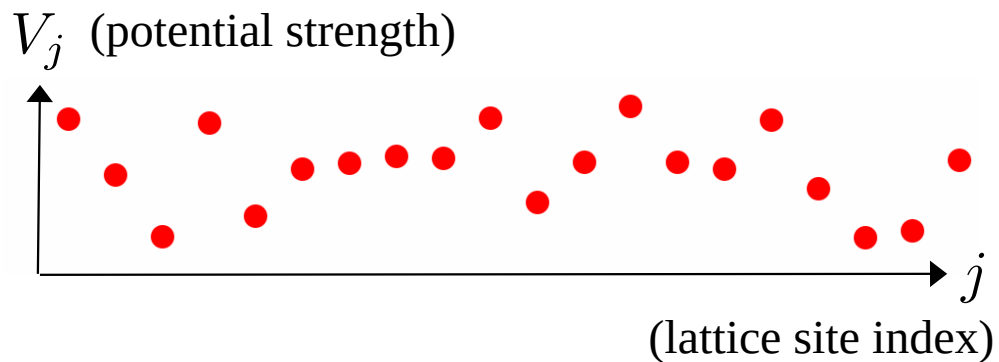


$$V_j = F(\omega j + \phi)$$

$$F(y + 2\pi) = F(y)$$

$$j \in \mathbb{Z}, \quad \frac{\omega}{2\pi} \notin \mathbb{Q}$$

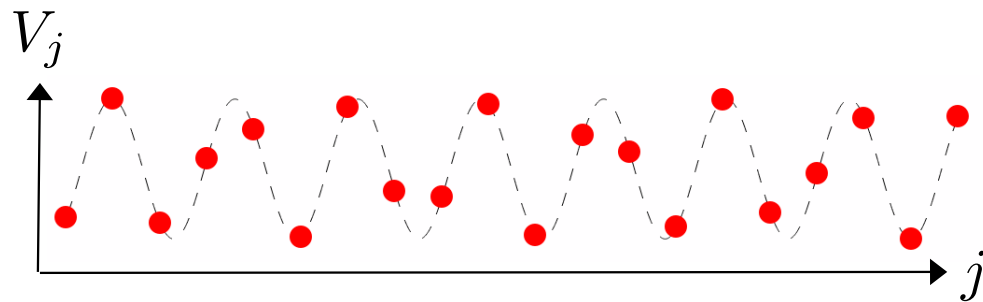
Random systems



V_j : random numbers
indep. for each j

Scale $V_j \in [-W, W]$, W : disorder strength

Quasiperiodic systems



$$V_j = F(\omega j + \phi)$$

$$F(y + 2\pi) = F(y)$$

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What's special about 1D quasiperiodic systems

Random (Anderson)

Quasiperiodic

Single
particle

Interacting
Many-body

What's special about 1D quasiperiodic systems

Random (Anderson)

Quasiperiodic

W (disorder strength)

Single
particle

Localized

E
(energy)

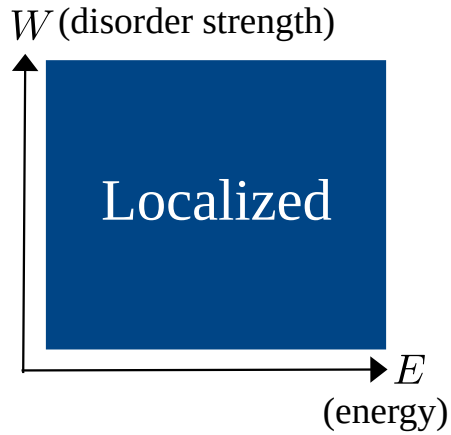
Interacting
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What's special about 1D quasiperiodic systems

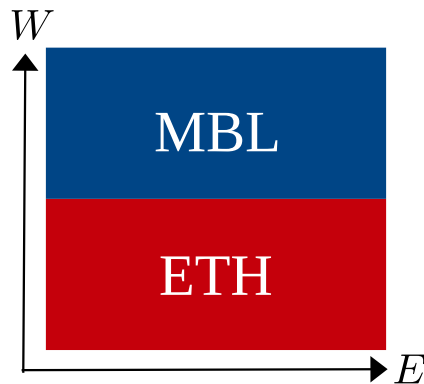
Random (Anderson)

Quasiperiodic

Single
particle

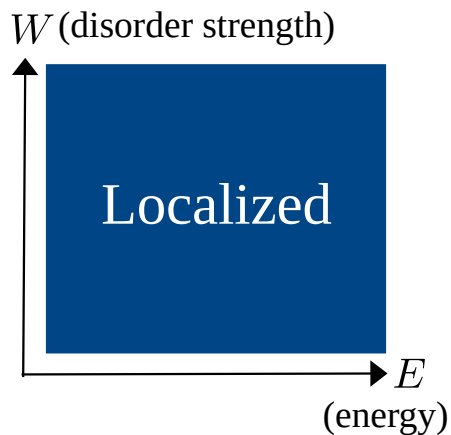


Interacting
Many-body
(at least for
small size)

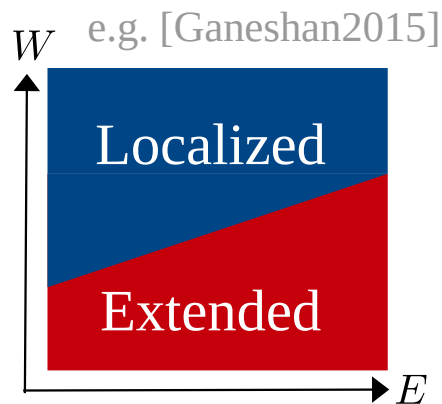


What's special about 1D quasiperiodic systems

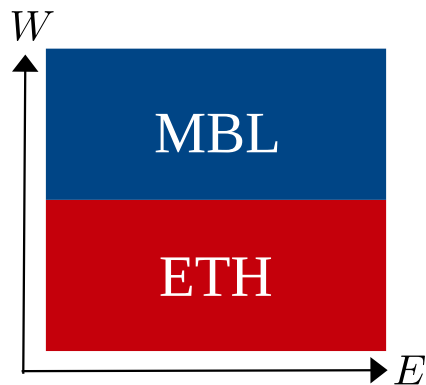
Random (Anderson)



Quasiperiodic

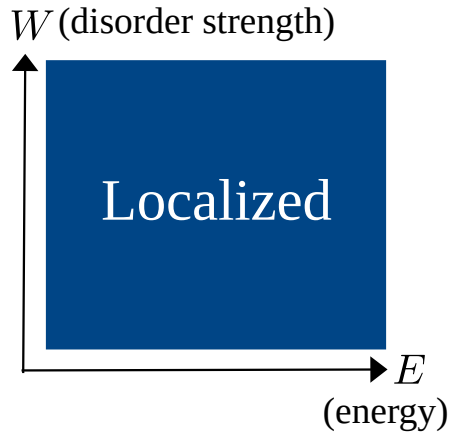


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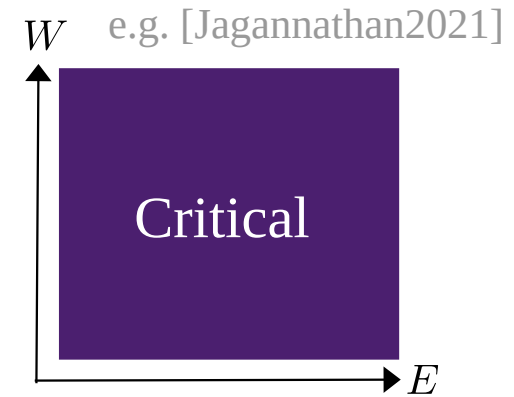
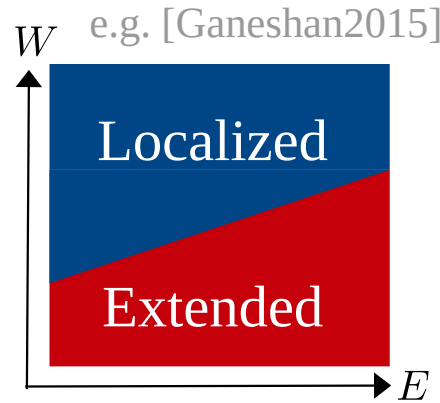


What's special about 1D quasiperiodic systems

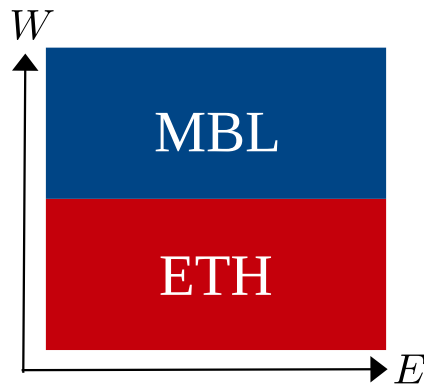
Random (Anderson)



Quasiperiodic

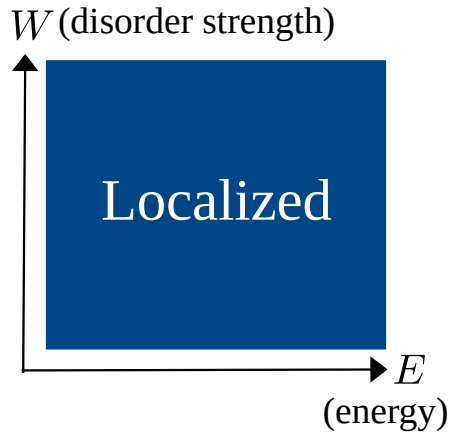


Interacting
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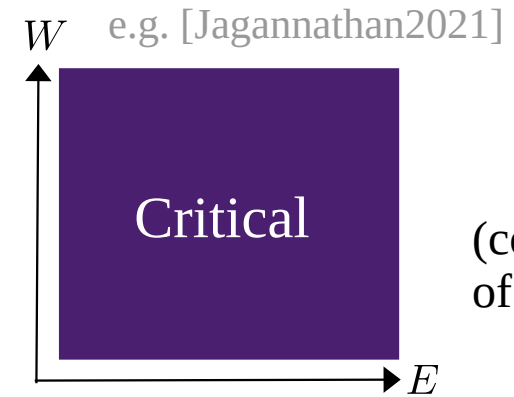
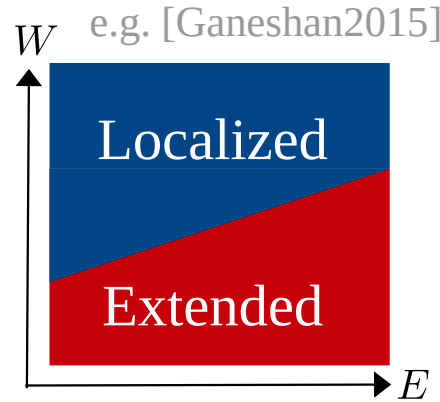


What's special about 1D quasiperiodic systems

Random (Anderson)



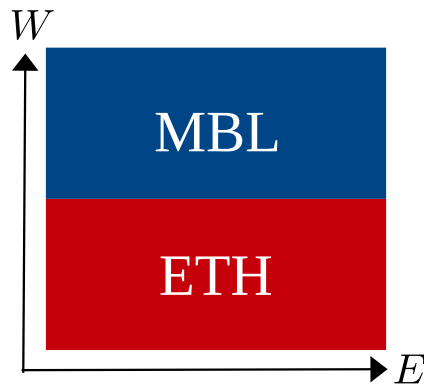
Quasiperiodic



...

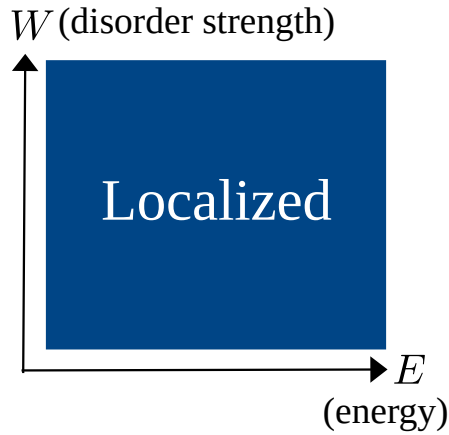
(combination of them)

Interacting
Many-body
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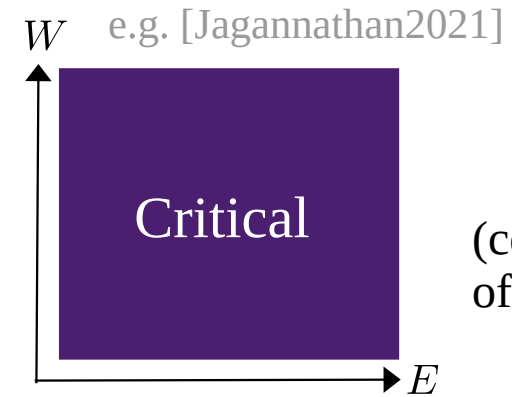
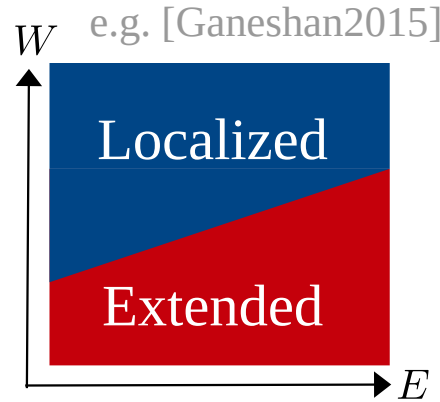


What's special about 1D quasiperiodic systems

Random (Anderson)



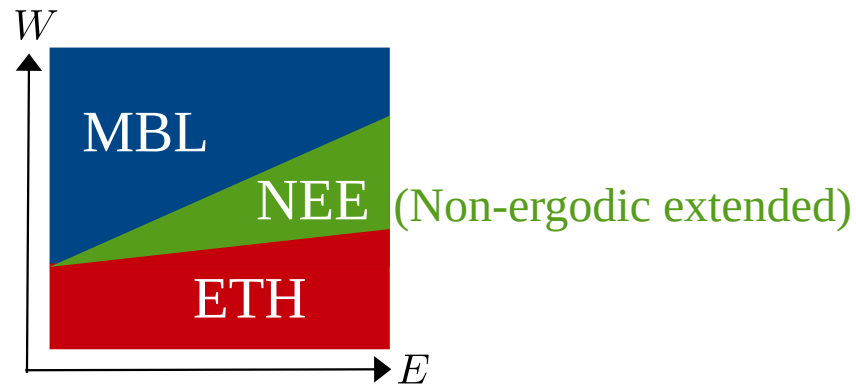
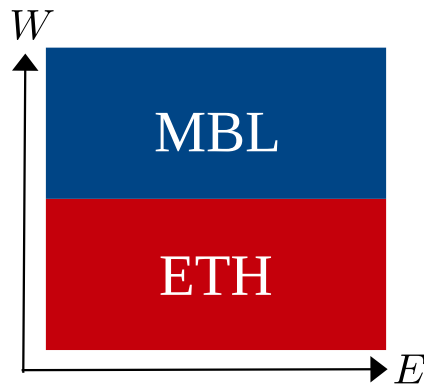
Quasiperiodic



...

(combination of them)

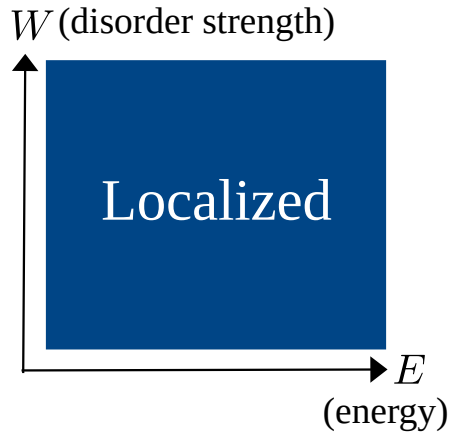
Interacting
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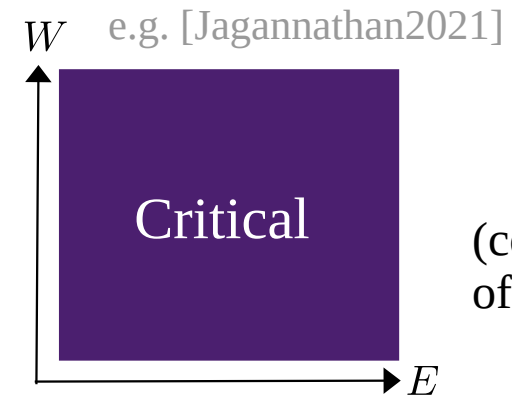
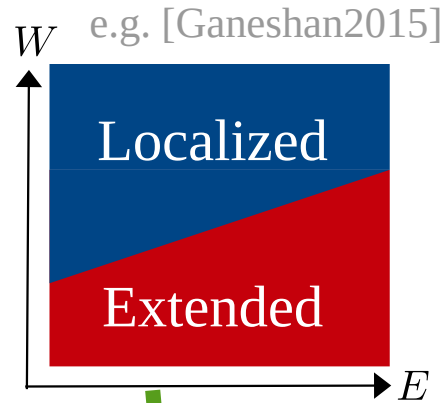
[Li2015, Hsu2018]

What's special about 1D quasiperiodic systems

Random (Anderson)



Quasiperiodic

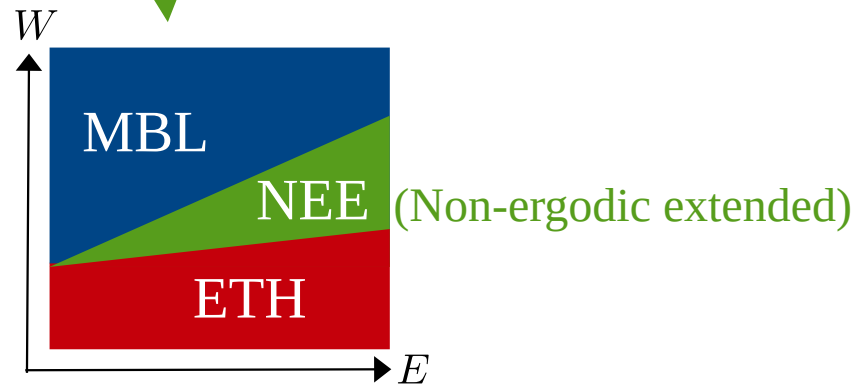
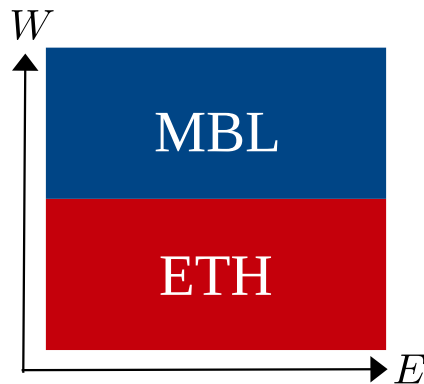


...

(combination of them)

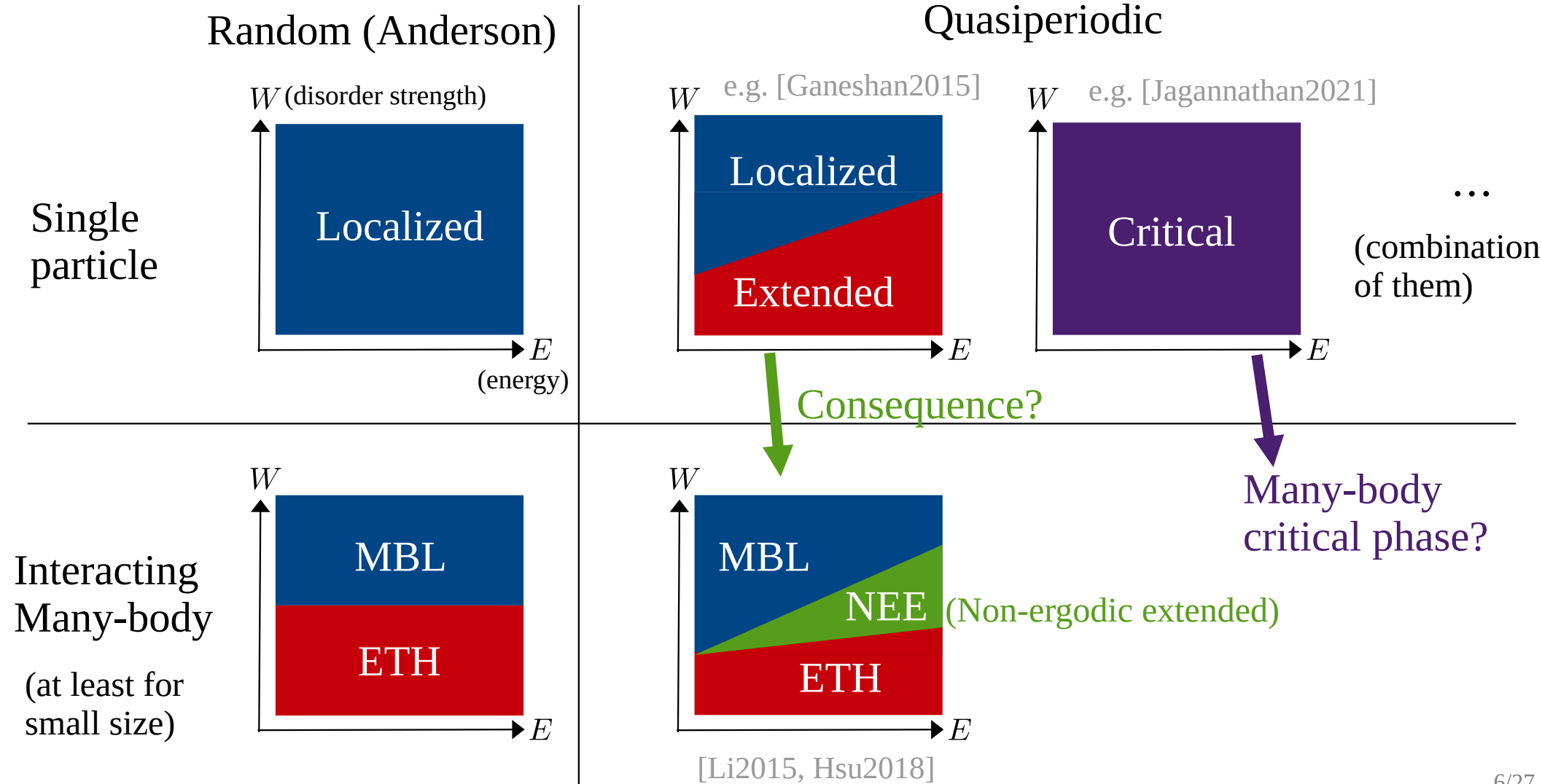
Consequence?

Interacting
Many-body
(at least for
small size)



[Li2015, Hsu2018]

What's special about 1D quasiperiodic systems





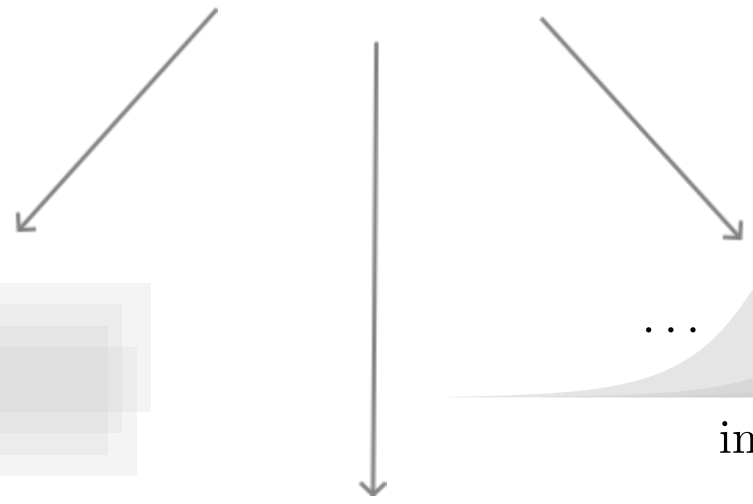
initial state



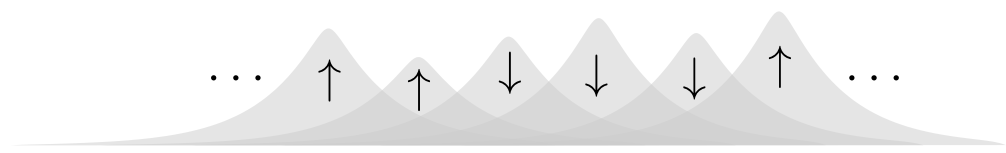
information remains local



initial state



thermalized



information remains local



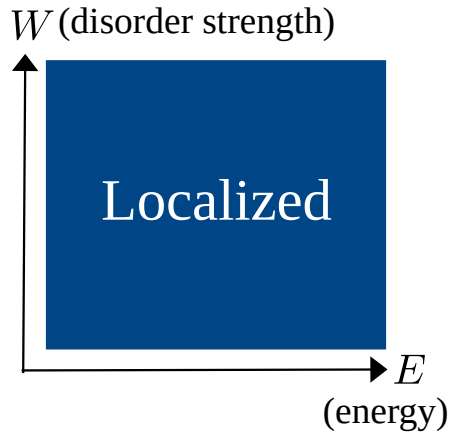
something in between?
critical phase?

Outline

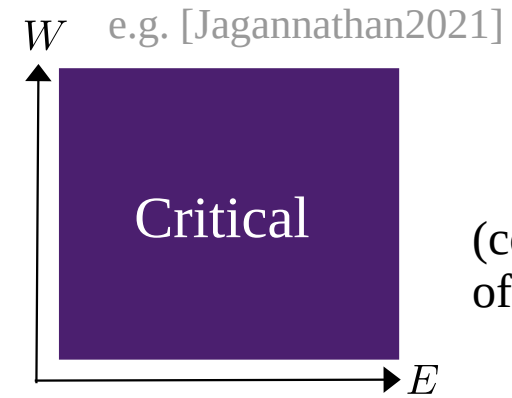
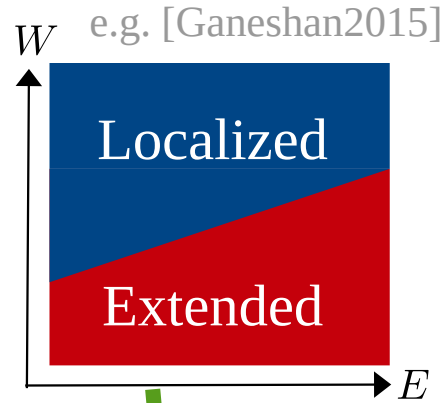
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What's special about 1D quasiperiodic systems

Random (Anderson)



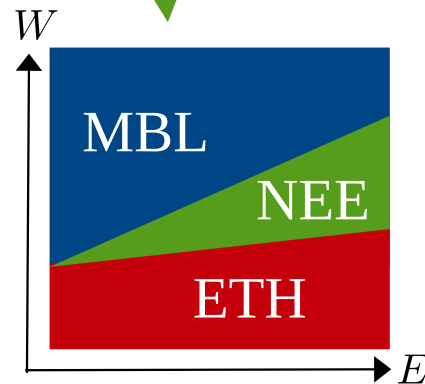
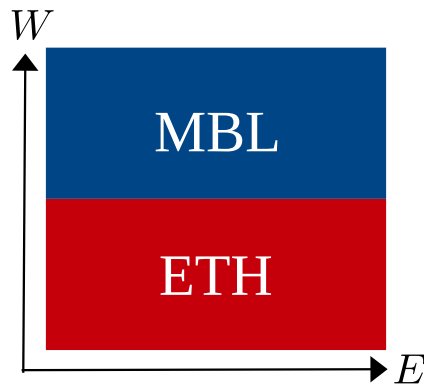
Quasiperiodic



...

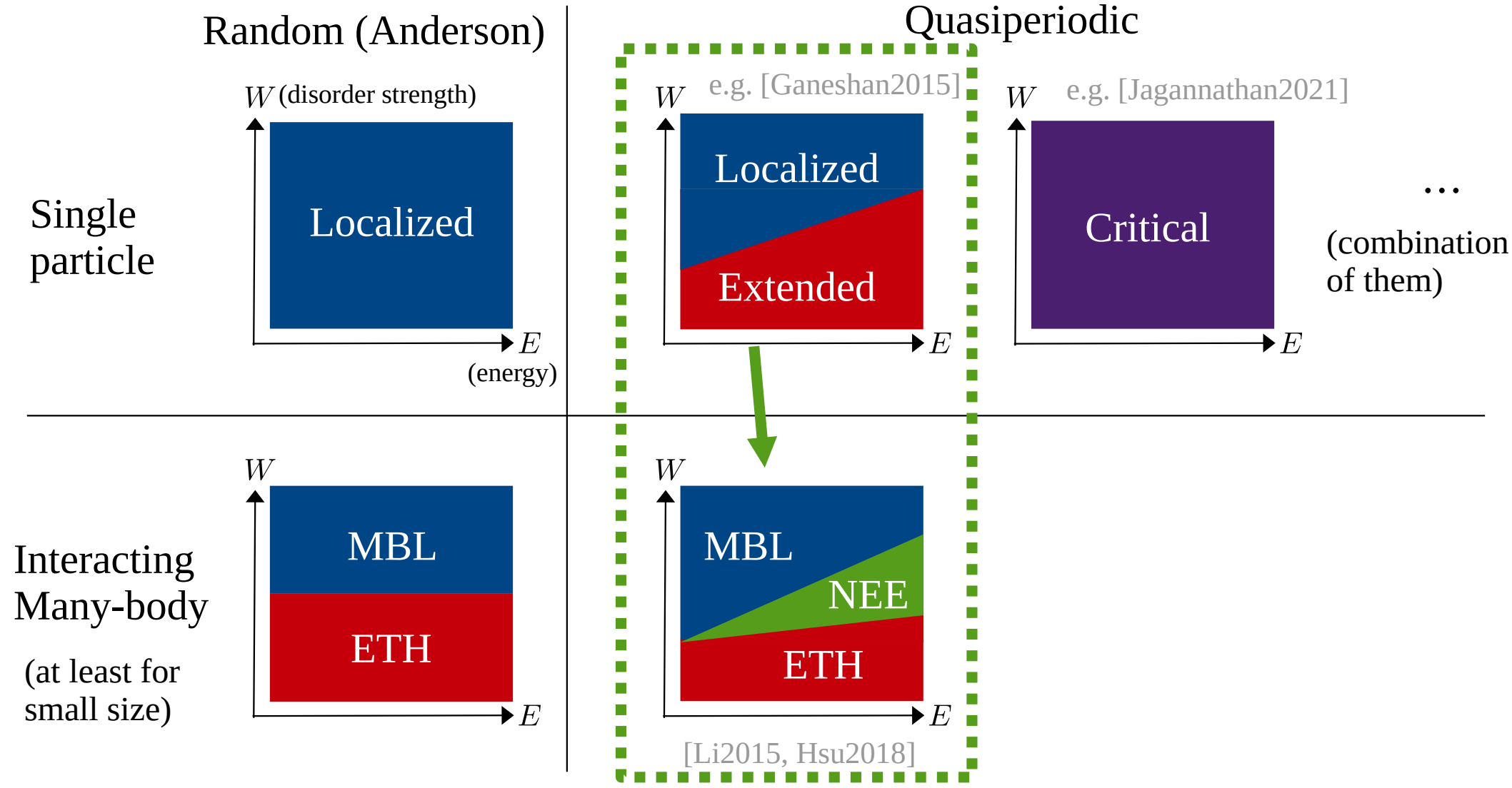
(combination of them)

Interacting
Many-body
(at least for
small size)



[Li2015, Hsu2018]

What's special about 1D quasiperiodic systems



Ganeshan-Pixley-Das Sarma (GPD) model

Ganeshan, Pixley, Das Sarma, PRL (2015)

$$H = \sum_{j=1}^{L-1} (S_j^x S_{j+1}^x + S_j^y S_{j+1}^y + U S_j^z S_{j+1}^z) + W \sum_{j=1}^L h_j S_j^z, \quad h_j = \frac{\cos(2\pi\varphi j + \phi)}{1 - \alpha \cos(2\pi\varphi j + \phi)}$$

Ganeshan-Pixley-Das Sarma (GPD) model

Ganeshan, Pixley, Das Sarma, PRL (2015)

$$H = \sum_{j=1}^{L-1} (S_j^x S_{j+1}^x + S_j^y S_{j+1}^y + U S_j^z S_{j+1}^z) + W \sum_{j=1}^L h_j S_j^z, \quad h_j = \frac{\cos(2\pi\varphi j + \phi)}{1 - \alpha \cos(2\pi\varphi j + \phi)}$$

↑
Spin-1/2

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Spin-1/2

Interaction strength
 $U=1$

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Spin-1/2

Interaction strength
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Initial phase (averaged over)

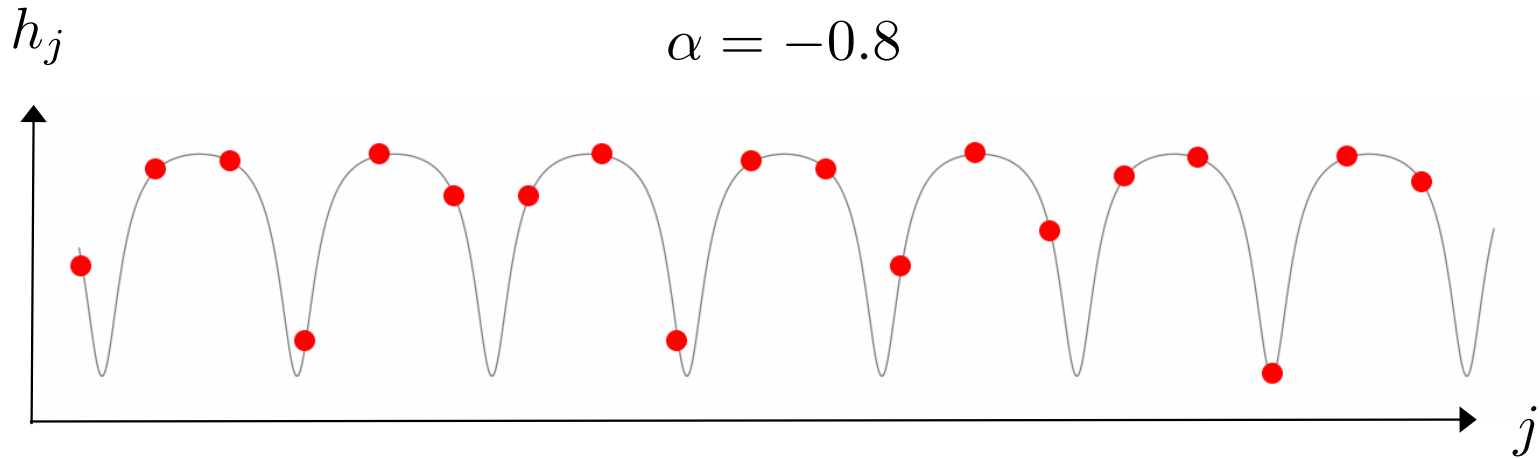
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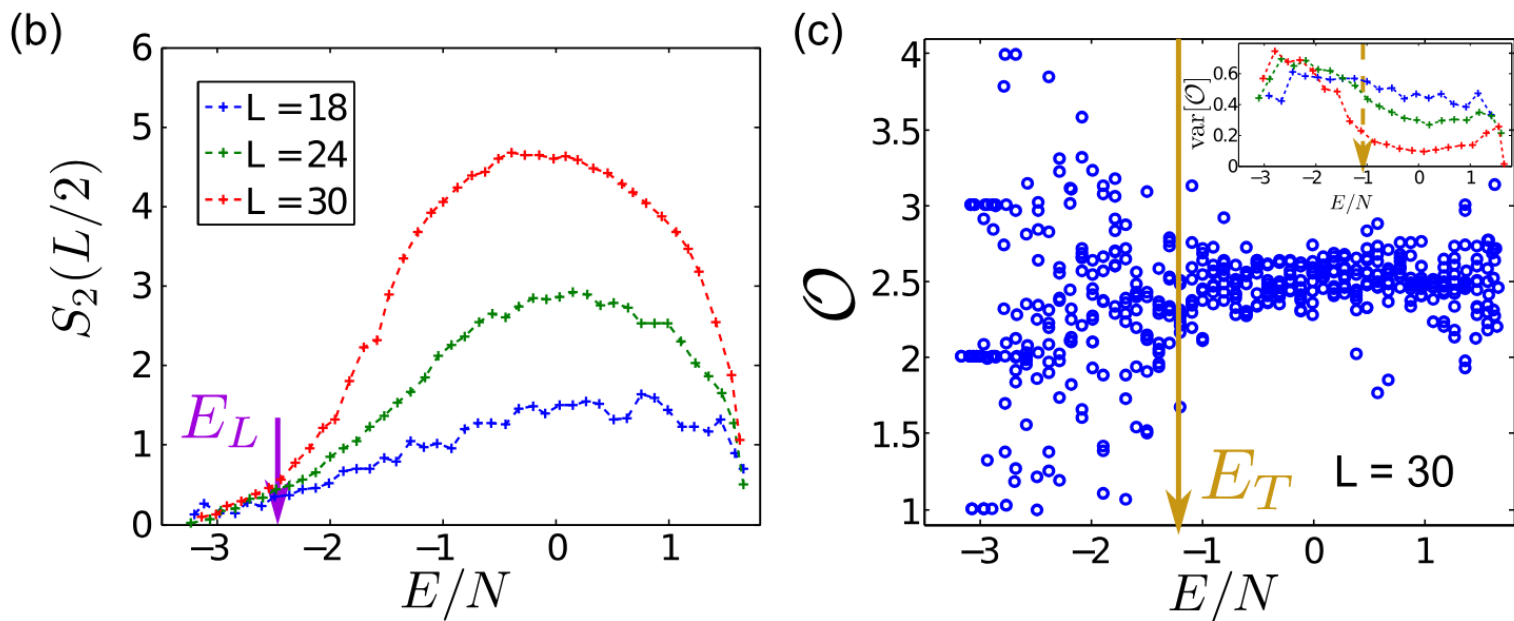


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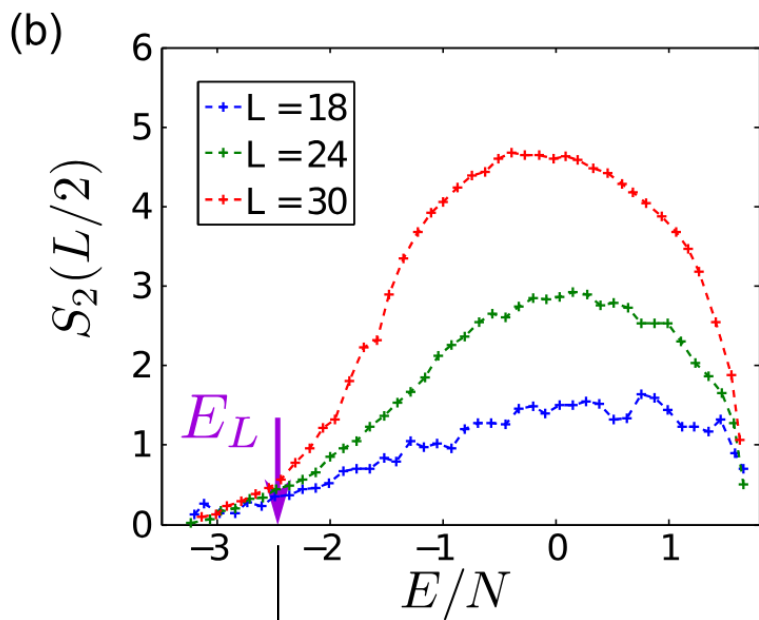
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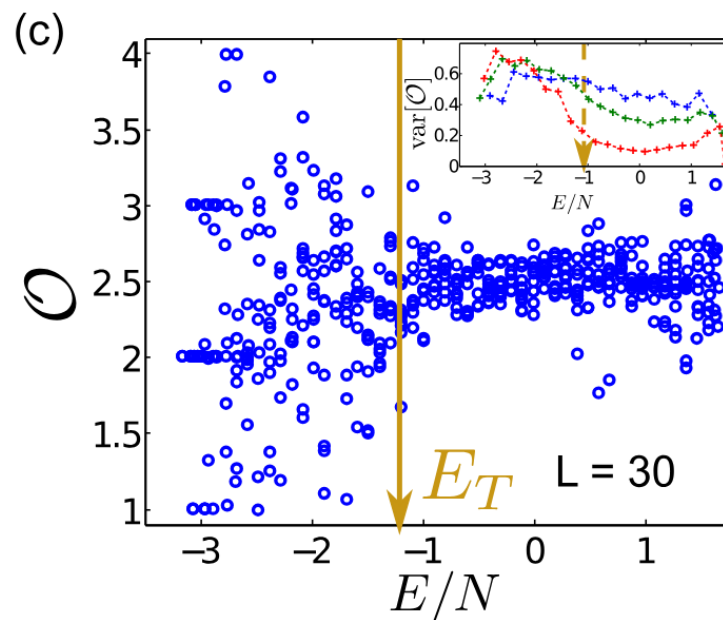
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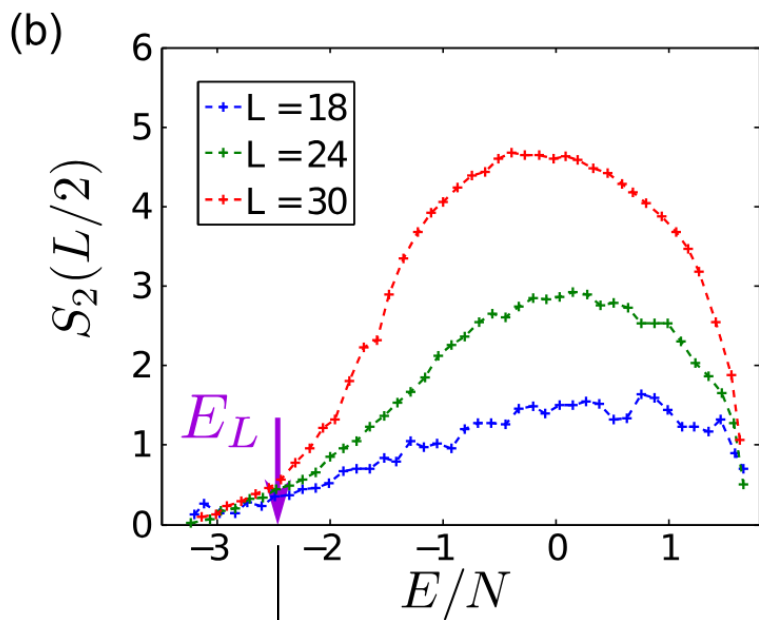
Area-law
(complete LIOMs)

Volume-law entanglement
("extended")



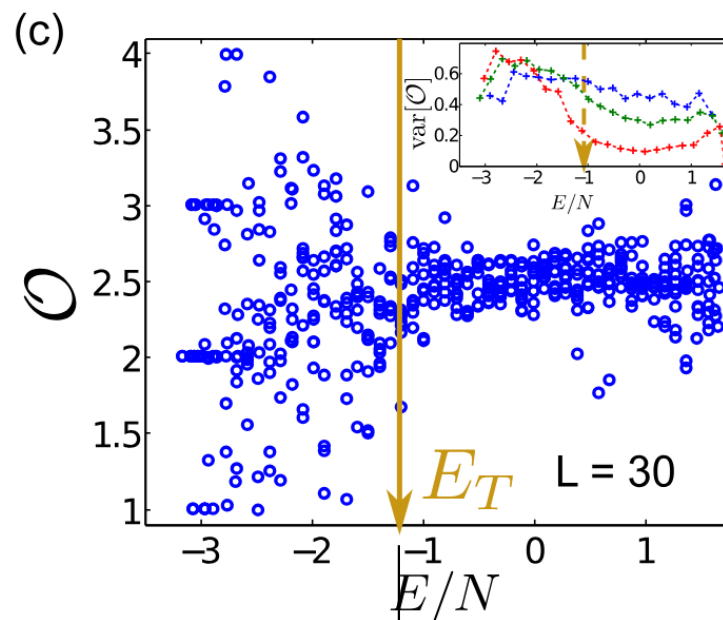
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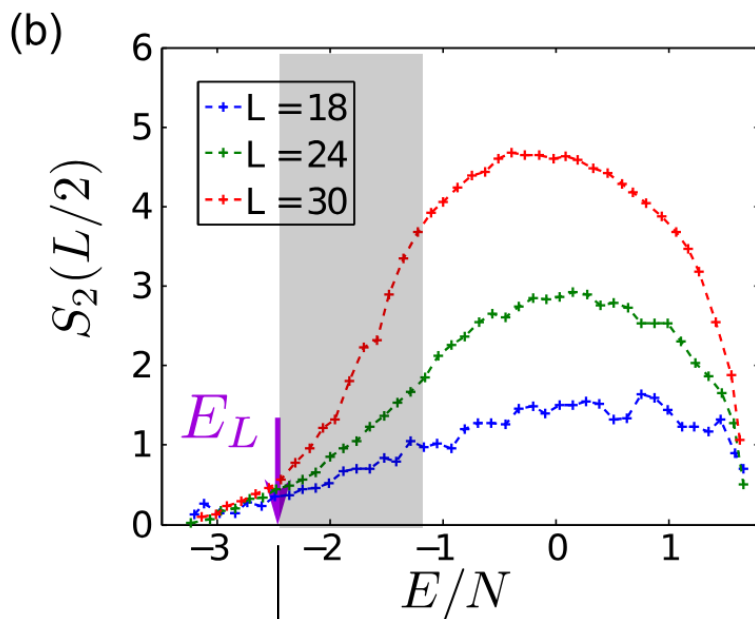


Violates ETH
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Satisfies ETH

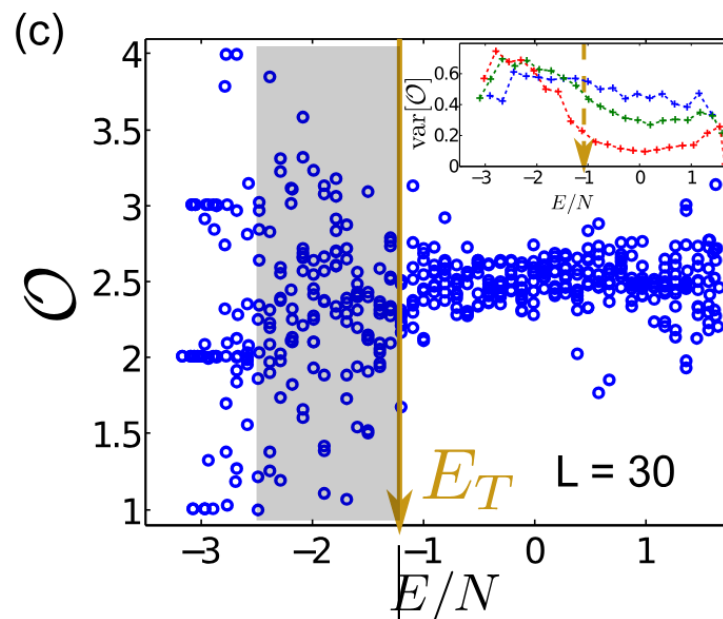
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Spin-spin correlator in the prethermal regime

$$C(E, t) = \left[\langle \psi_E | S_{L/2}^z(t) S_{L/2}^z(0) | \psi_E \rangle \right] - \left[\langle \psi_E | S_{L/2}^z(0) | \psi_E \rangle \right]^2$$

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Model of random accidental resonance predicts:

Long et al. PRL '23; Long et al. SciPost '23

$$C(E, t) \approx A(E) \exp \left[- \left(\frac{t}{\tau(E)} \right)^{\beta(E)} \right]$$

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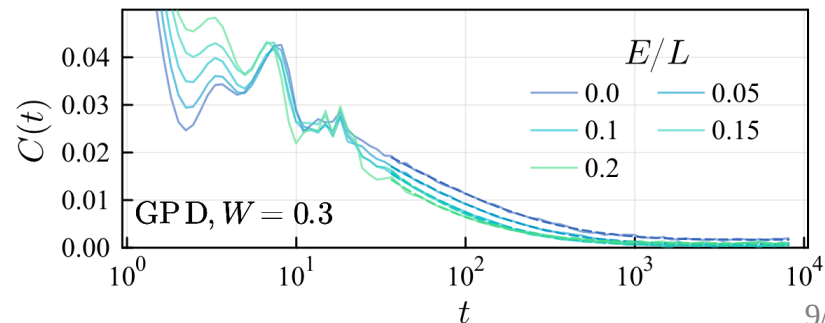
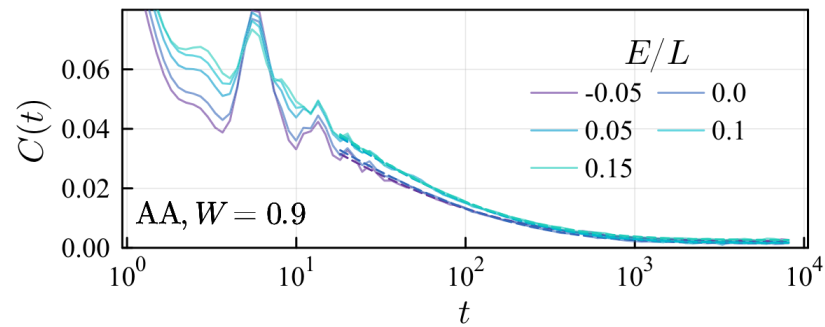
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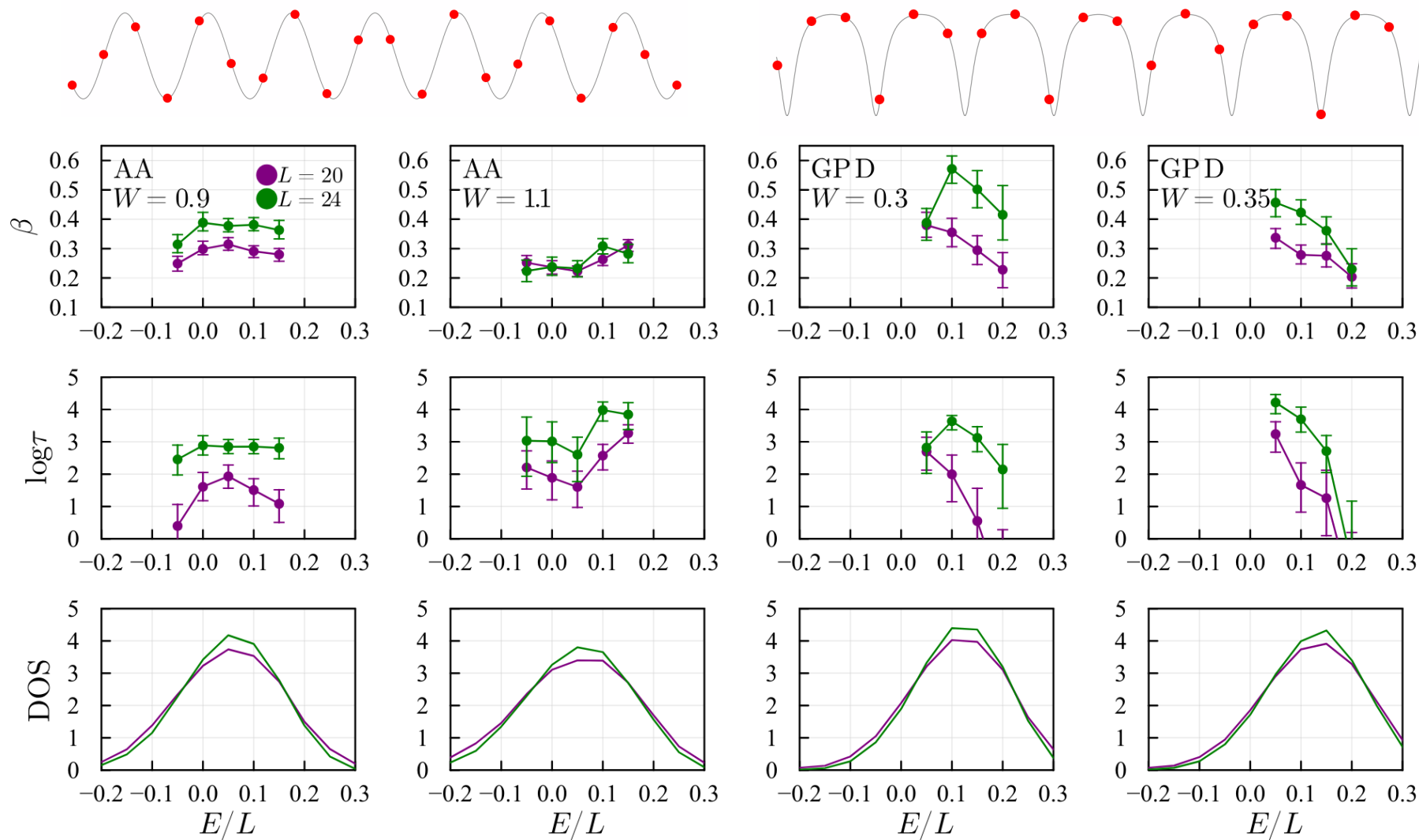


Stretched exponential fit

YTT, Long, Das Sarma PRB (2024)

$\alpha = 0$

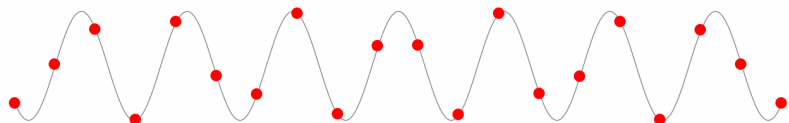
$\alpha = -0.8$



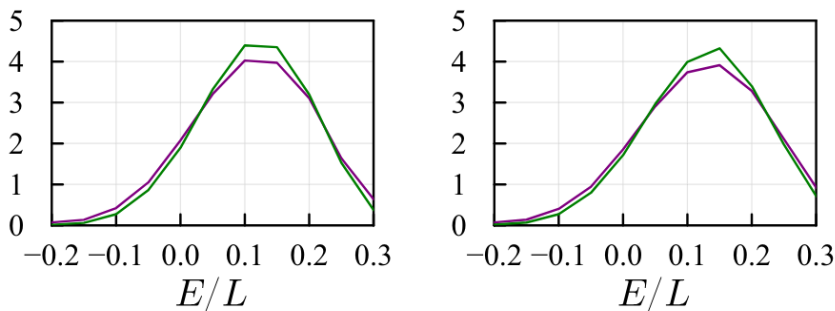
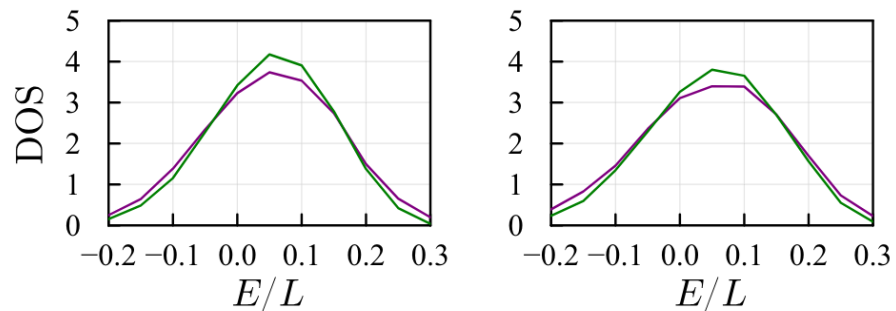
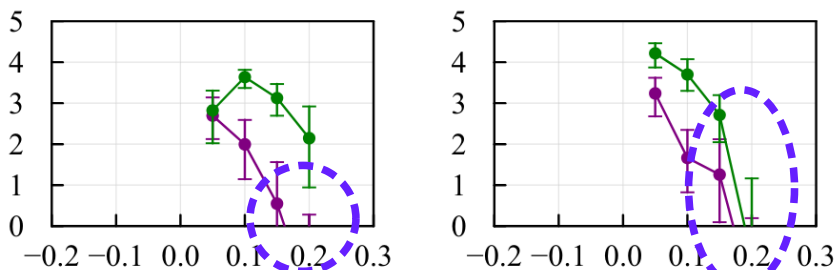
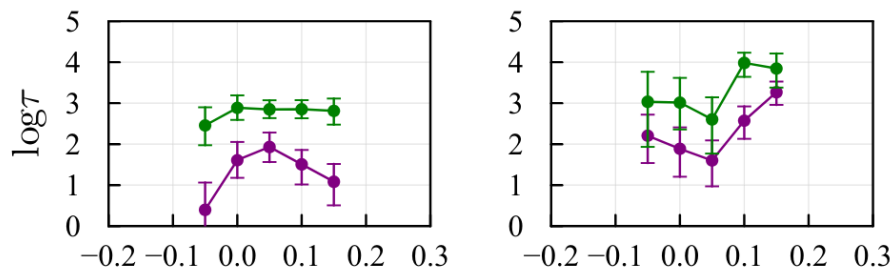
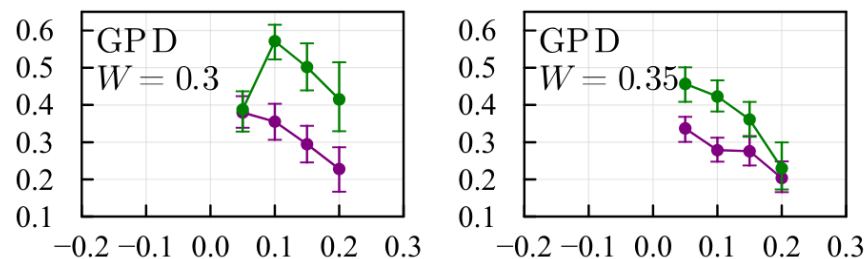
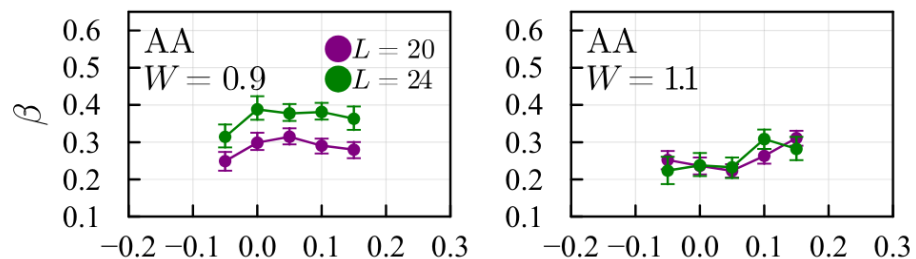
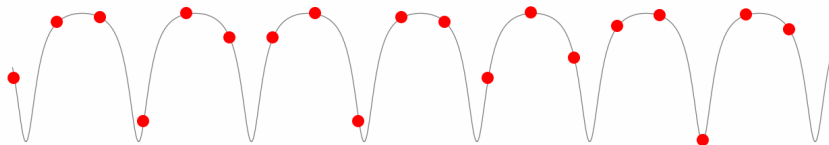
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YTT, Long, Das Sarma PRB (2024)

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Resonance model failed!

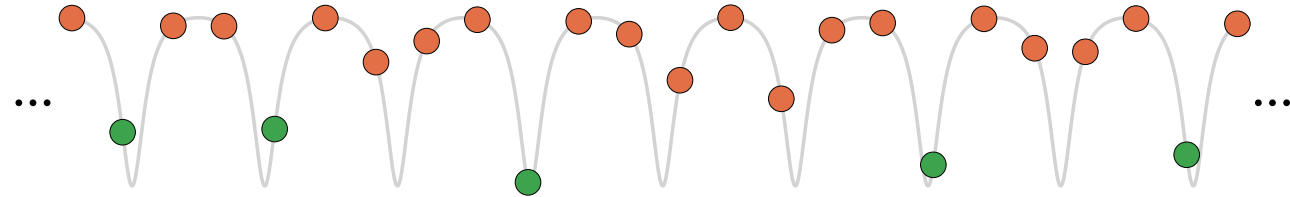
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Effects of regularly spaced deep wells

YTT, Long, Das Sarma PRB (2024)

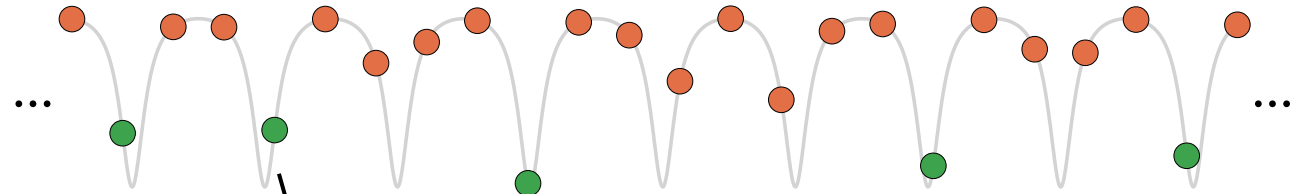
GPD potential:



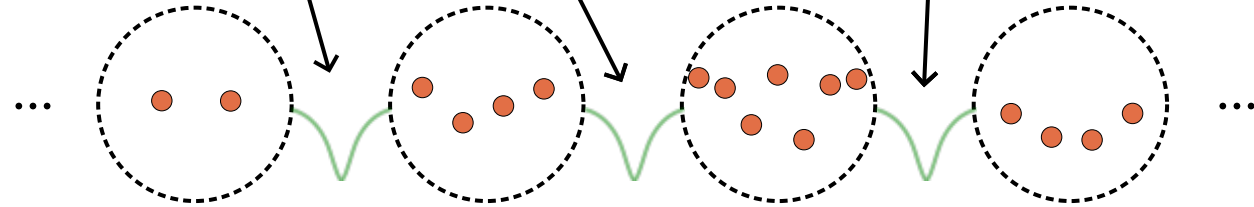
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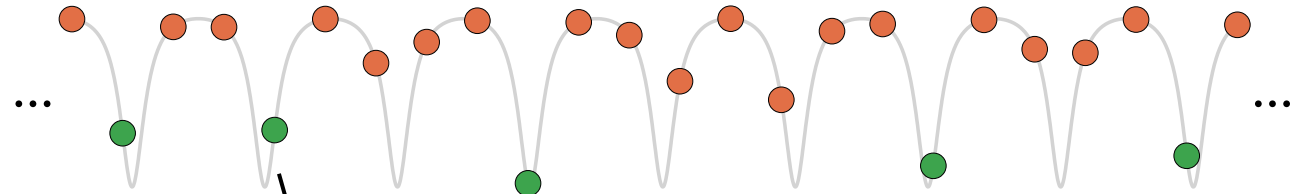
Supersites:



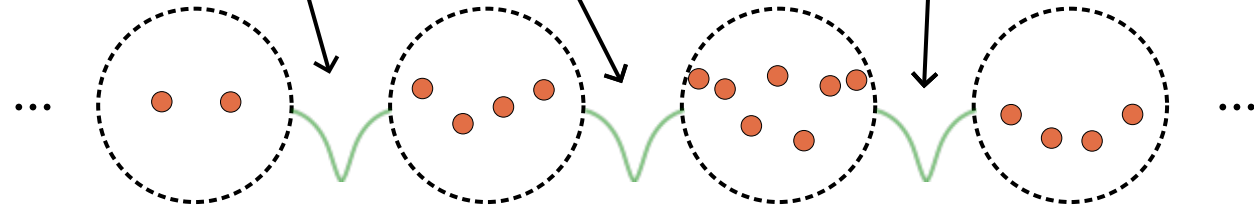
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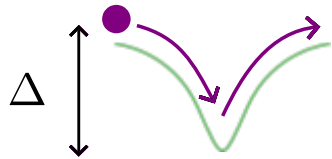
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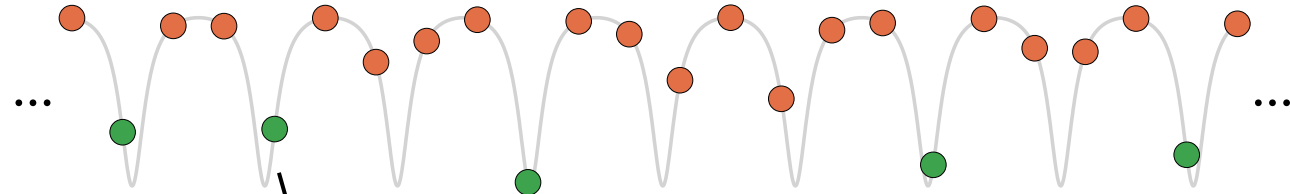
Inter-supersite hopping:



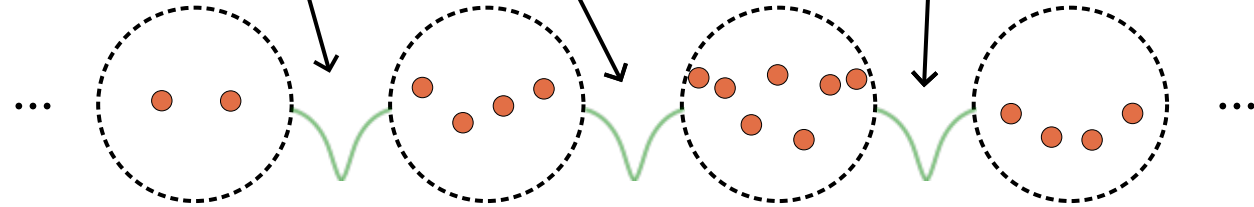
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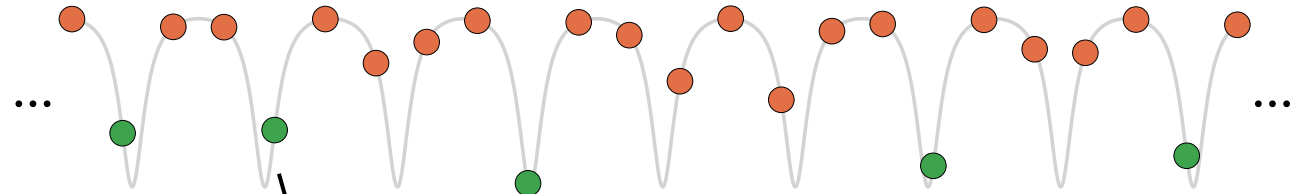
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$$\Delta \sim \frac{J^2}{\Delta}$$

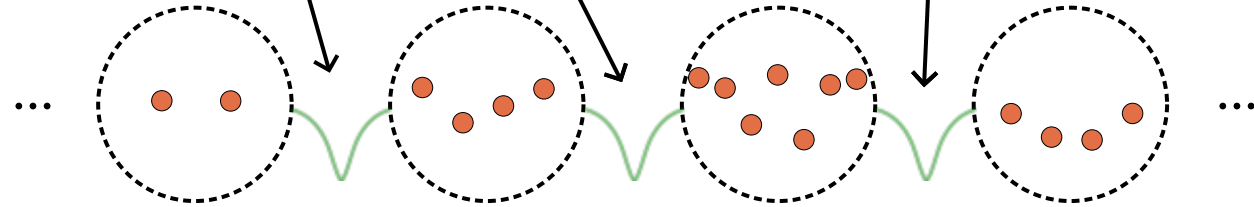
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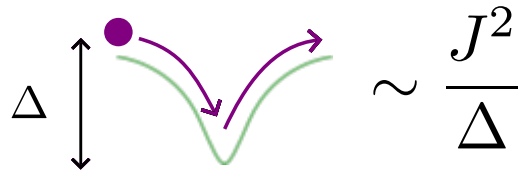
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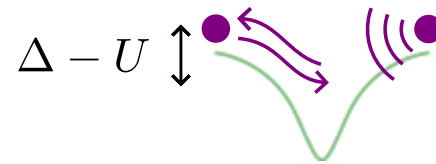
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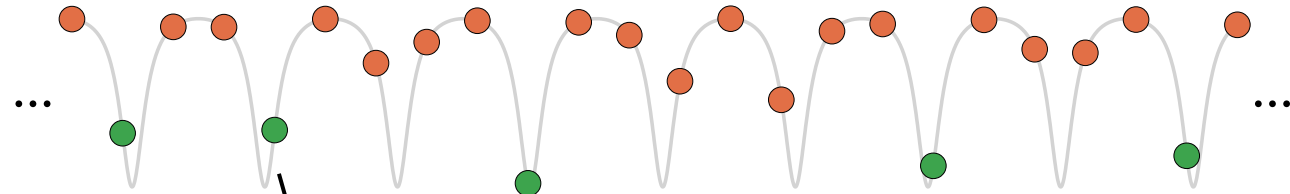
Inter-supersite interaction:



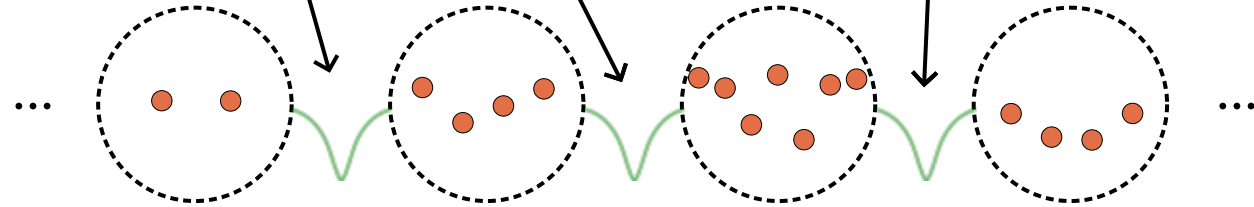
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GPD potential:



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Inter-supersite hopping:

The diagram illustrates a particle (purple dot) moving from one well to an adjacent one, indicated by a purple arrow. A vertical double-headed arrow labeled Δ represents the energy difference between the wells. The hopping process is associated with the equation:

$$\sim \frac{J^2}{\Delta}$$

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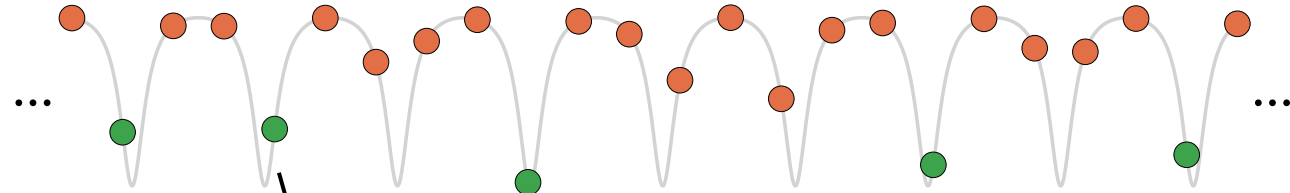
The diagram shows two particles (purple dots) in the same well, with a purple arrow indicating interaction. A vertical double-headed arrow labeled $\Delta - U$ represents the energy difference. The interaction process is associated with the equation:

$$\sim \frac{J^2}{\Delta - U} - \frac{J^2}{\Delta}$$

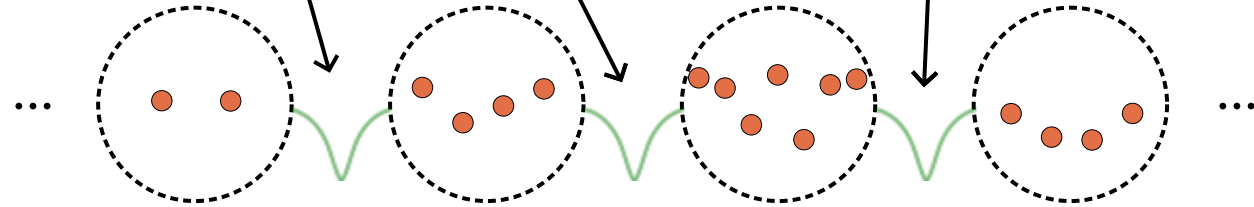
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GPD potential:



Supersites:



Inter-supersite hopping:

The diagram shows a purple dot representing a particle in a potential well. A double-headed purple arrow indicates the hopping between two adjacent wells. A vertical double-headed arrow to the left is labeled Δ , representing the energy gap. The hopping is given by the equation:

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Inter-supersite interaction:

The diagram shows two purple dots representing particles in a potential well. Double-headed purple arrows between them indicate interaction. A vertical double-headed arrow to the left is labeled $\Delta - U$, representing the energy gap minus the interaction energy. The interaction is given by the equation:

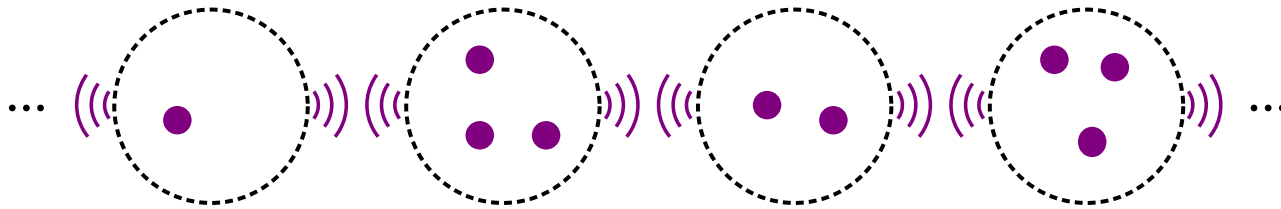
$$\sim \frac{J^2}{\Delta - U} - \frac{J^2}{\Delta}$$

Can be much larger!

Non-ergodic extended (NEE) behavior

YTT, Long, Das Sarma PRB (2024)

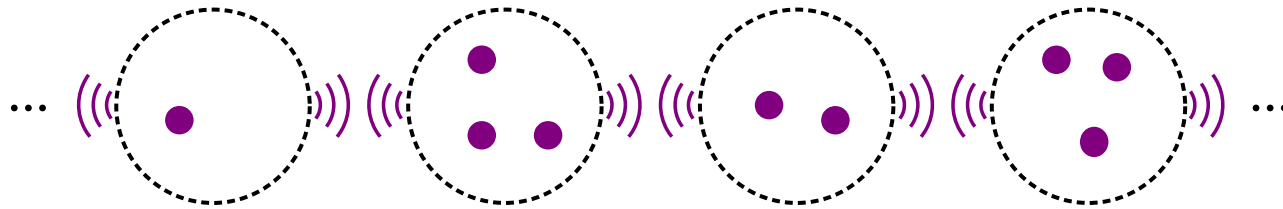
Approximate inter-supersite hopping as 0 but interaction being finite.



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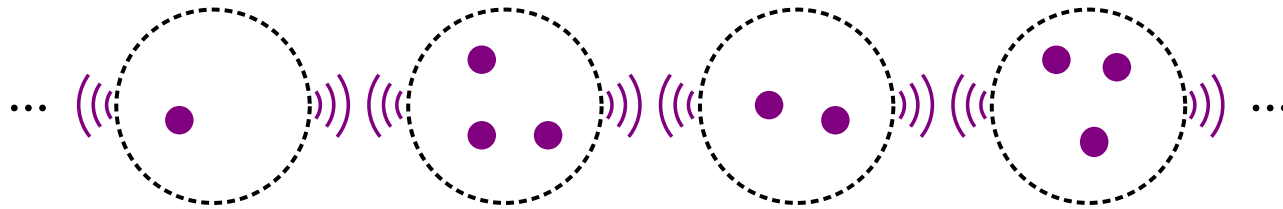


- For each supersite, we have one LIOM (the number of particles in it).

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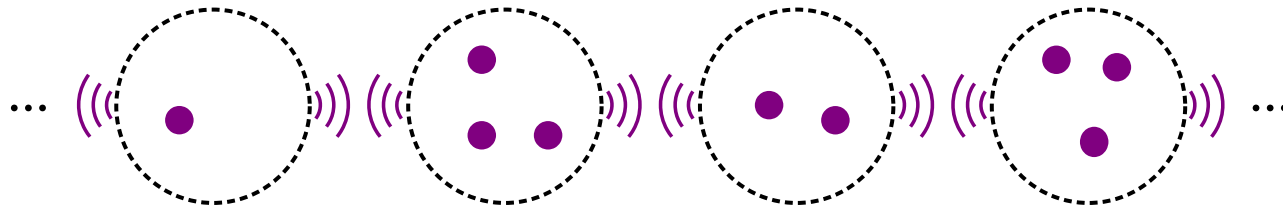


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- But this does not completely describe the state.

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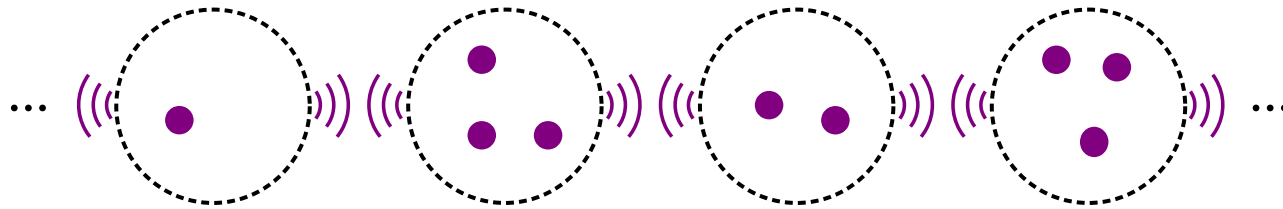


- For each supersite, we have one LIOM (the number of particles in it).
- But this does not completely describe the state.
- The remaining DOFs in the super-chain can still thermalize.

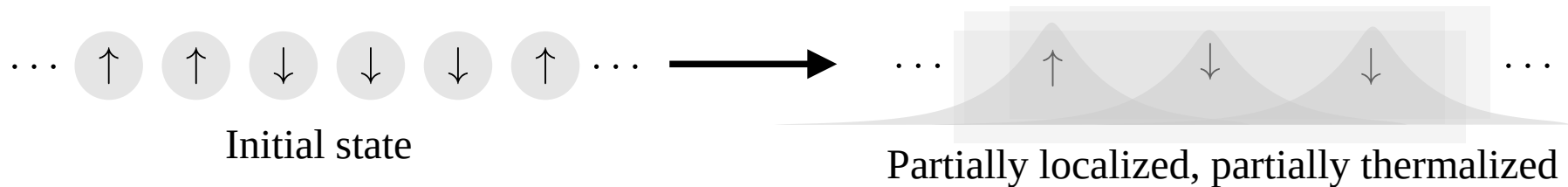
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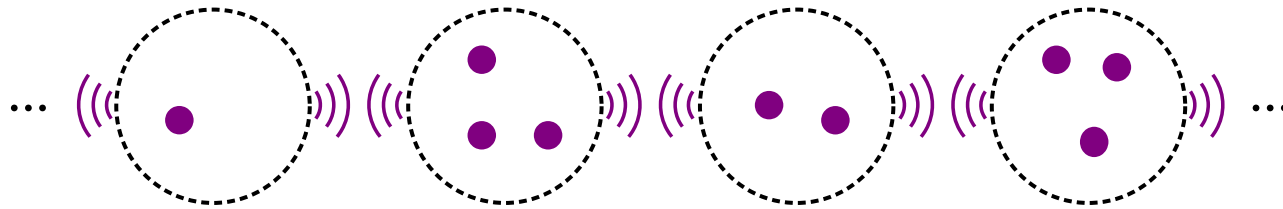
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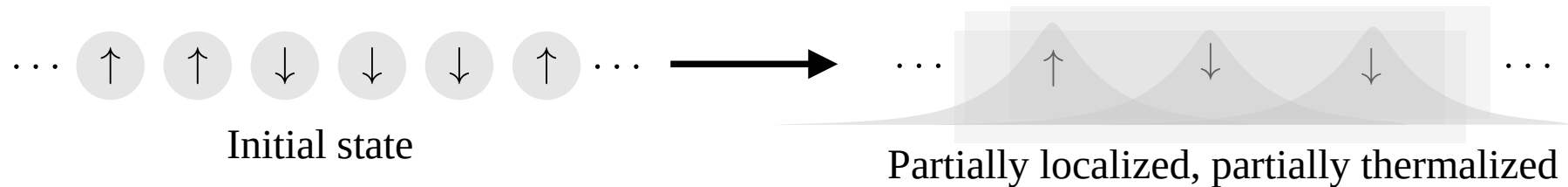
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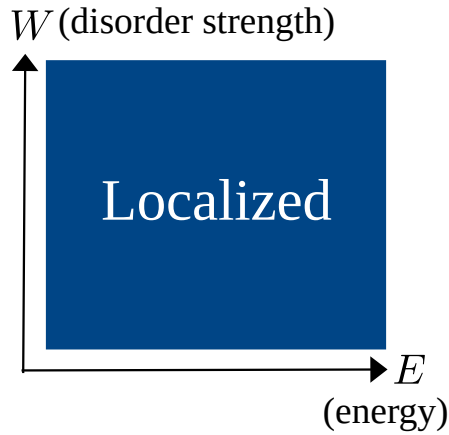
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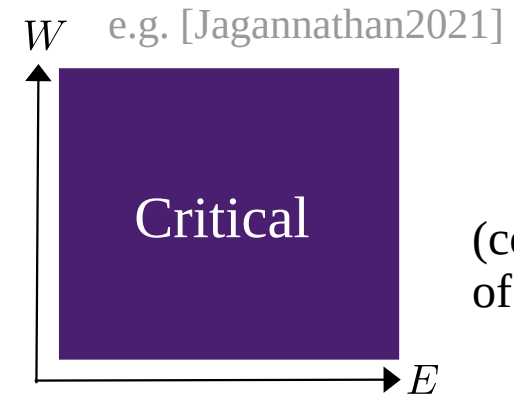
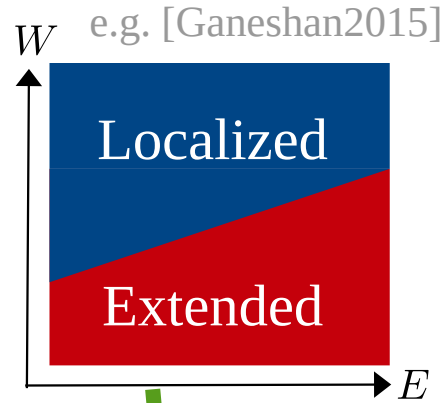
Expected to be a finite-size/time behavior (eventually everything thermalizes)

What's special about 1D quasiperiodic systems

Random (Anderson)



Quasiperiodic

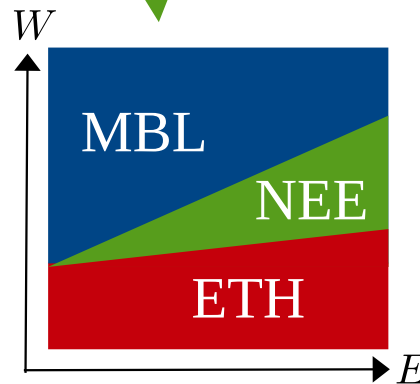
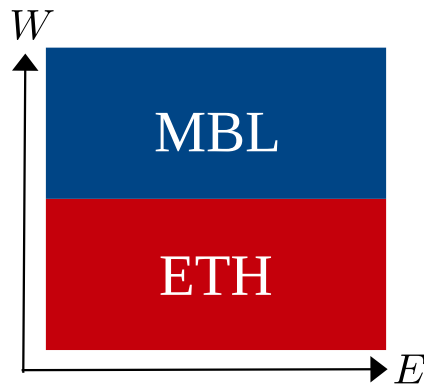


...

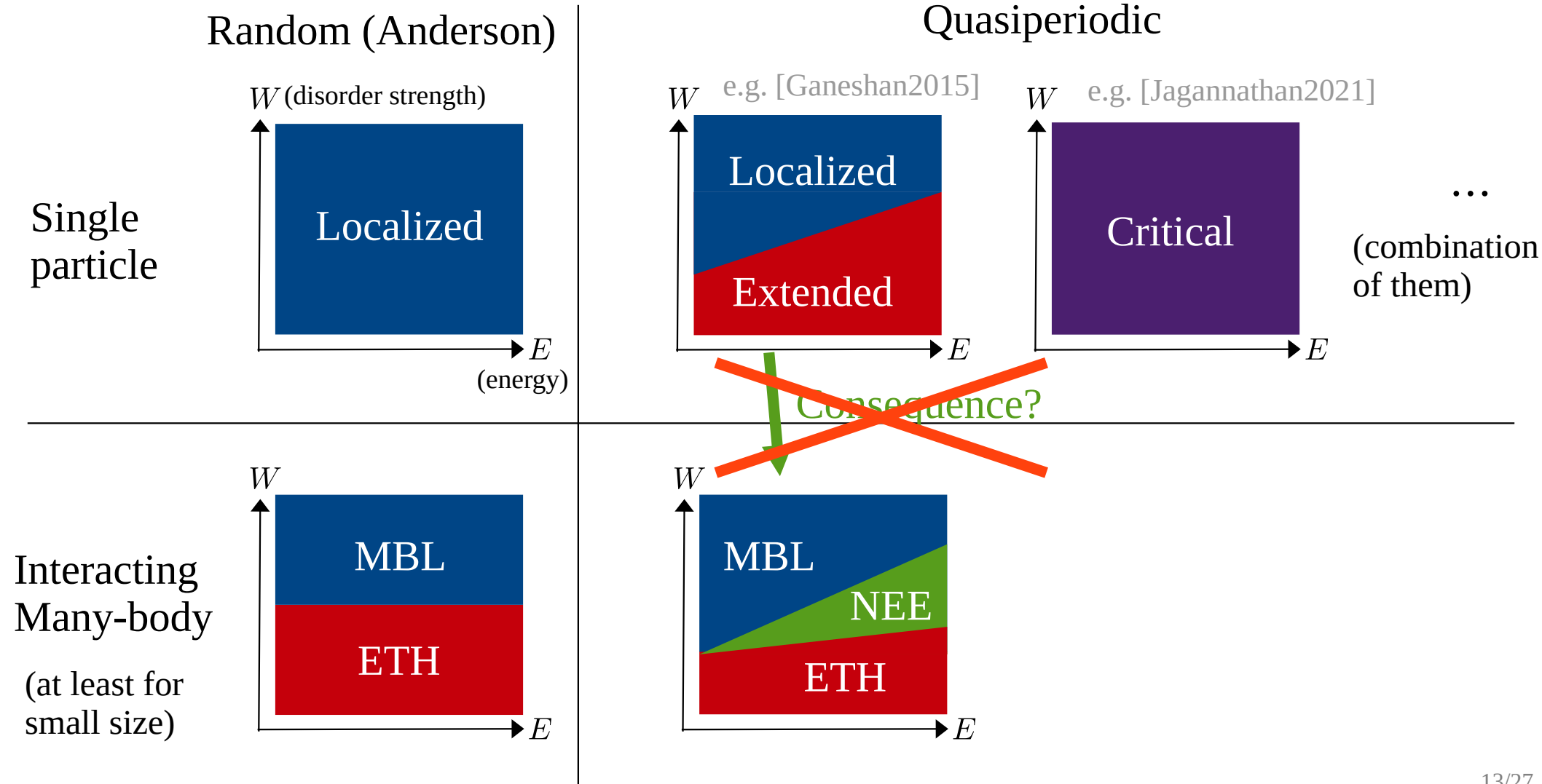
(combination of them)

Consequence?

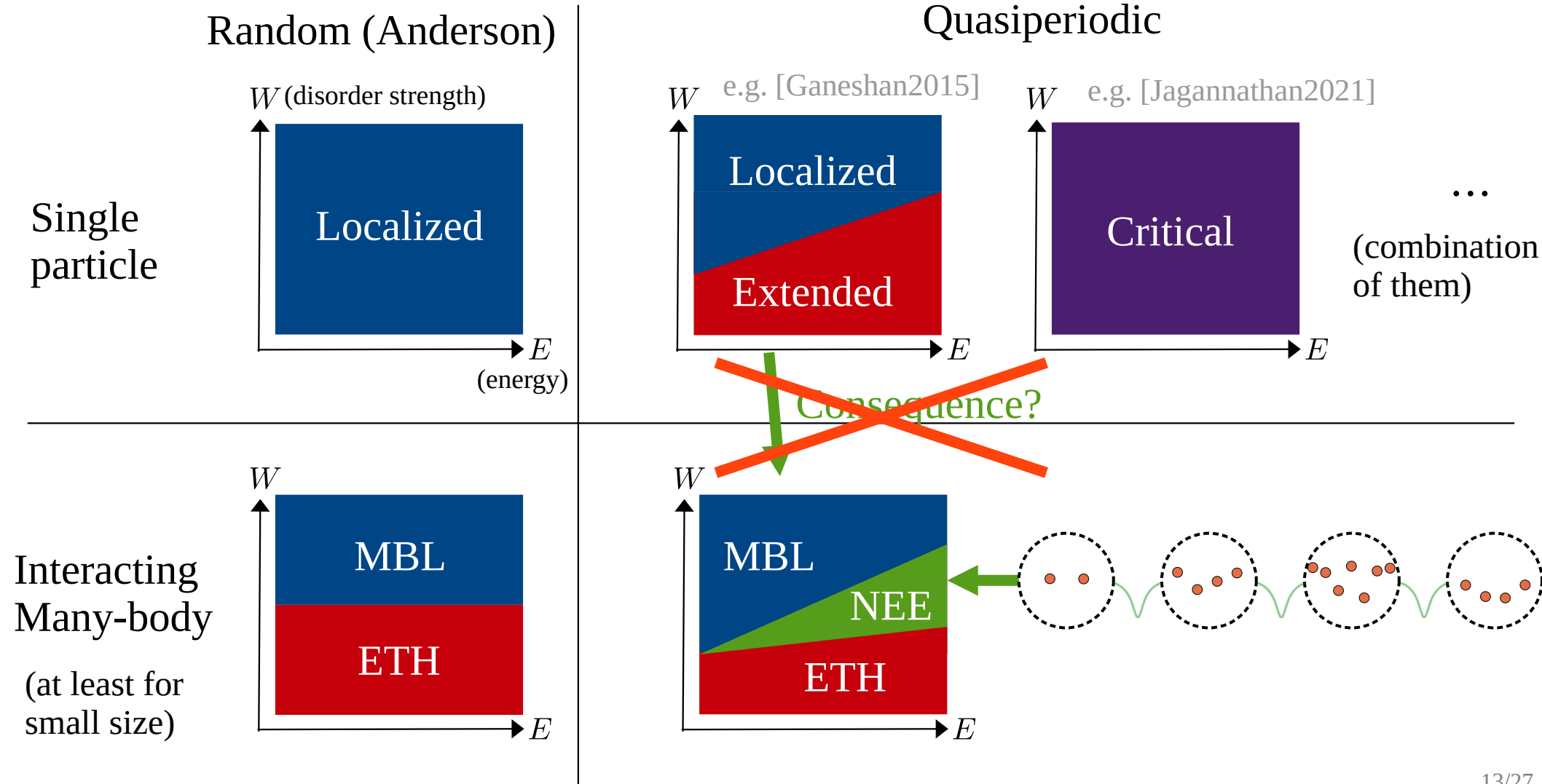
Interacting
Many-body
(at least for
small size)



What's special about 1D quasiperiodic systems

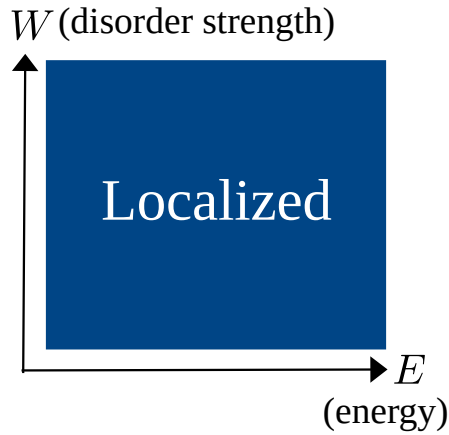


What's special about 1D quasiperiodic systems



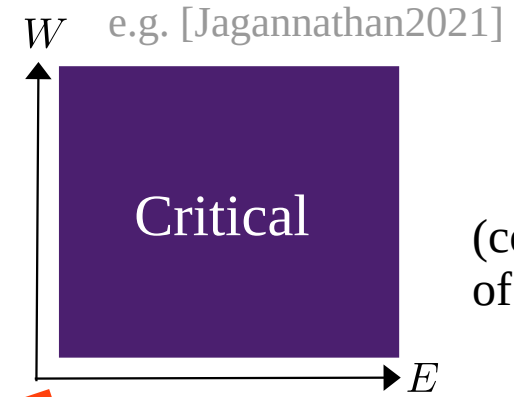
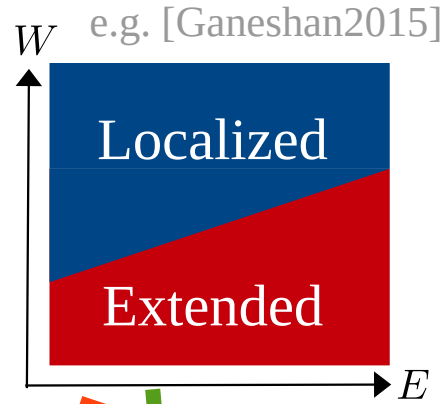
What's special about 1D quasiperiodic systems

Random (Anderson)



Single
particle

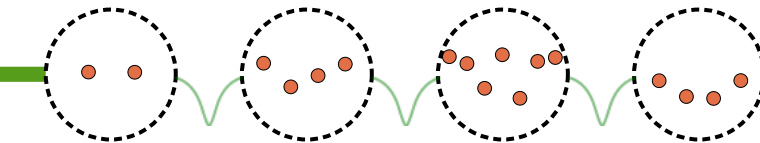
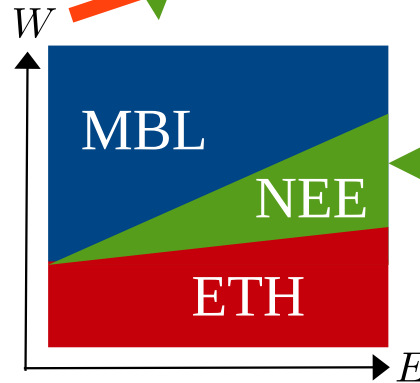
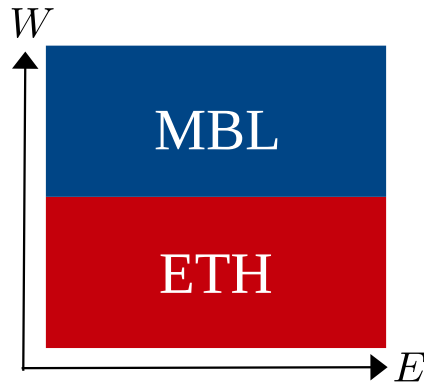
Quasiperiodic



...
(combination
of them)

Consequence?

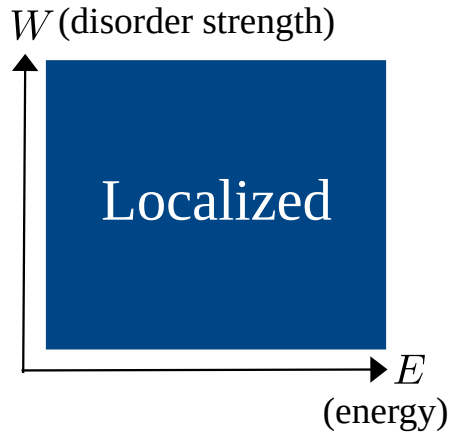
Interacting
Many-body
(at least for
small size)



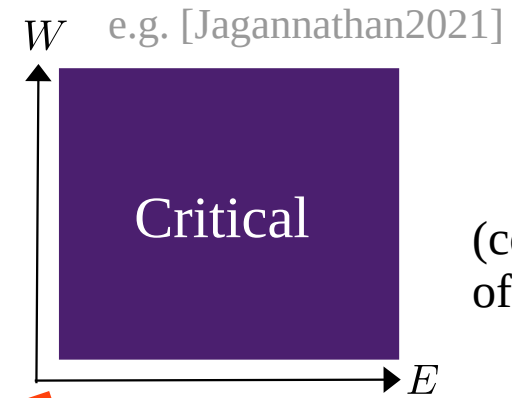
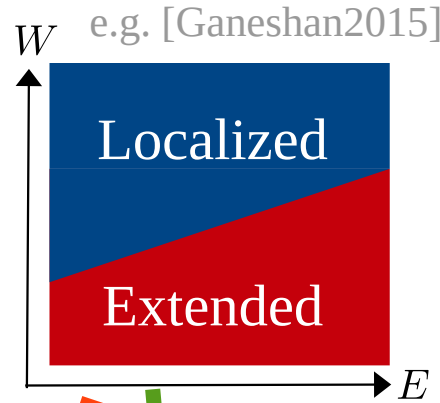
Due to spatial bottleneck at finite time

What's special about 1D quasiperiodic systems

Random (Anderson)



Quasiperiodic

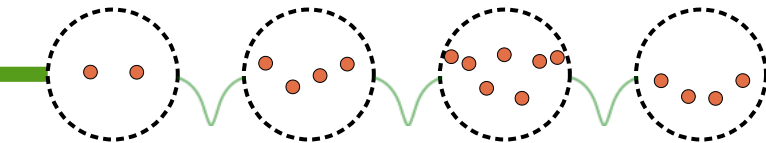
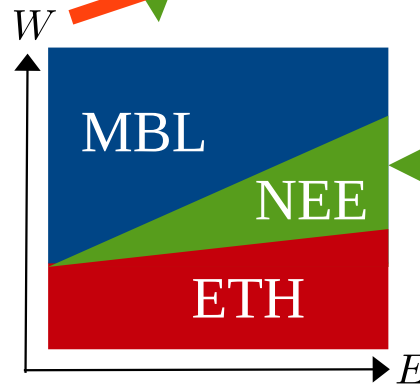
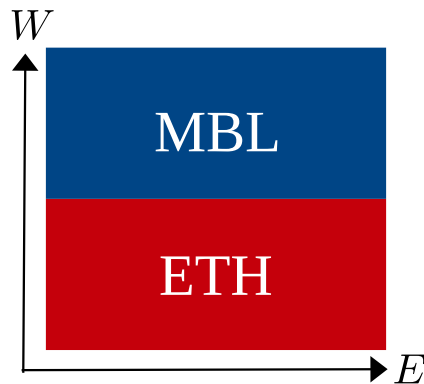


...

(combination of them)

Consequence?

Interacting Many-body
(at least for small size)



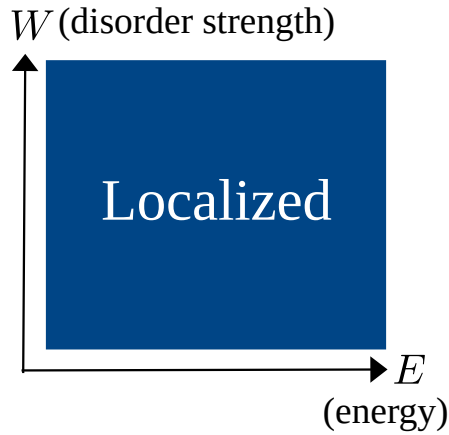
Due to spatial bottleneck at finite time
Not a thermodynamic phase

Outline

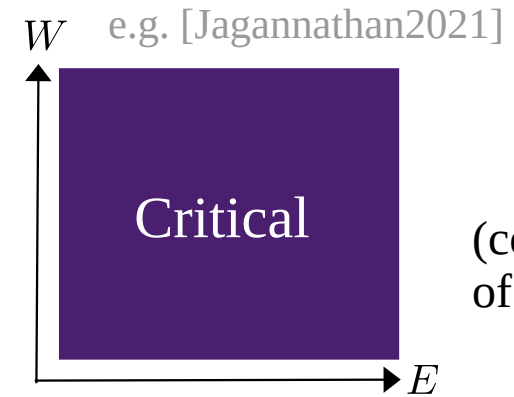
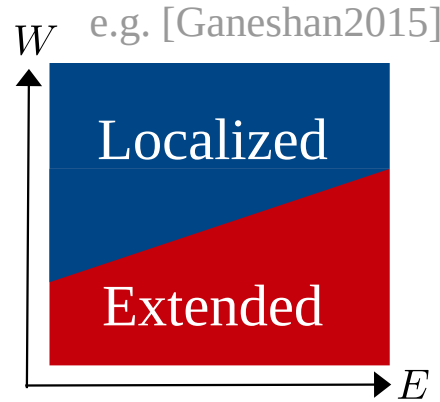
- Introduction
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 - Asymptotically solvable model

What's special about 1D quasiperiodic systems

Random (Anderson)



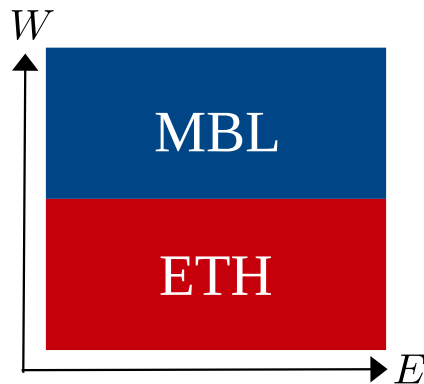
Quasiperiodic



...

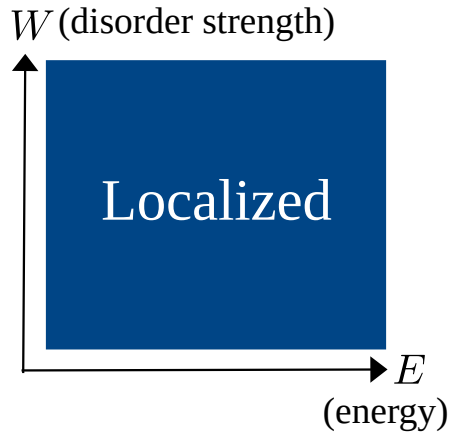
(combination of them)

Interacting
Many-body
(at least for
small size)

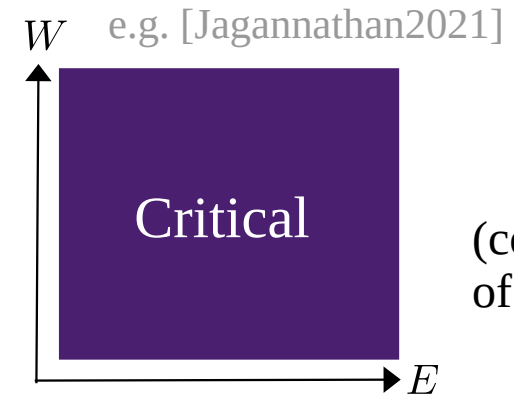
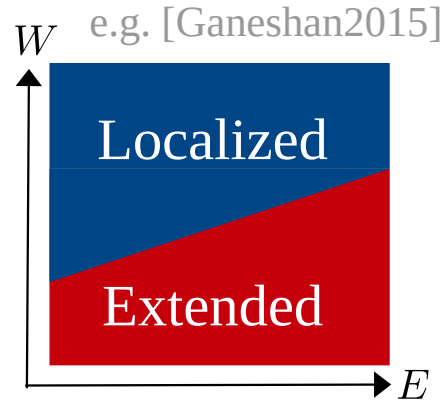


What's special about 1D quasiperiodic systems

Random (Anderson)



Quasiperiodic

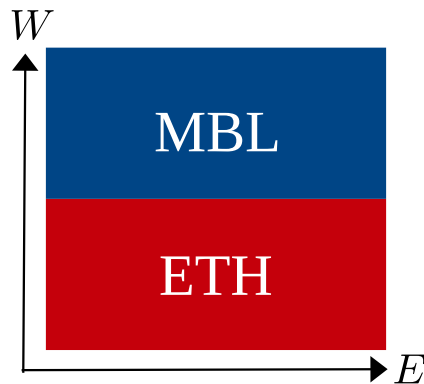


...

(combination of them)

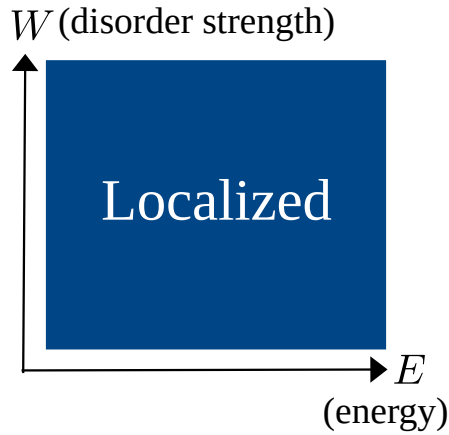


Interacting
Many-body
(at least for
small size)

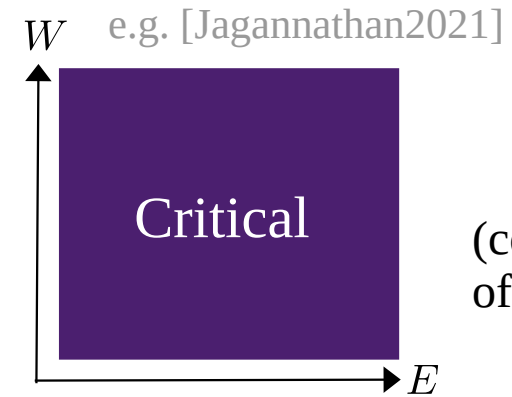
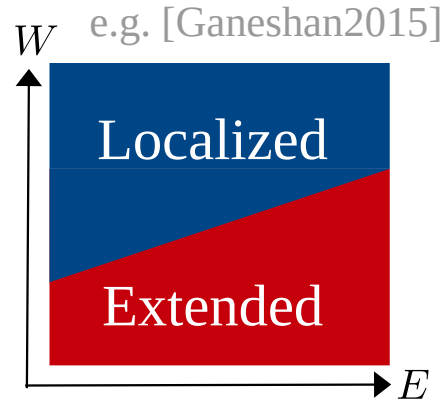


What's special about 1D quasiperiodic systems

Random (Anderson)



Quasiperiodic



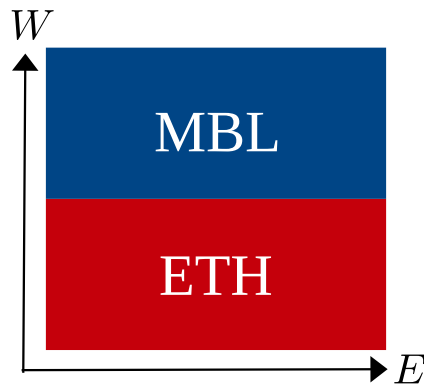
...

(combination of them)



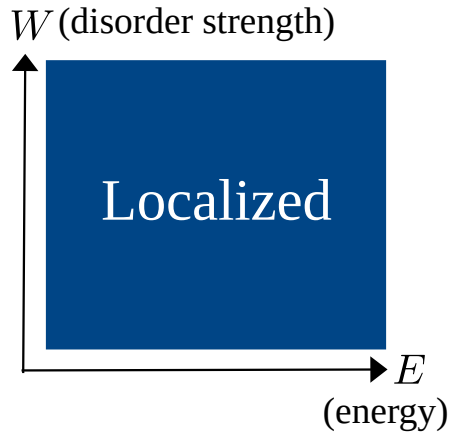
Becomes MBL?

Interacting
Many-body
(at least for
small size)

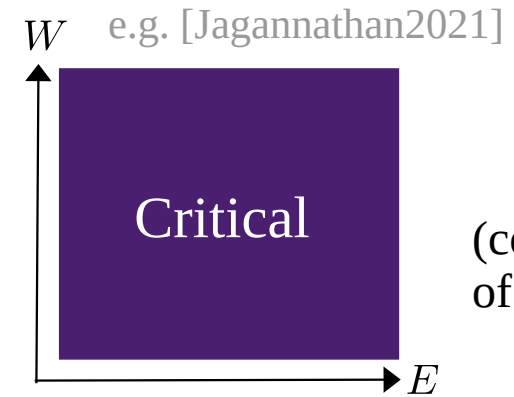
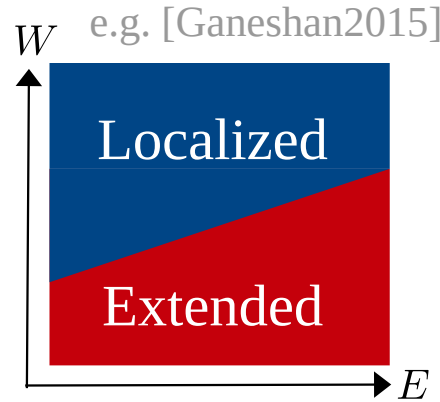


What's special about 1D quasiperiodic systems

Random (Anderson)



Quasiperiodic



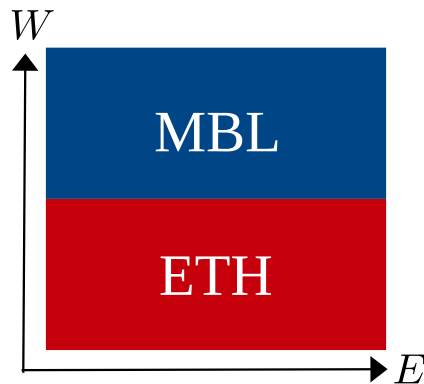
...

(combination of them)



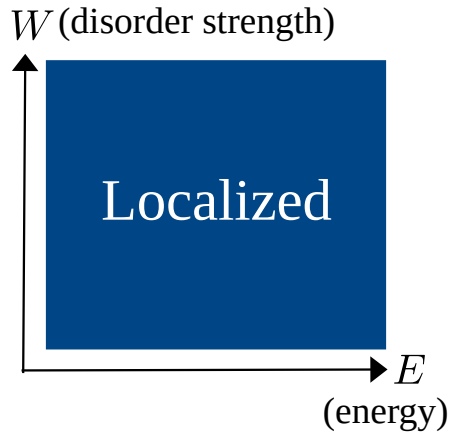
Becomes MBL?
Becomes ETH?

Interacting
Many-body
(at least for
small size)

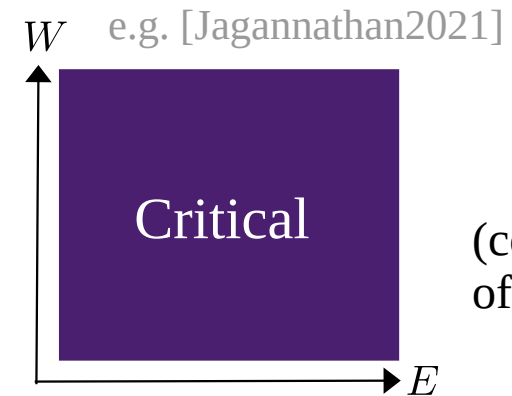
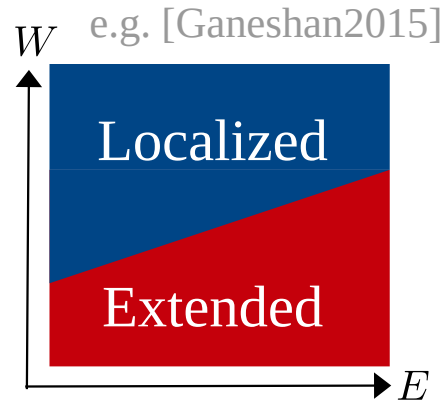


What's special about 1D quasiperiodic systems

Random (Anderson)



Quasiperiodic



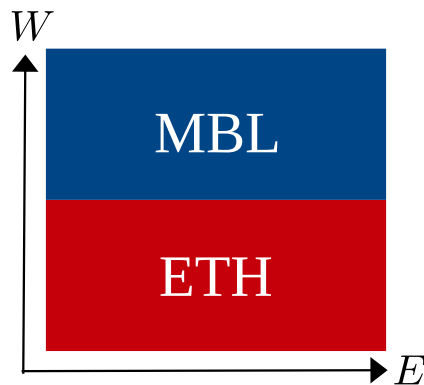
...

(combination of them)



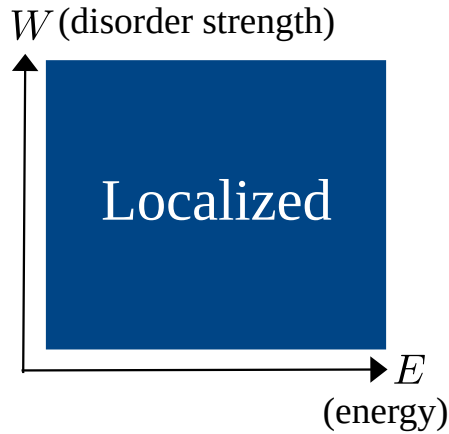
Becomes MBL?
Becomes ETH?
Many-body critical phase?

Interacting
Many-body
(at least for
small size)

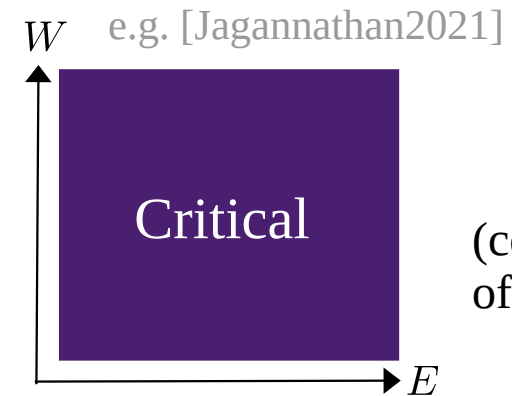
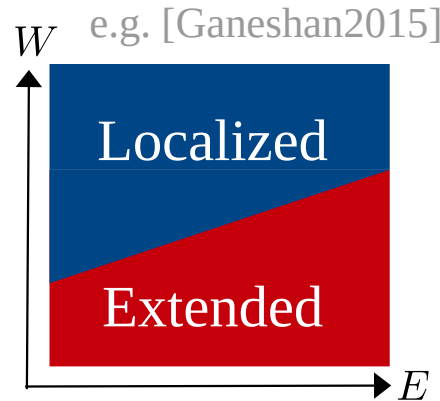


What's special about 1D quasiperiodic systems

Random (Anderson)



Quasiperiodic



...

(combination of them)

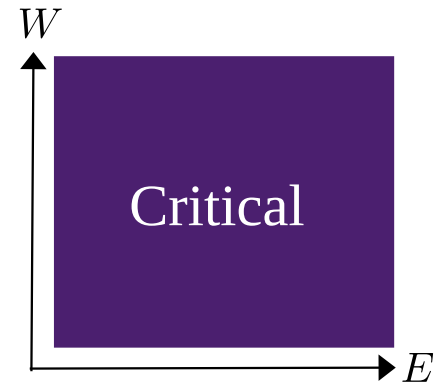
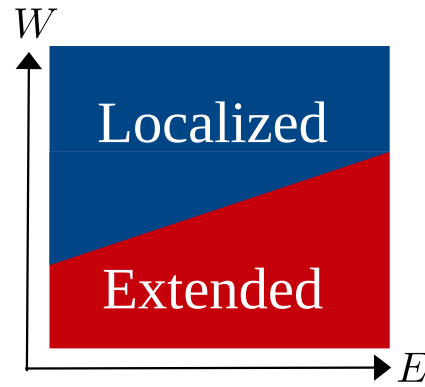


Becomes MBL?
Becomes ETH?
Many-body critical phase?

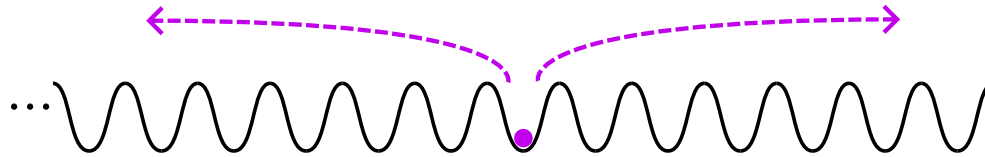
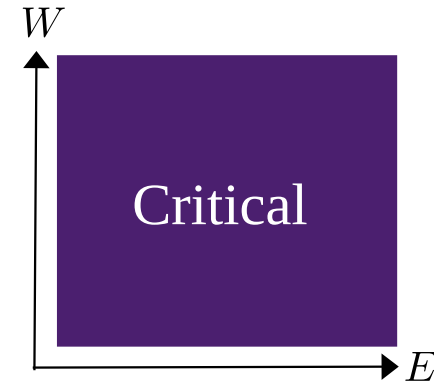
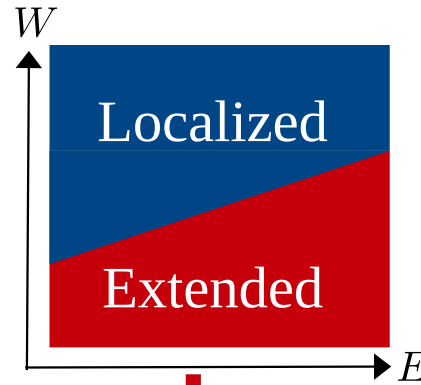
Existing literature are contradictory
e.g. [Mace2019, Varma2019, Wang2021]

And often based on small size numerics

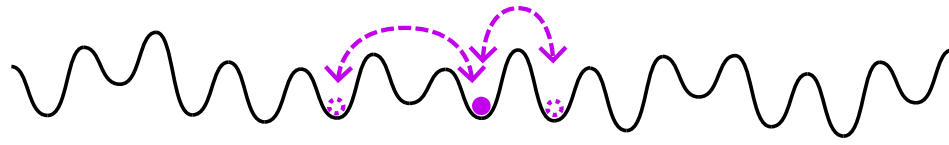
Single
particle



Single
particle

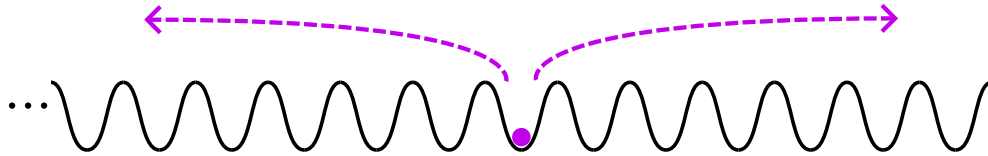
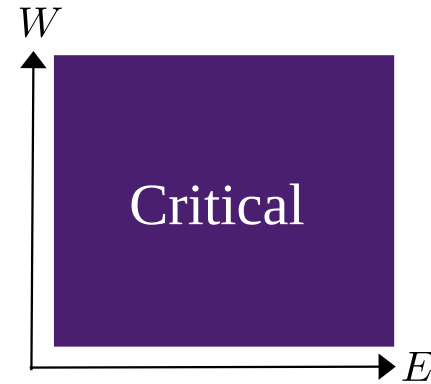
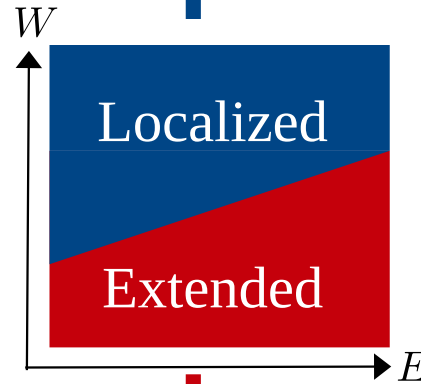


Travels further and further away
and never comes back

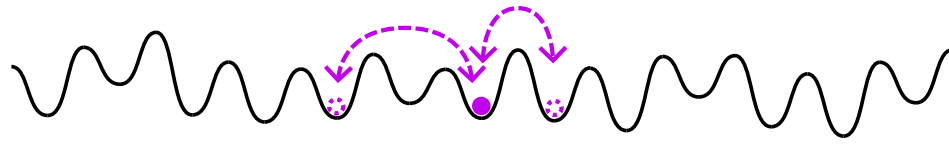


Comes back regularly
and does not go arbitrarily far away

Single
particle

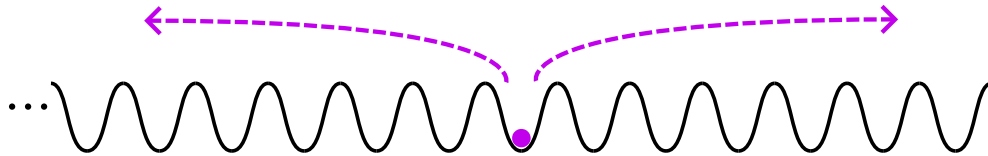
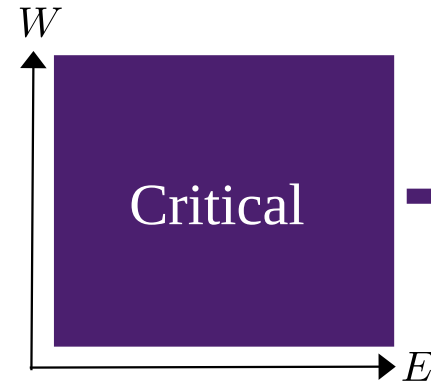
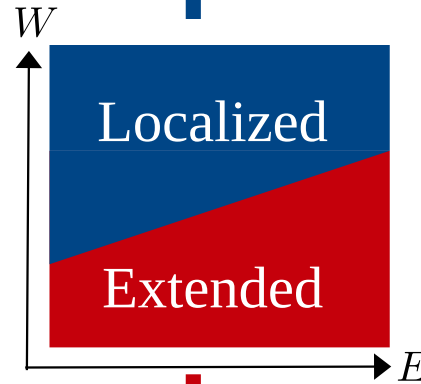


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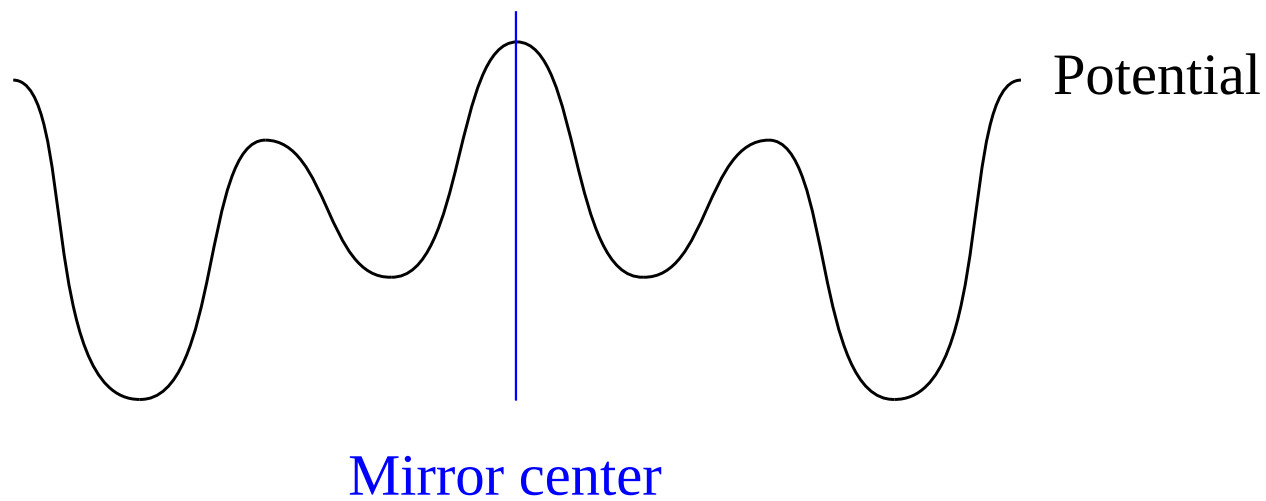


Travels further and further away
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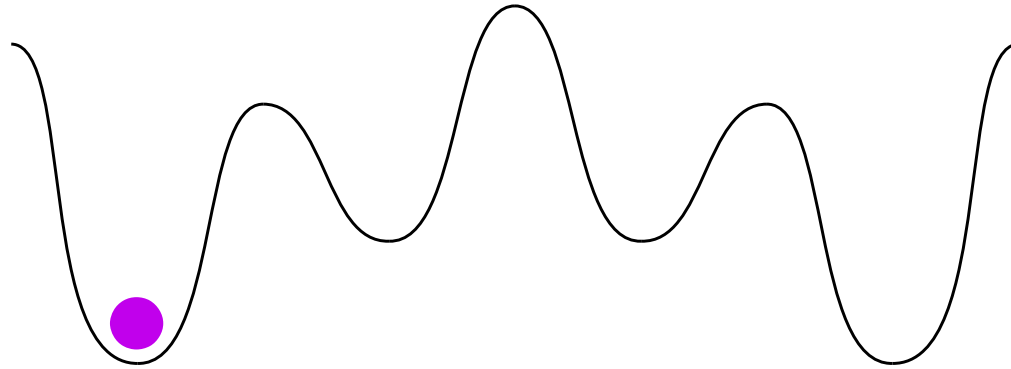
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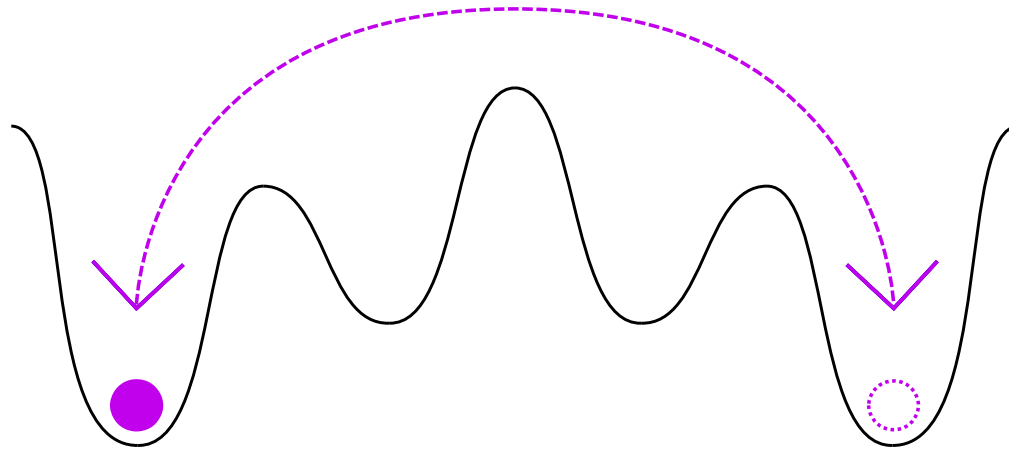
Hierarchical mirror structure (HMS)



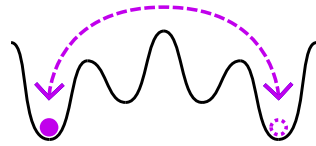
Hierarchical mirror structure (HMS)



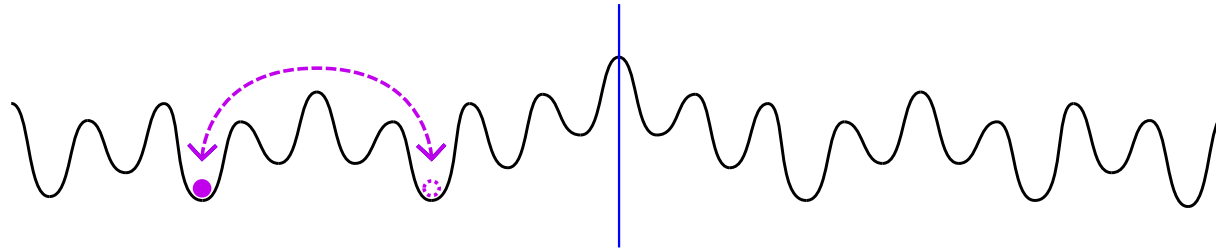
Hierarchical mirror structure (HMS)



Hierarchical mirror structure (HMS)

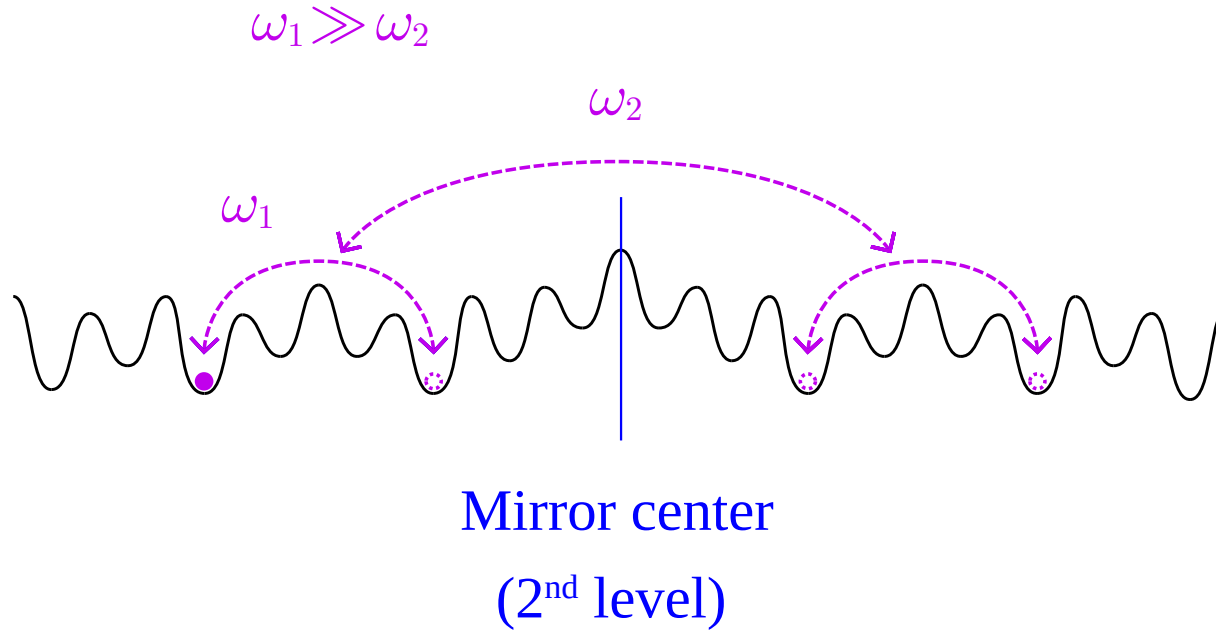


Hierarchical mirror structure (HMS)

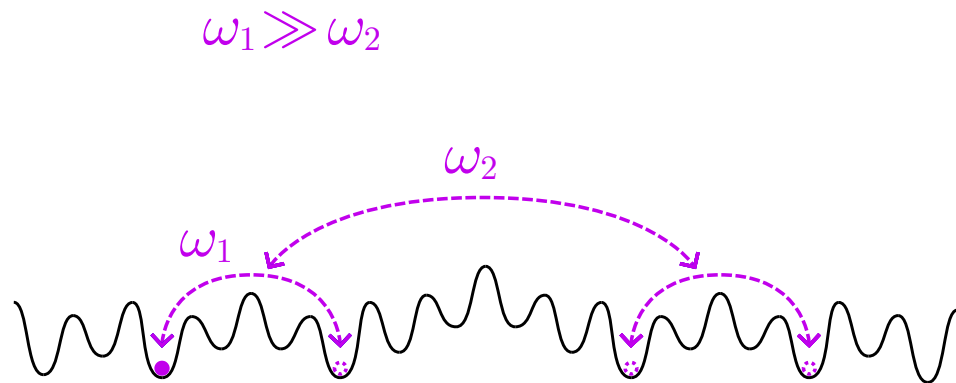


Mirror center
(2nd level)

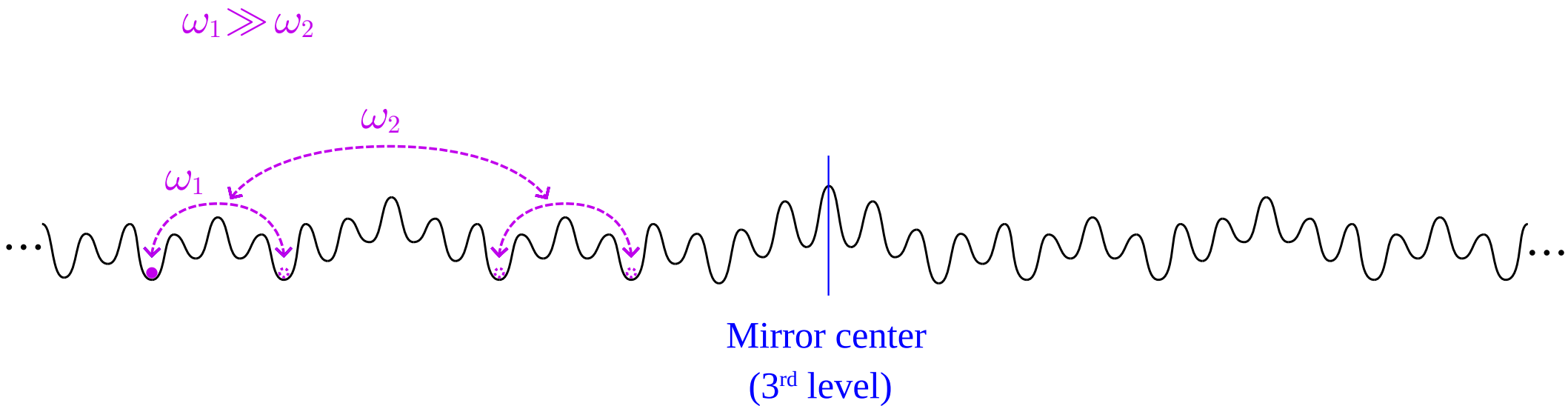
Hierarchical mirror structure (HMS)



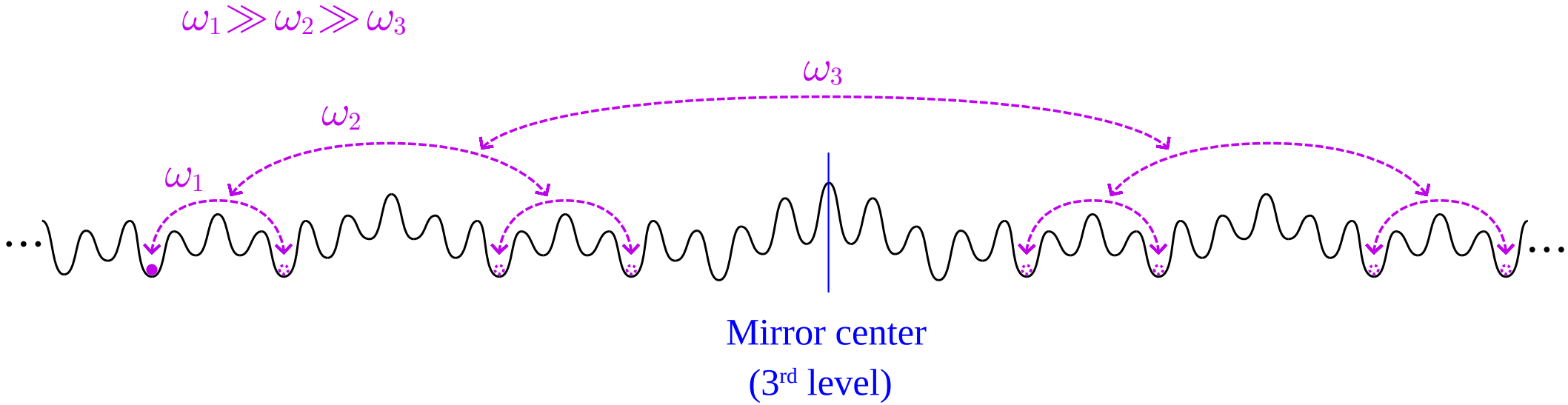
Hierarchical mirror structure (HMS)



Hierarchical mirror structure (HMS)



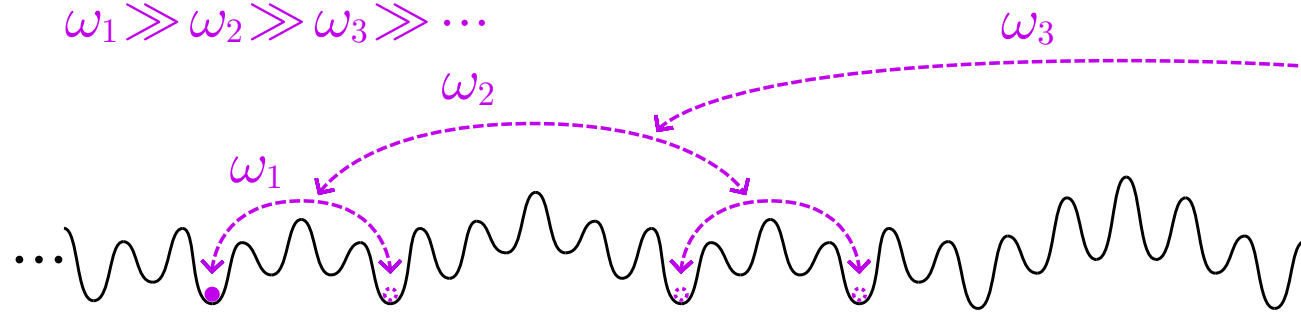
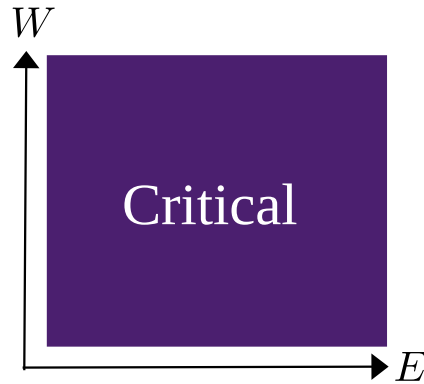
Hierarchical mirror structure (HMS)



Critical phases in many quasiperiodic systems has HMS

- Fibonacci quasicrystal [Jagannathan2021]
- Extended Aubry-Andre model [Avila2017]

Single
particle



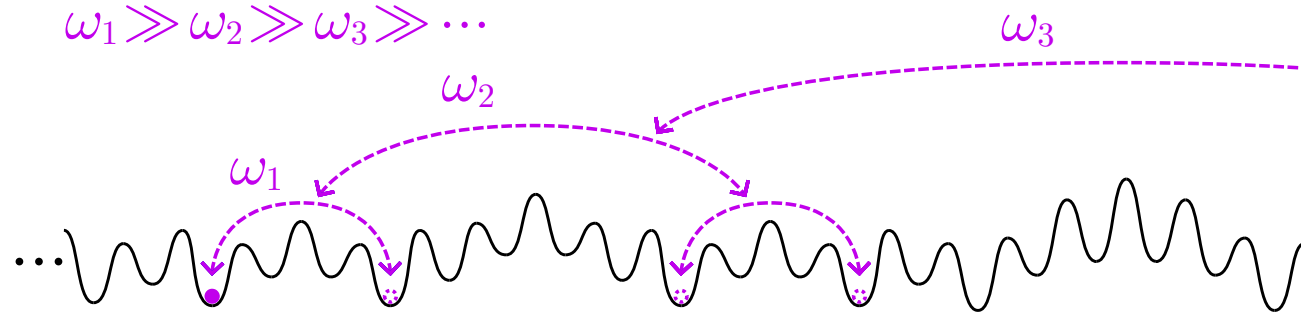
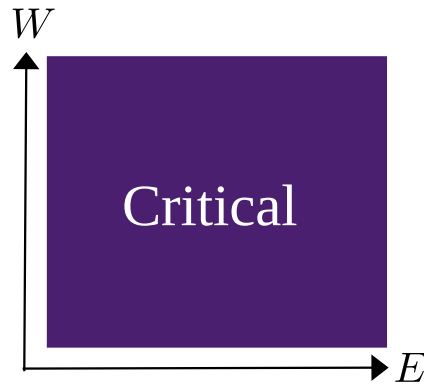
The particle goes arbitrarily far away,
but also comes back after arbitrarily long time

Interacting
Many-body

Critical phases in many quasiperiodic systems has HMS

- Fibonacci quasicrystal [Jagannathan2021]
- Extended Aubry-Andre model [Avila2017]

Single
particle



The particle goes arbitrarily far away,
but also comes back after arbitrarily long time

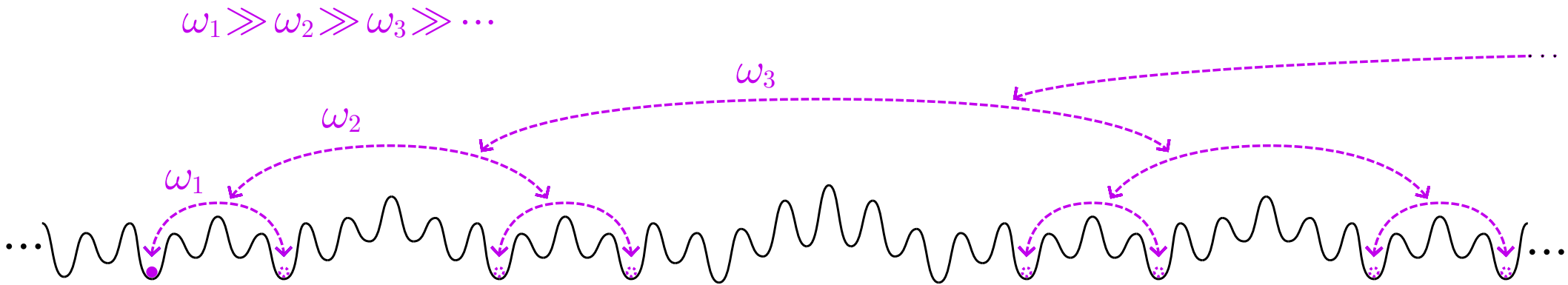
Interacting
Many-body

Becomes MBL?
Becomes ETH?
Many-body critical phase?

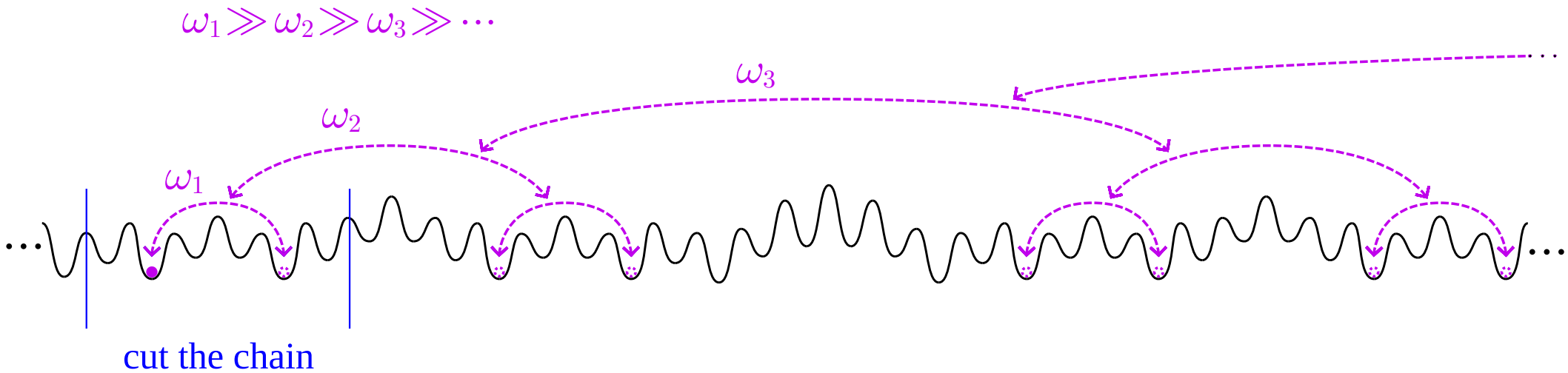
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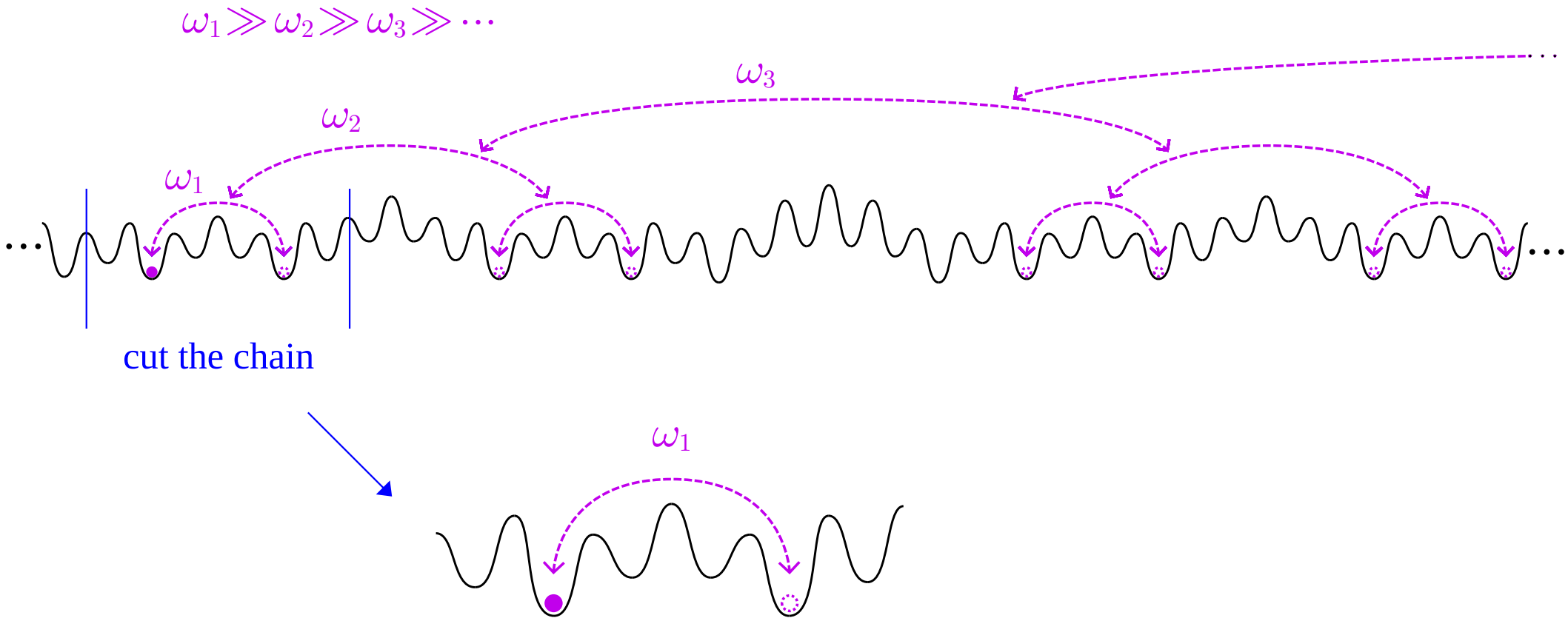
Restriction to the first level



Restriction to the first level



Restriction to the first level

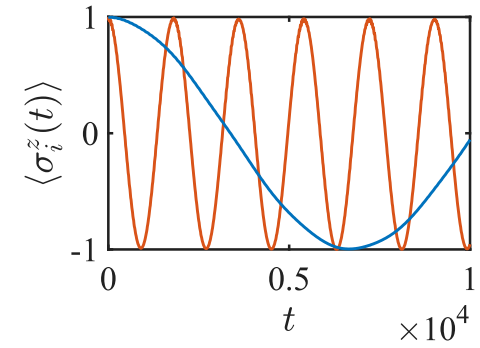
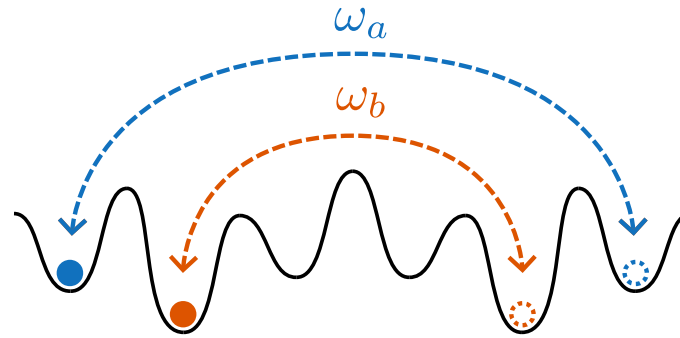


Resonance (oscillation) in the first level

Effects of interaction in the first level

Li, YTT, Das Sarma PRB (2026)

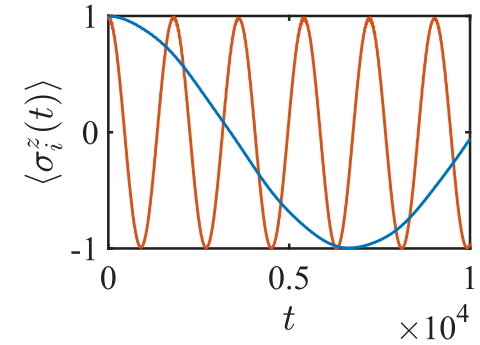
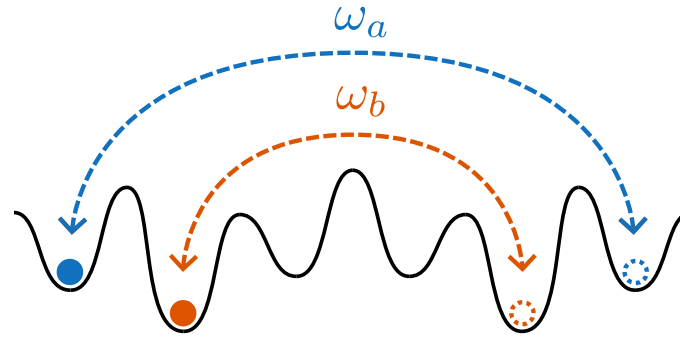
Non-interacting:



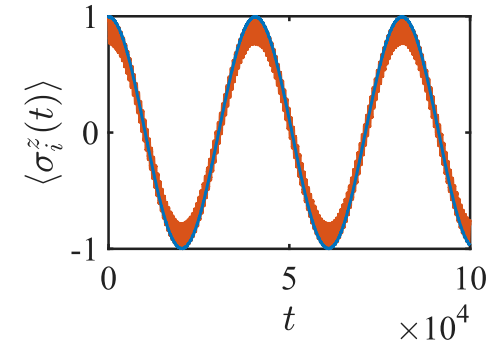
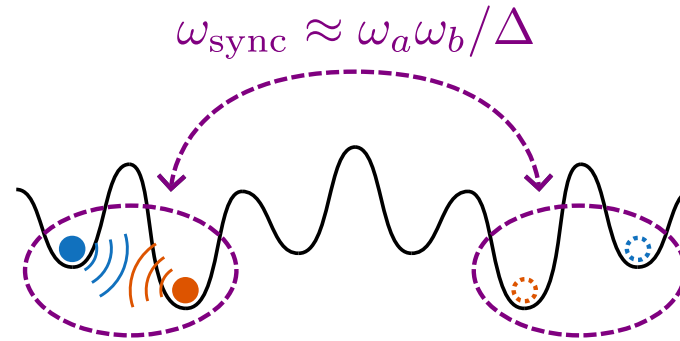
Effects of interaction in the first level

Li, YTT, Das Sarma PRB (2026)

Non-interacting:



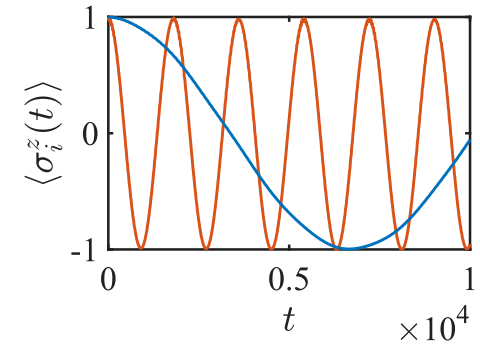
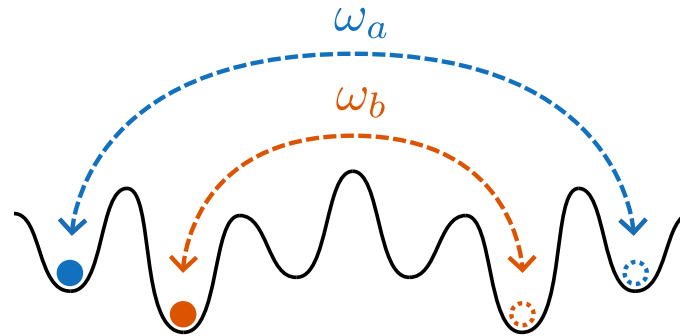
Interacting:



Effects of interaction in the first level

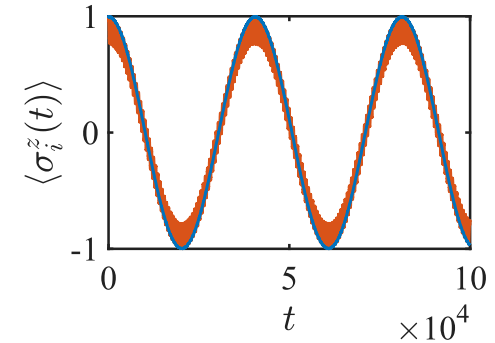
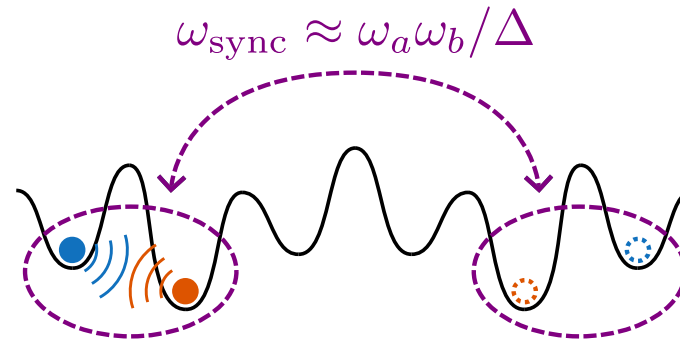
Li, YTT, Das Sarma PRB (2026)

Non-interacting:



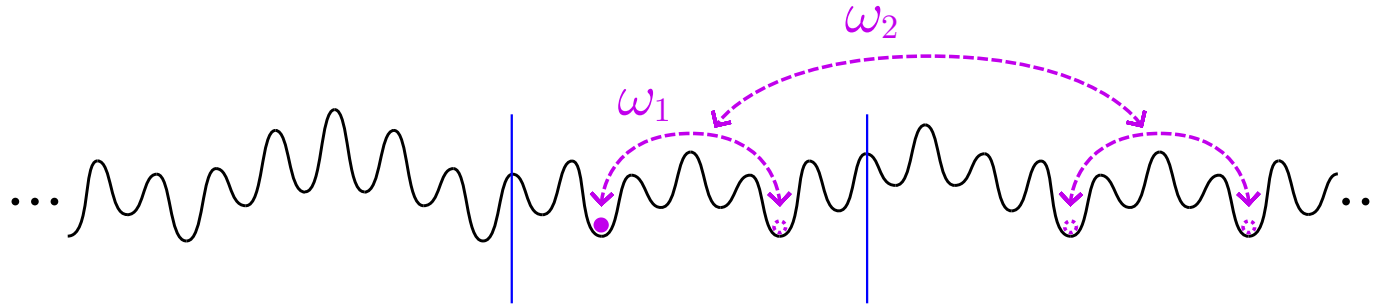
Unsynchronized

Interacting:

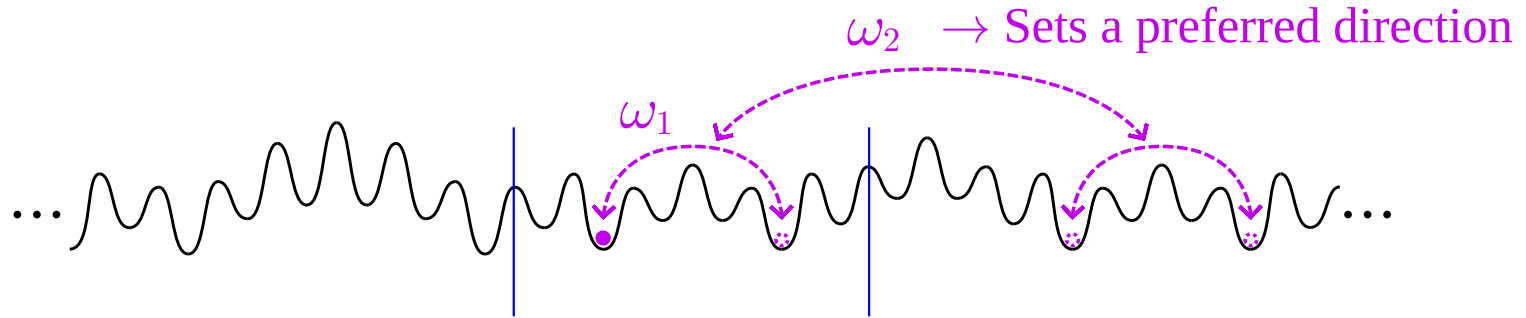


Synchronized

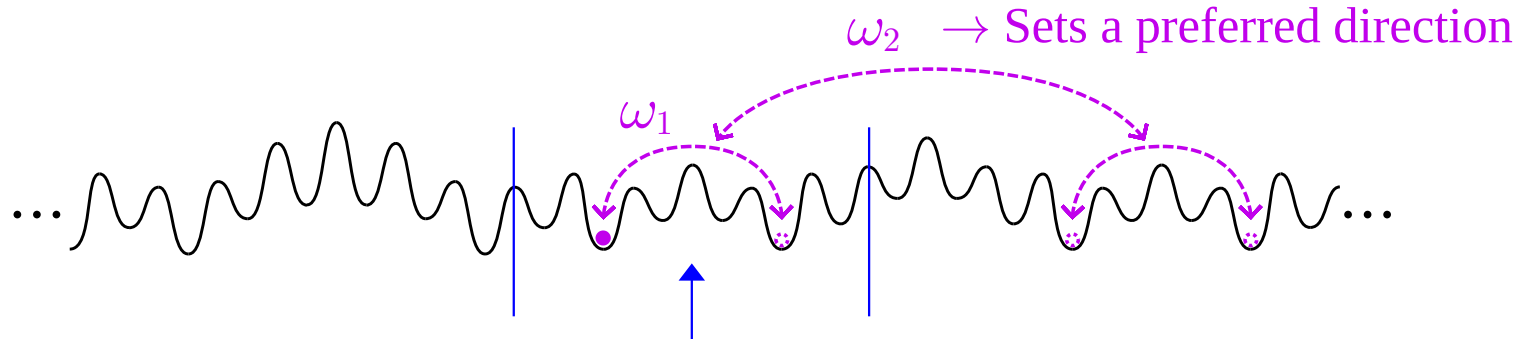
Symmetry breaking



Symmetry breaking

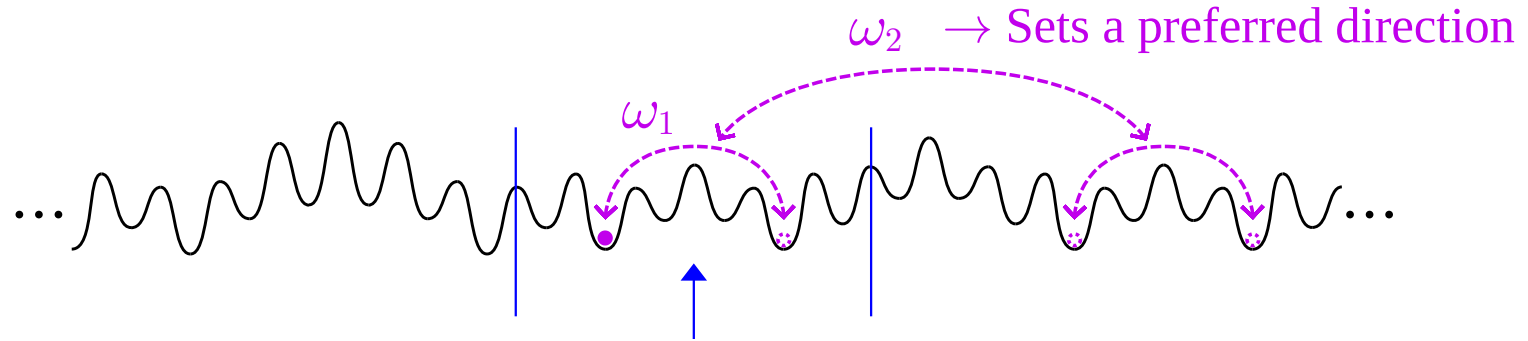


Symmetry breaking

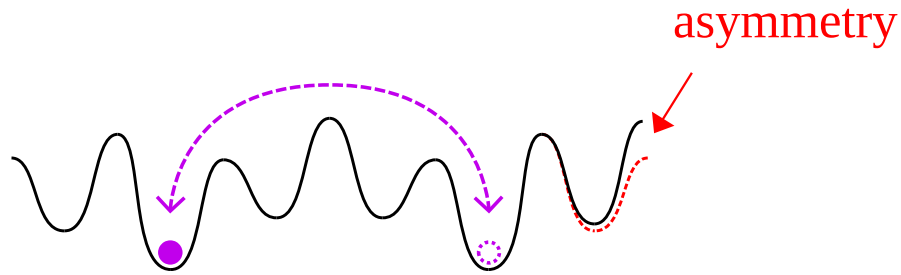


This mirror symmetry cannot be exact

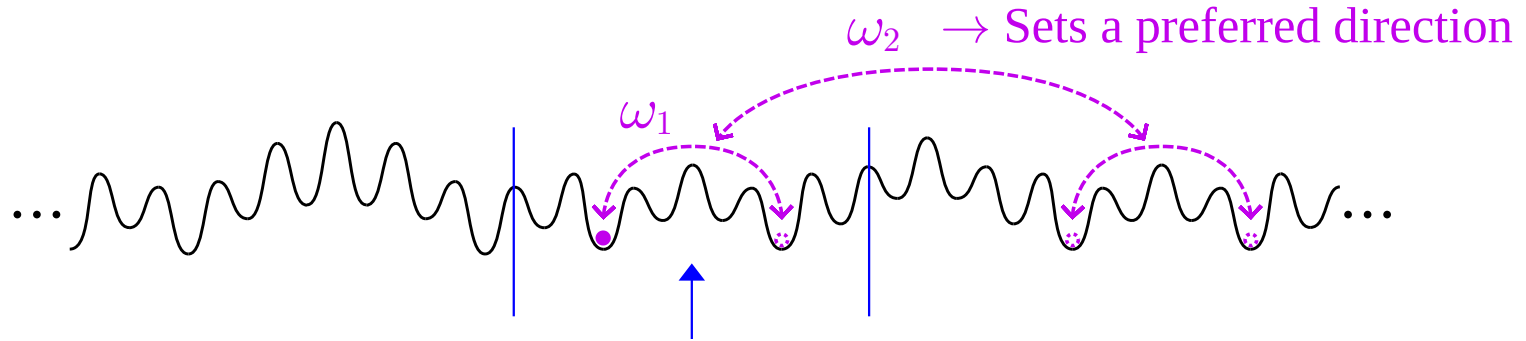
Symmetry breaking



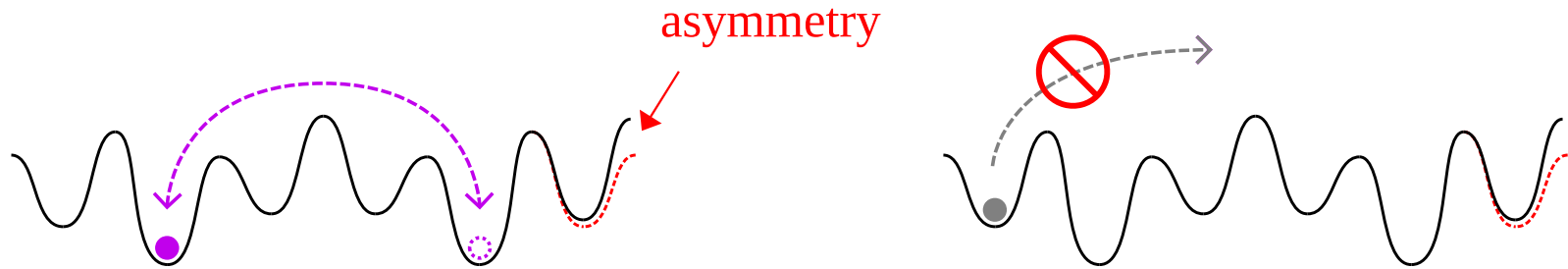
This mirror symmetry cannot be exact



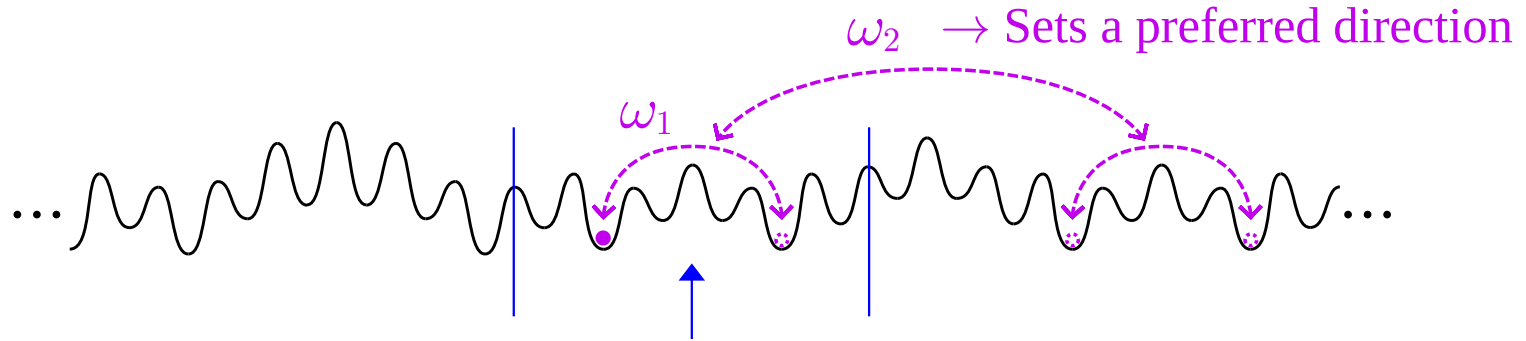
Symmetry breaking



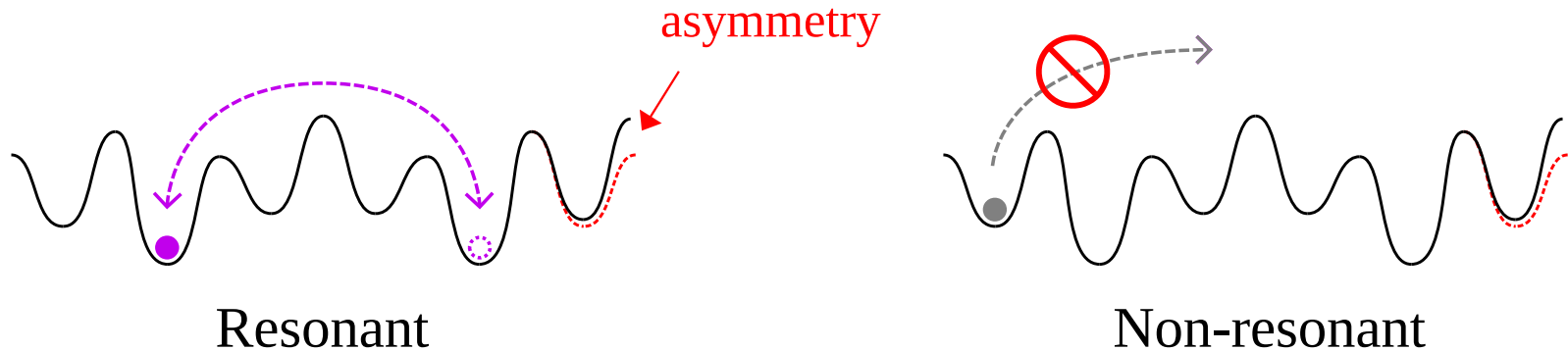
This mirror symmetry cannot be exact



Symmetry breaking

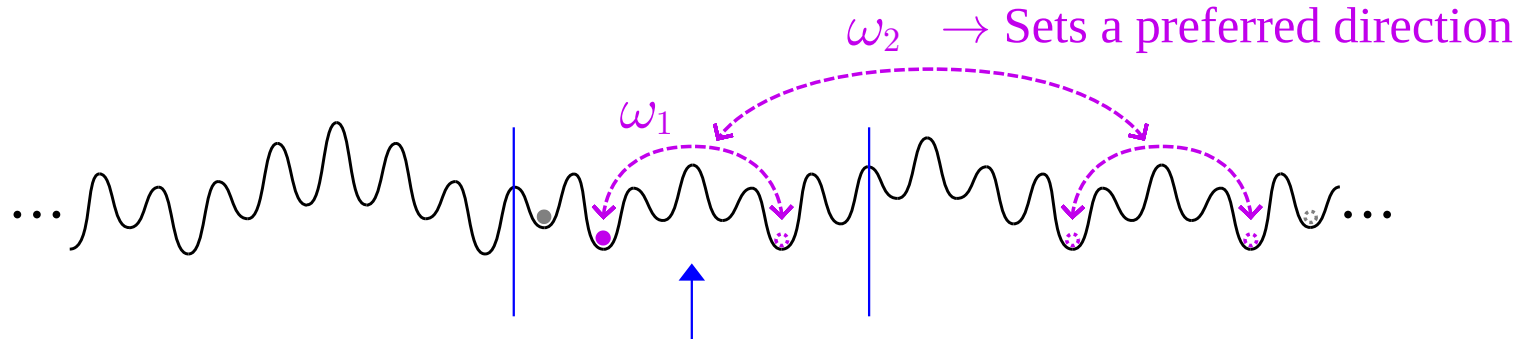


This mirror symmetry cannot be exact

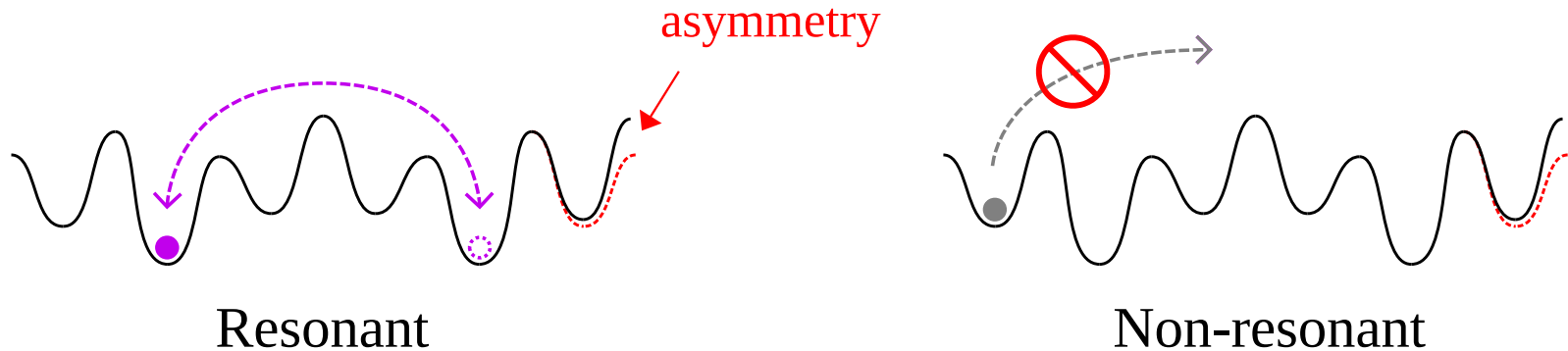


A particle too far away from the mirror center becomes *non-resonant*

Symmetry breaking

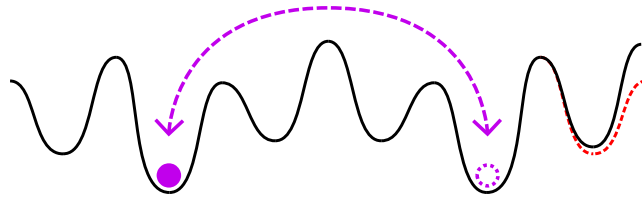


This mirror symmetry cannot be exact

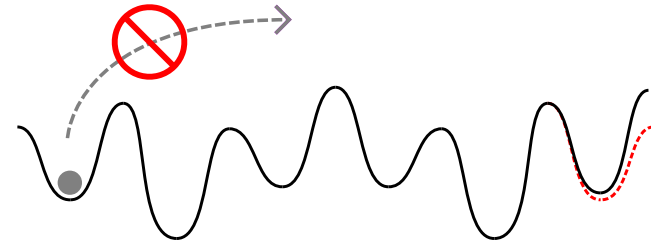


A particle too far away from the mirror center becomes *non-resonant*

Interplay between interaction and symmetry breaking

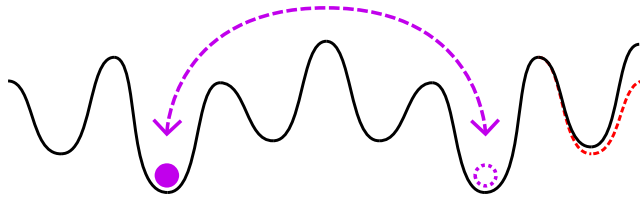


Resonant

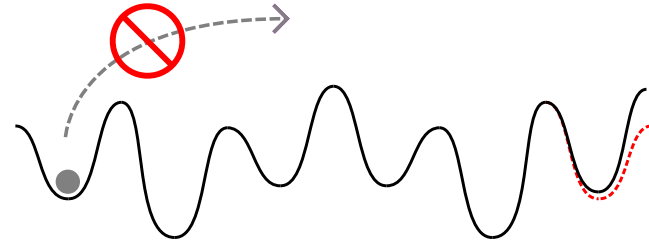


Non-resonant

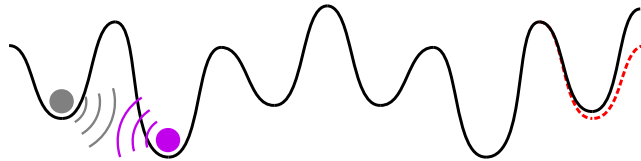
Interplay between interaction and symmetry breaking



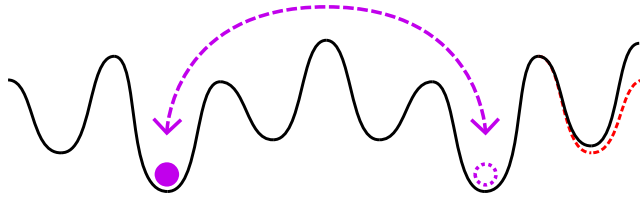
Resonant



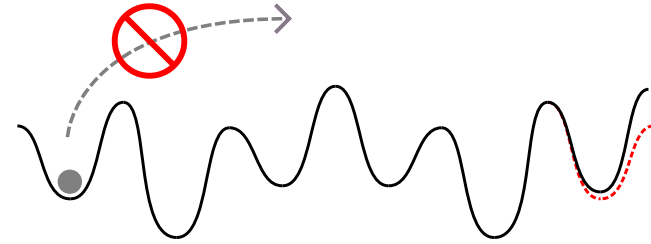
Non-resonant



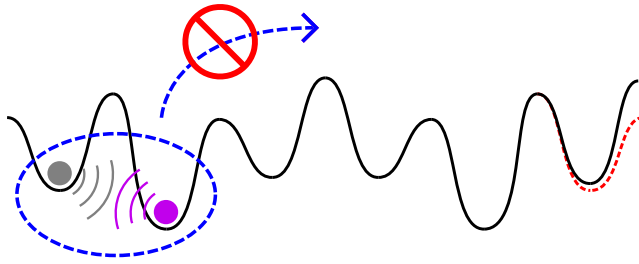
Interplay between interaction and symmetry breaking



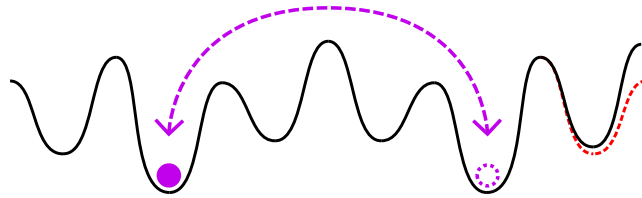
Resonant



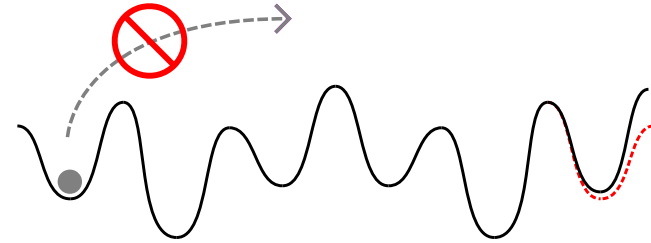
Non-resonant



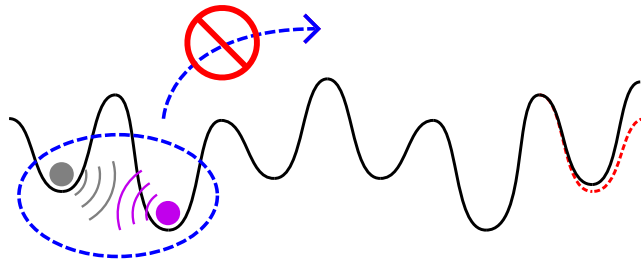
Interplay between interaction and symmetry breaking



Resonant

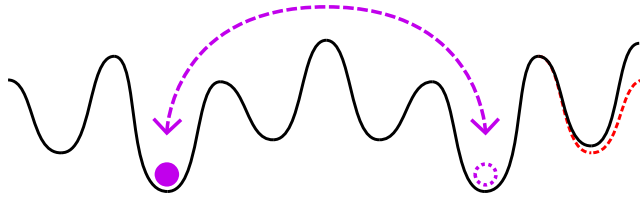


Non-resonant

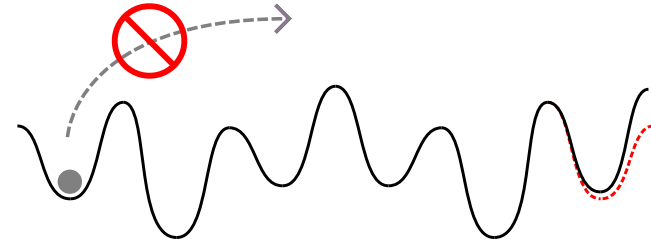


Frozen

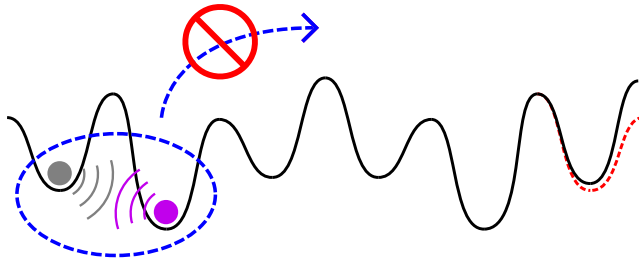
Interplay between interaction and symmetry breaking



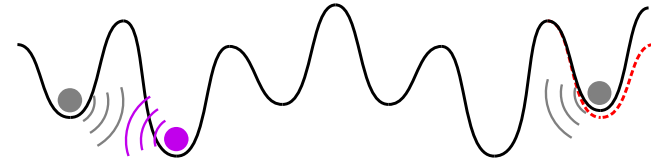
Resonant



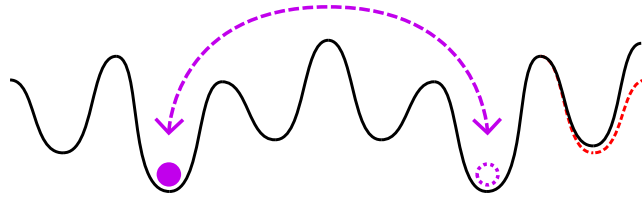
Non-resonant



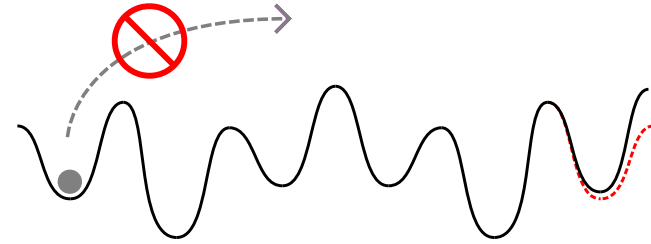
Frozen



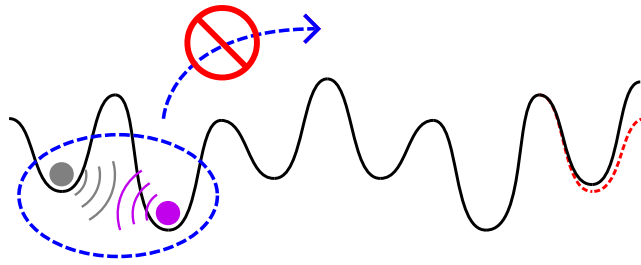
Interplay between interaction and symmetry breaking



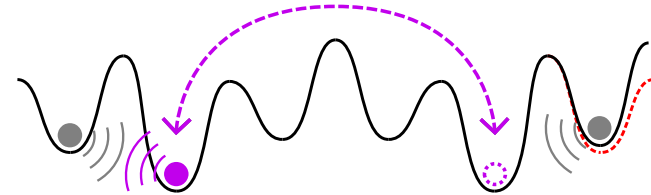
Resonant



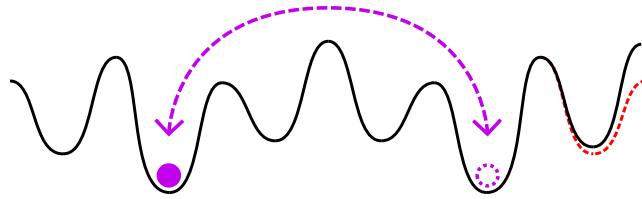
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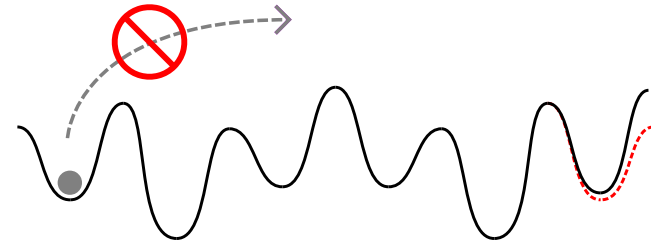
Frozen



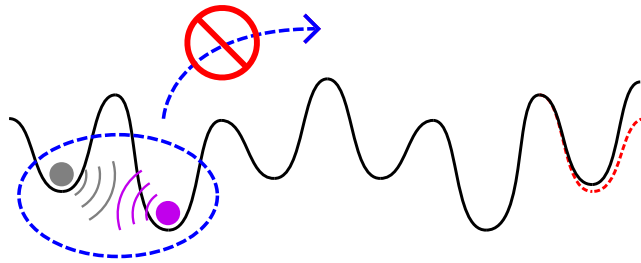
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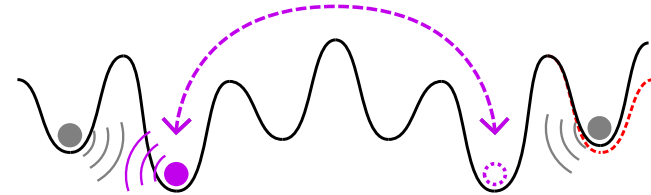
Resonant



Non-resonant

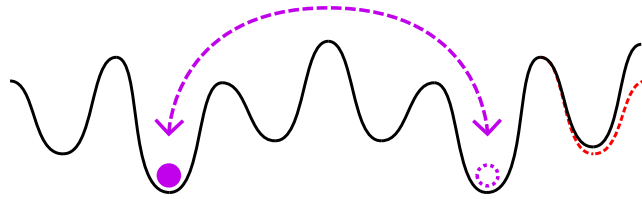


Frozen

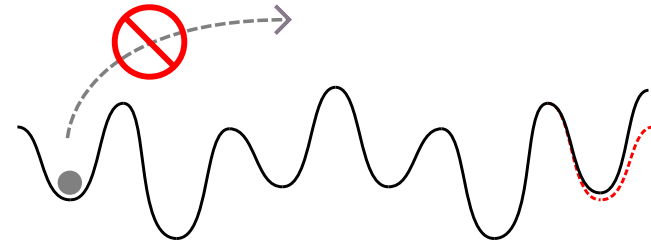


Protected

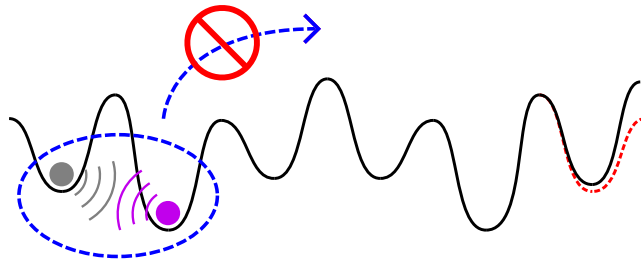
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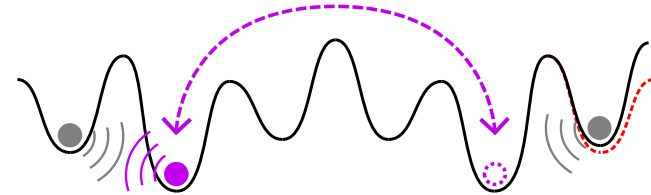
Resonant



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Frozen

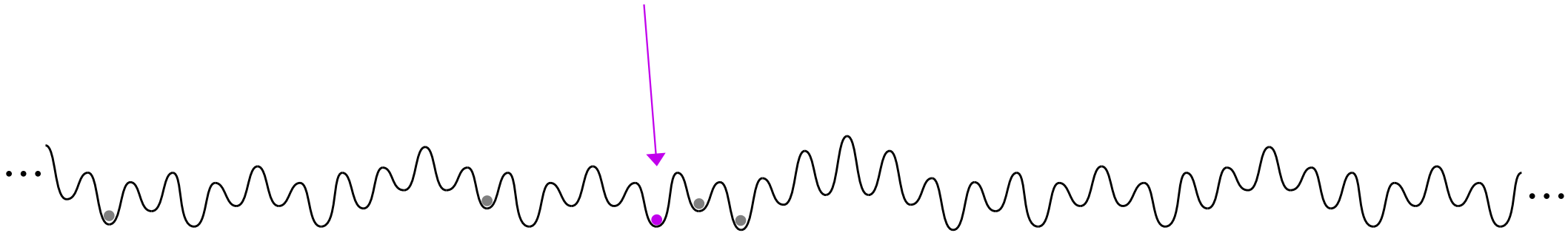


Protected

Resonances are *protected* iff the configs of other particles are symmetric enough

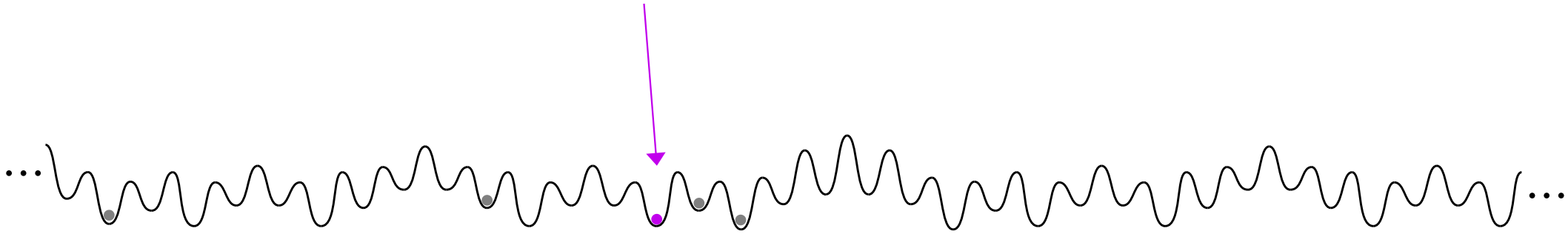
Condition of delocalization in HMS with interactions

Given an initial many-body configuration, can the purple particle go arbitrarily far away?



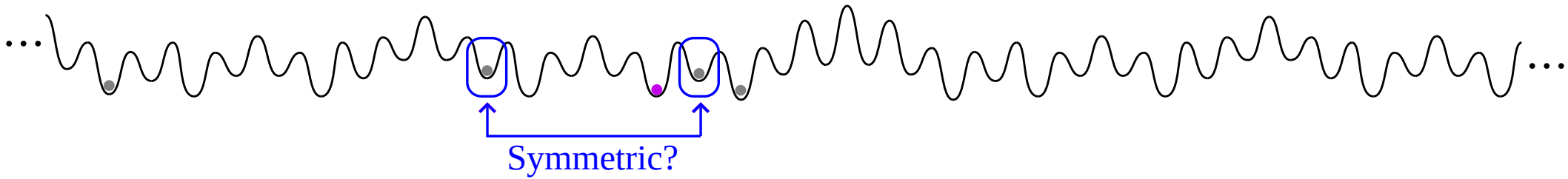
Condition of delocalization in HMS with interactions

Given an initial many-body configuration, can the purple particle go arbitrarily far away?
i.e. Is it in a localized phase or not?



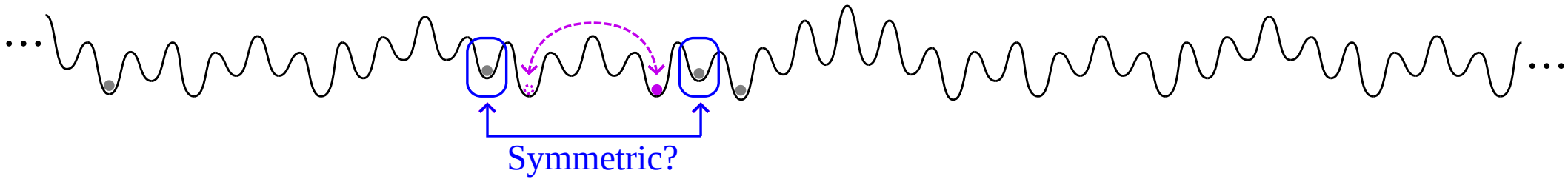
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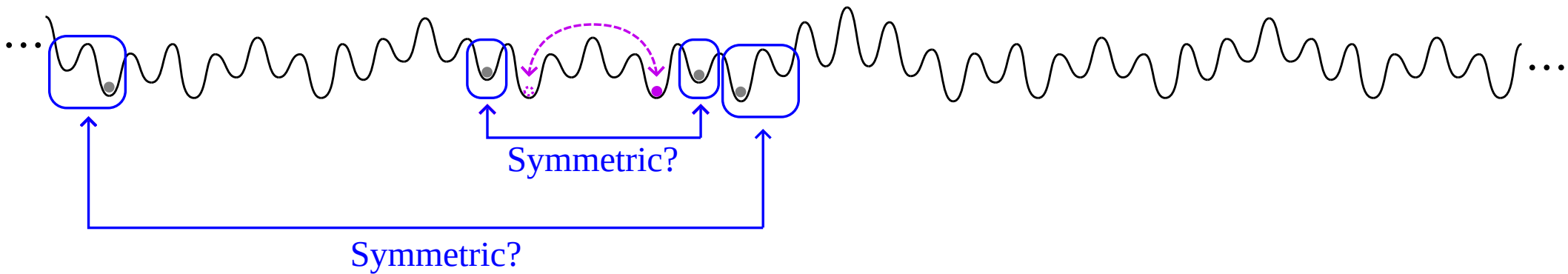
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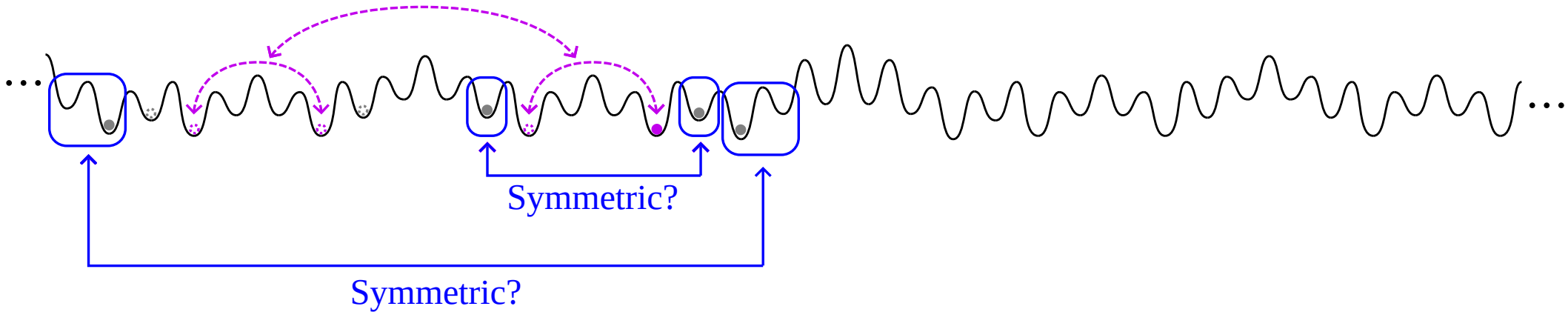
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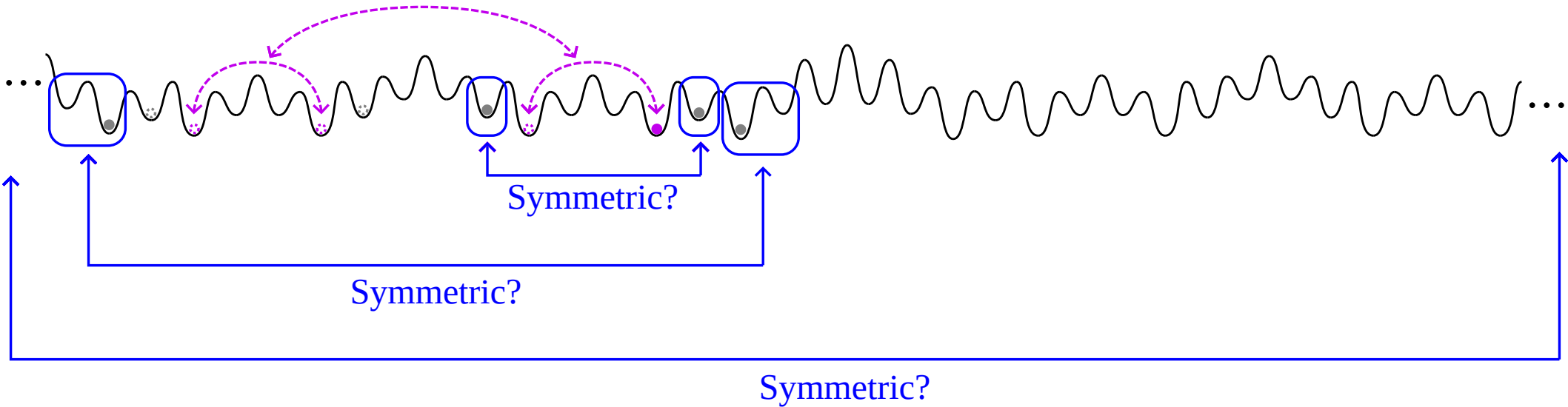
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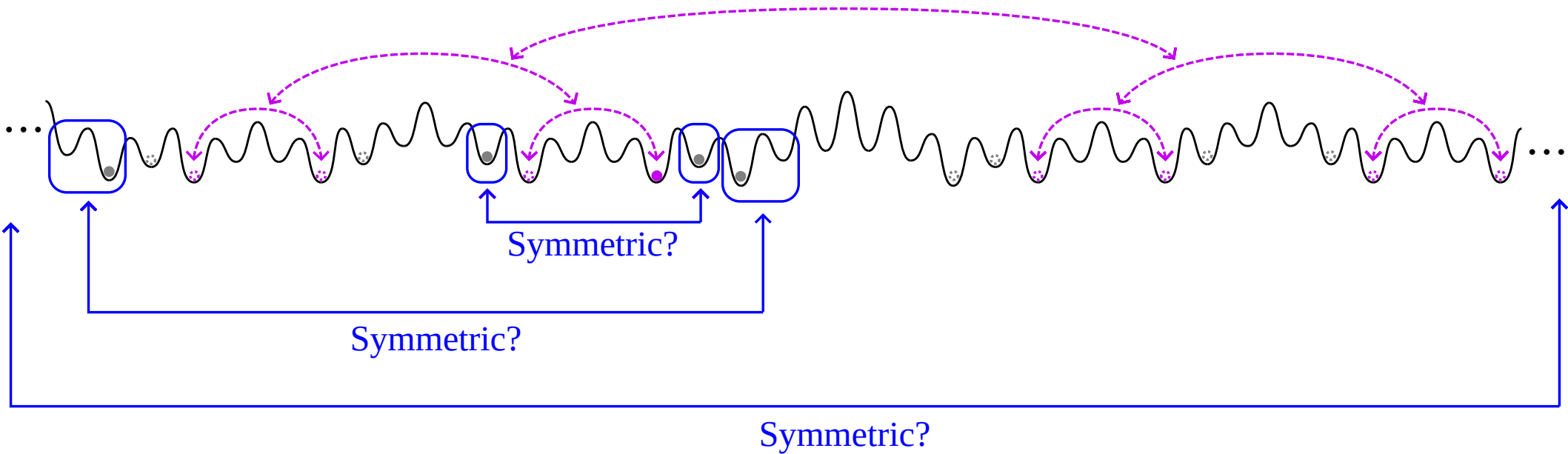
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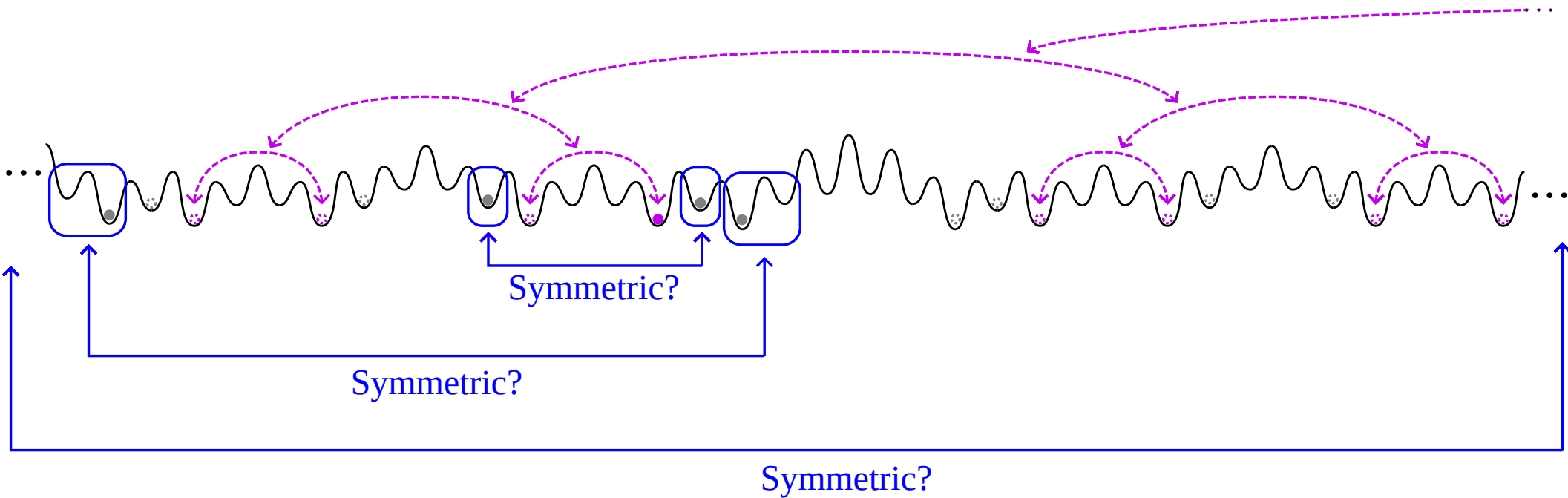
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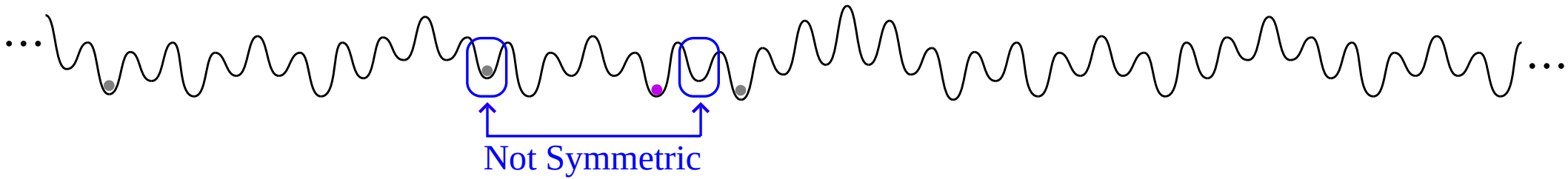
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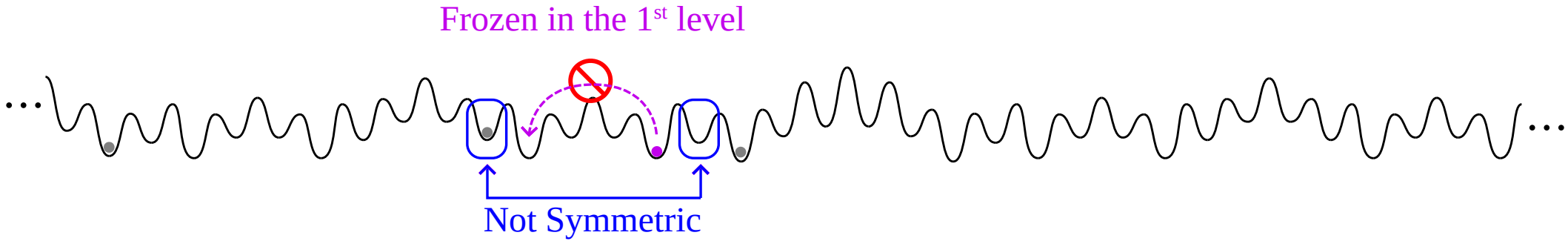
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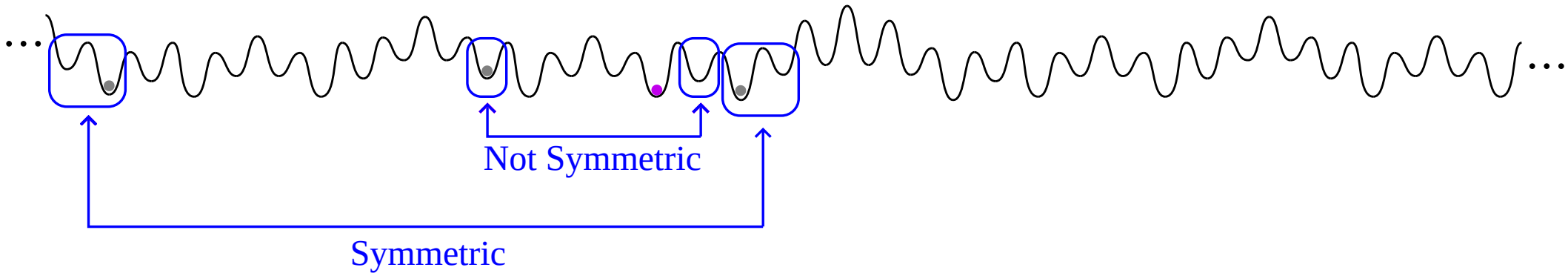
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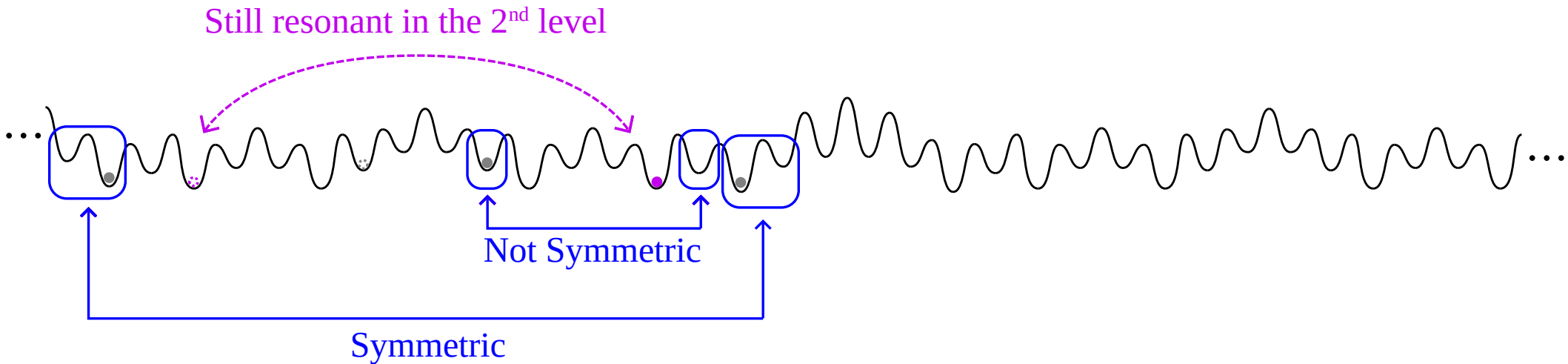
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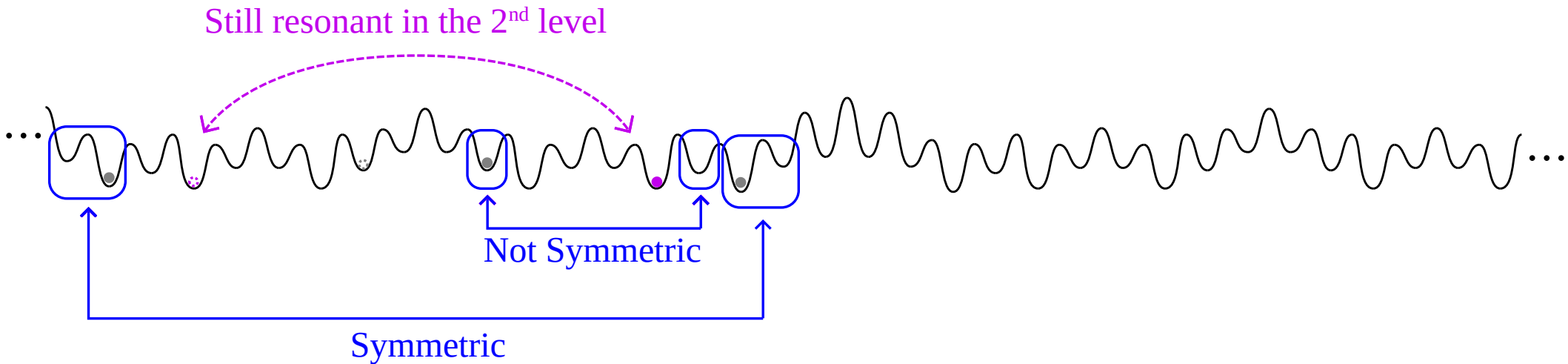
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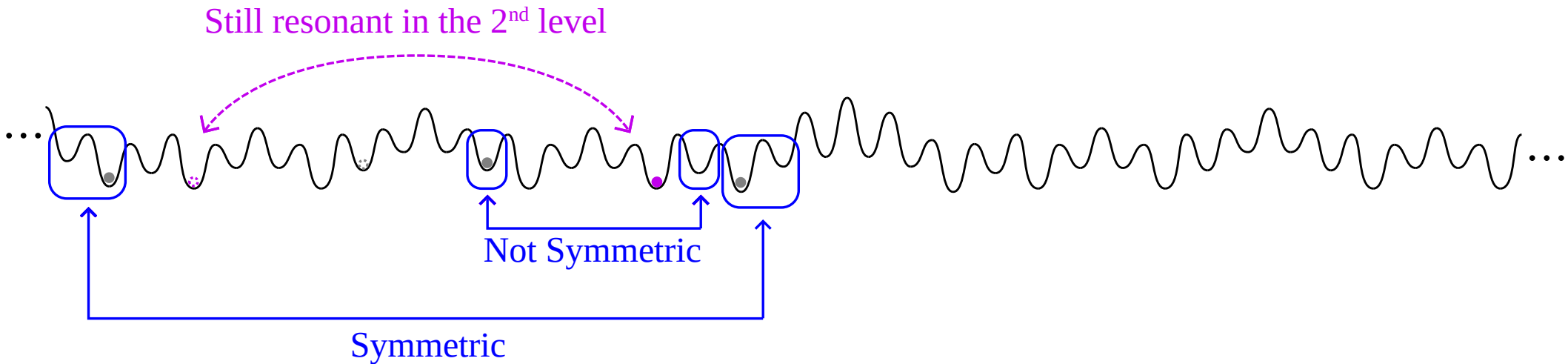
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As long as there are resonances at level n for arbitrarily large n ,

Condition of delocalization in HMS with interactions

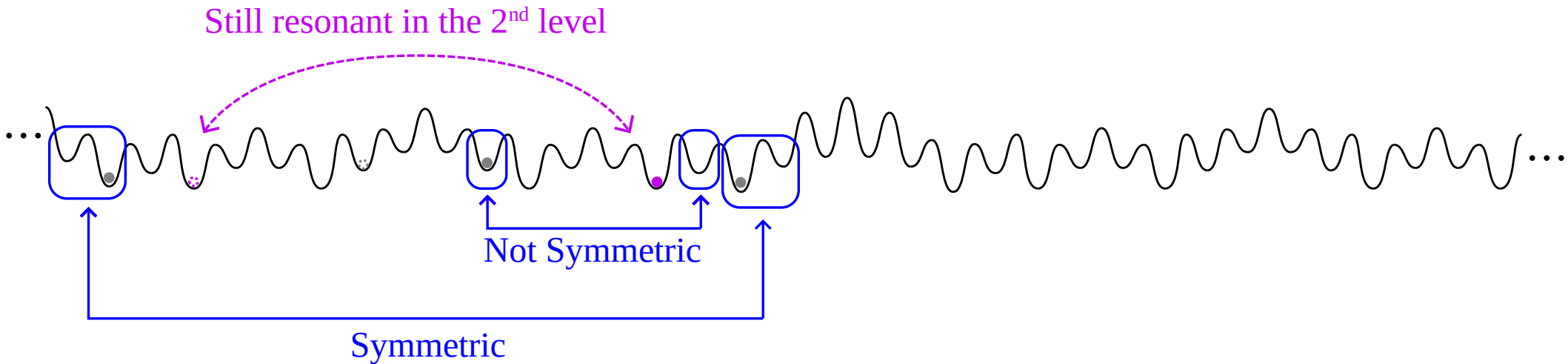
Given an initial many-body configuration, can the purple particle go arbitrarily far away?
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As long as there are resonances at level n for arbitrarily large n , the particle will go arbitrarily far away.

Condition of delocalization in HMS with interactions

Given an initial many-body configuration, can the purple particle go arbitrarily far away?
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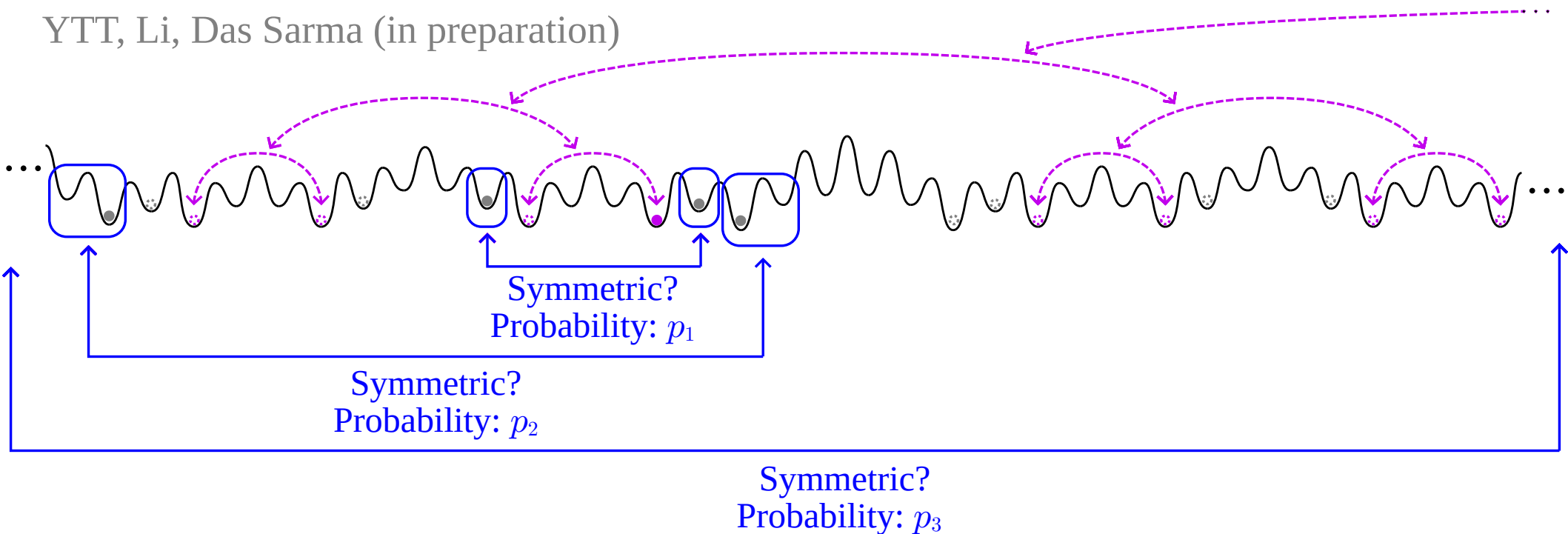


As long as there are resonances at level n for arbitrarily large n ,
the particle will go arbitrarily far away.

(No need to resonate at every n)

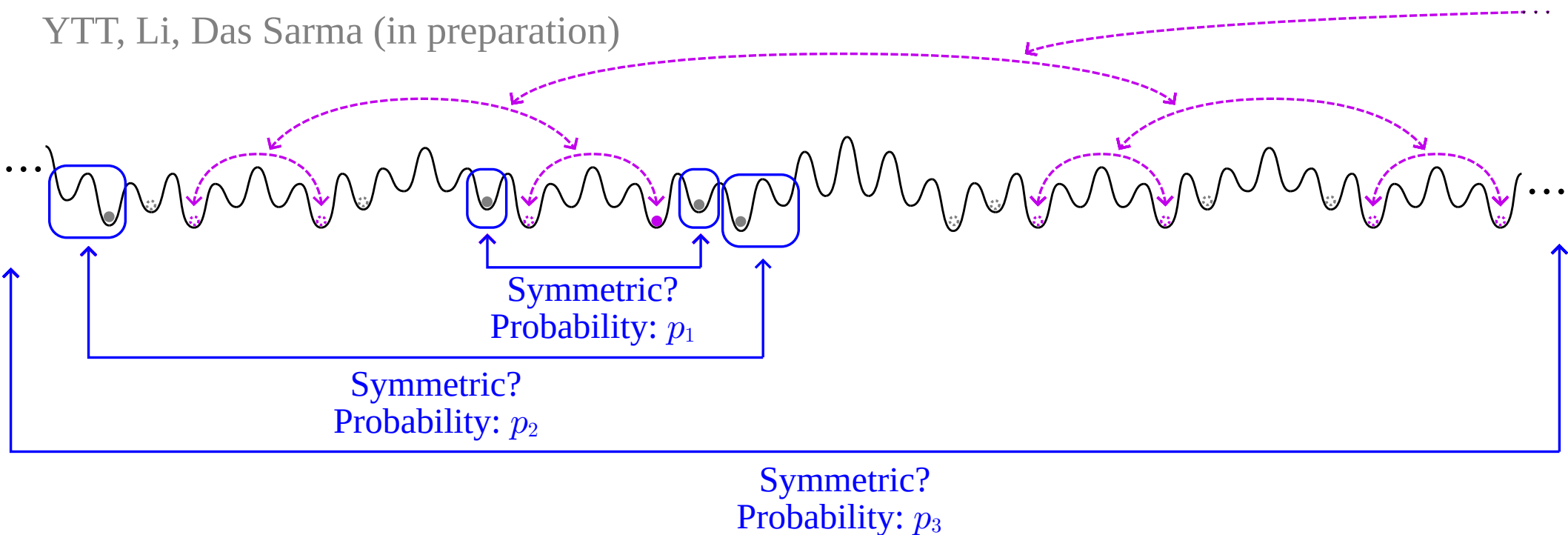
Condition of delocalization in HMS with interactions

YTT, Li, Das Sarma (in preparation)



Condition of delocalization in HMS with interactions

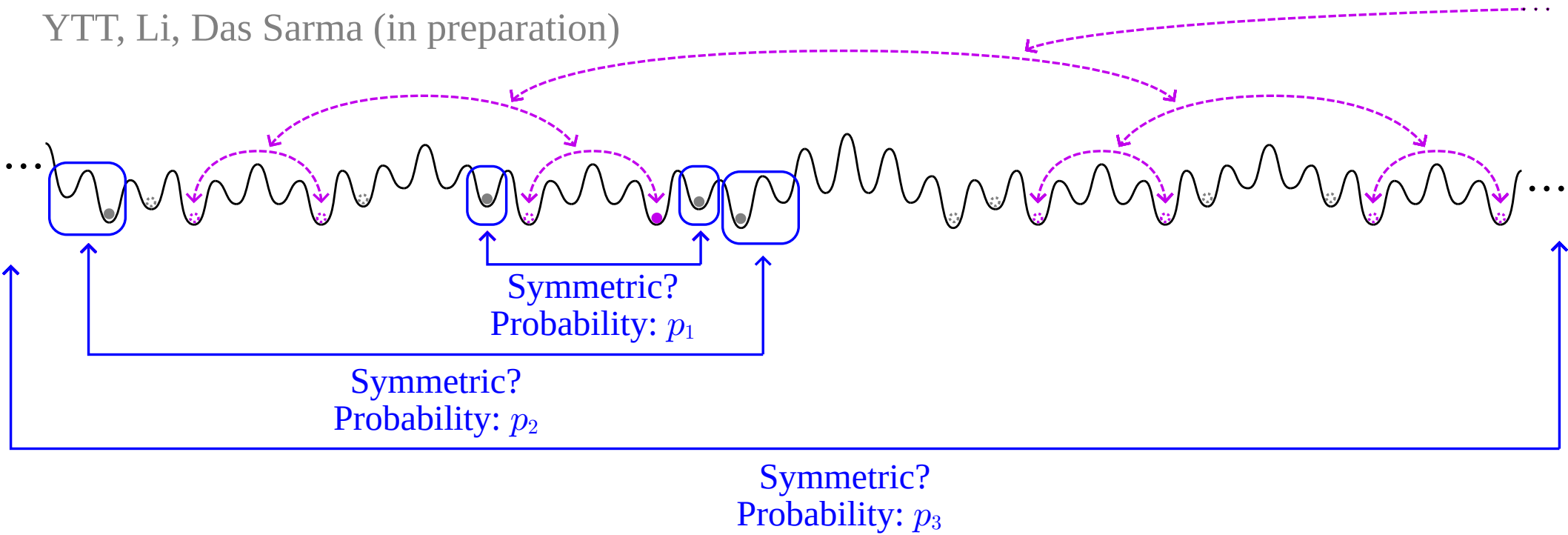
YTT, Li, Das Sarma (in preparation)



Expected number of resonances (0 or 1) at level n : p_n

Condition of delocalization in HMS with interactions

YTT, Li, Das Sarma (in preparation)

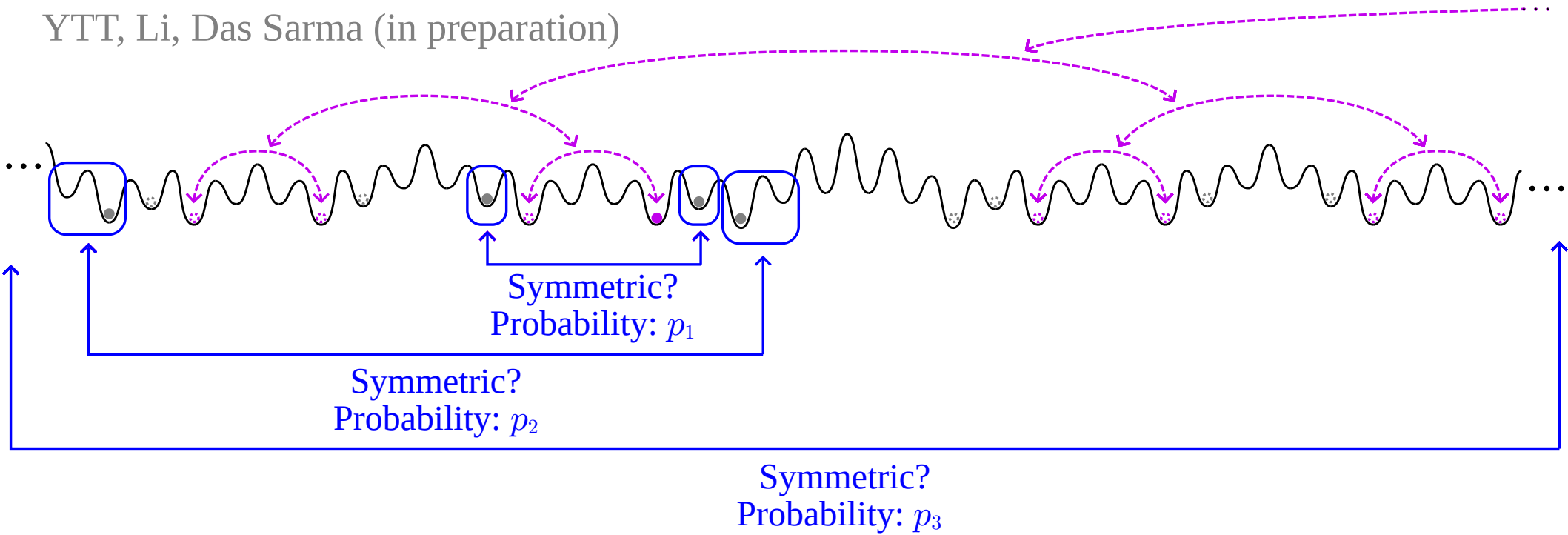


Expected number of resonances (0 or 1) at level n : p_n

Expected number of resonances from level 1 to n : $p_1 + p_2 + \dots + p_n$

Condition of delocalization in HMS with interactions

YTT, Li, Das Sarma (in preparation)



Expected number of resonances (0 or 1) at level n : p_n

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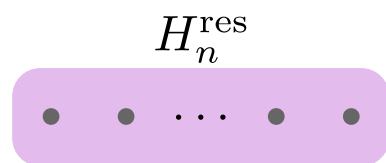
Taking $n \rightarrow \infty$, we get:

The particle goes arbitrarily far away iff $\sum p_n$ diverges.

Outline

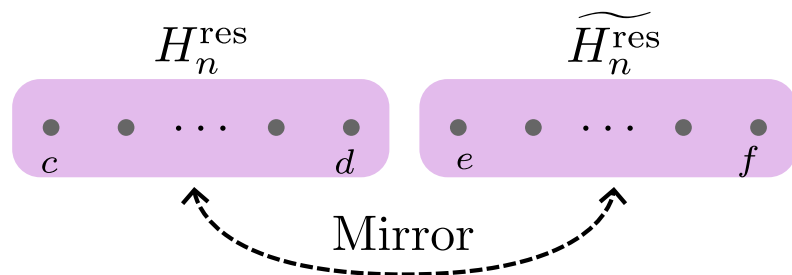
- Introduction
 - Many-body localization (MBL)
 - Random v.s. Quasiperiodic systems
- Non-ergodic extended (NEE) regime
 - Numerical results
 - Theoretical explanation
- Many-body critical (MBC) phases
 - Single-particle hierarchical mirror structure (HMS)
 - Effects of interaction on HMS
 - Asymptotically solvable model

Construction of the solvable model of MBC



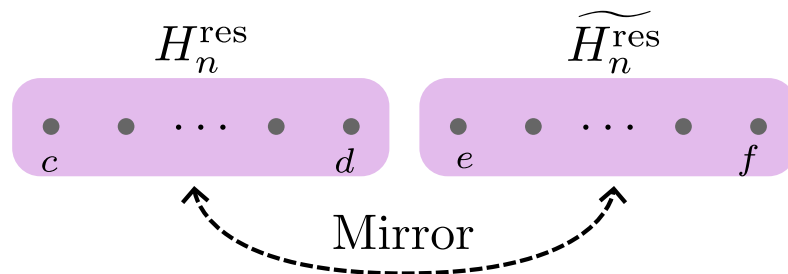
Random mixed-field Ising chain

Construction of the solvable model of MBC

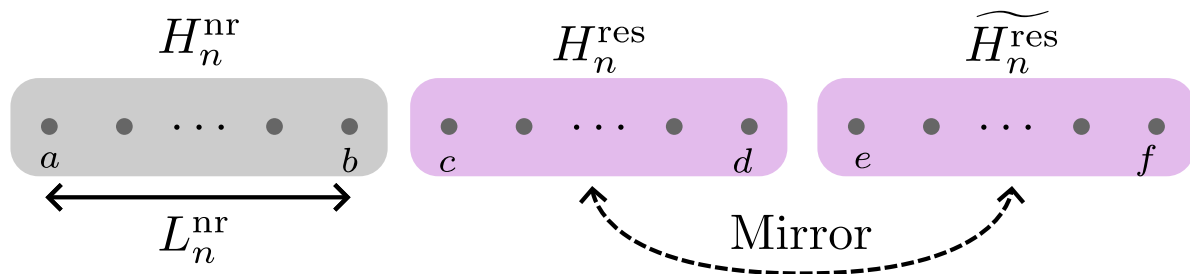


Construction of the solvable model of MBC

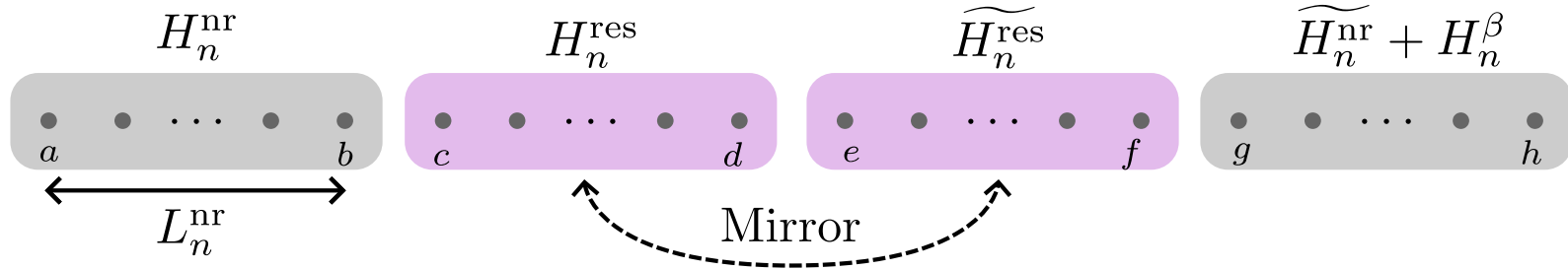
$\widetilde{X}$: mirror reflection of X



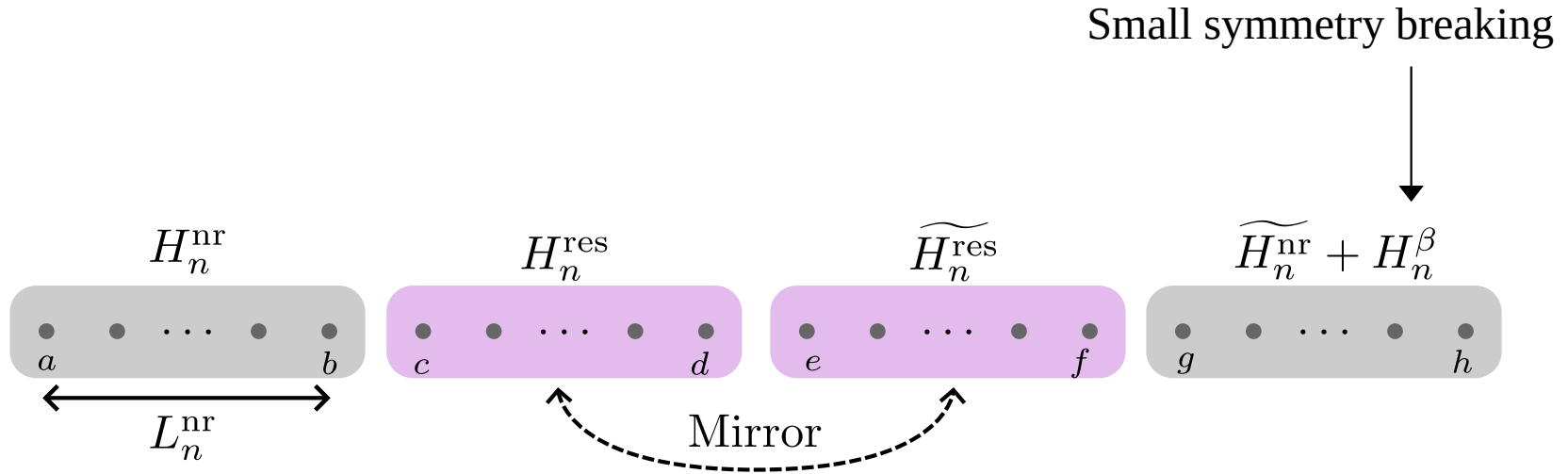
Construction of the solvable model of MBC



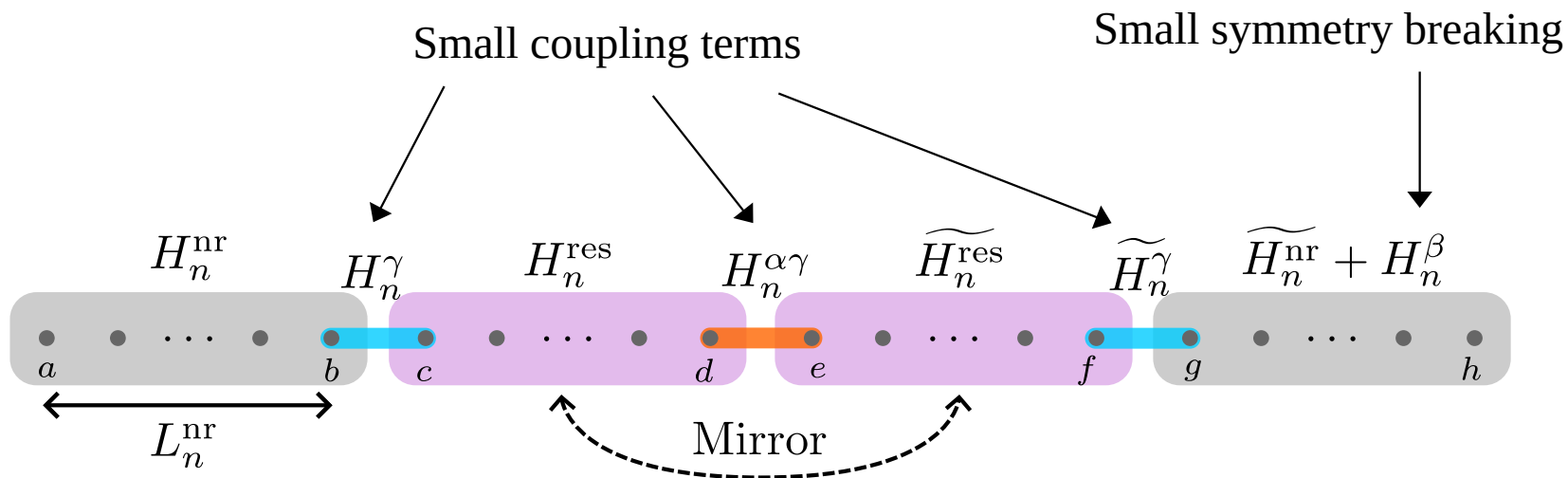
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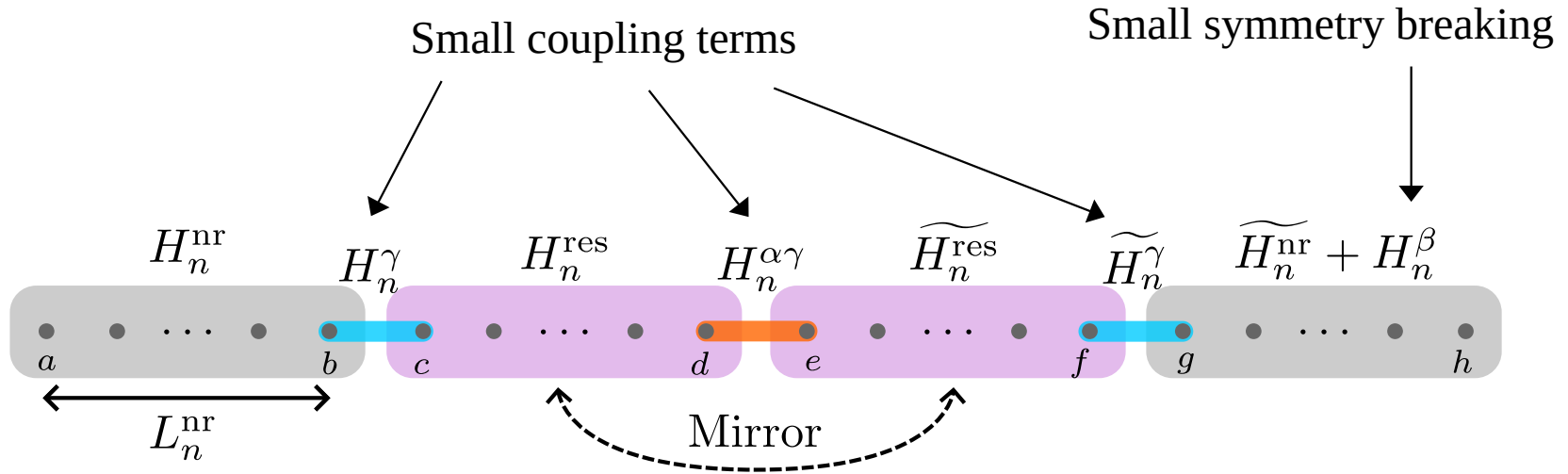
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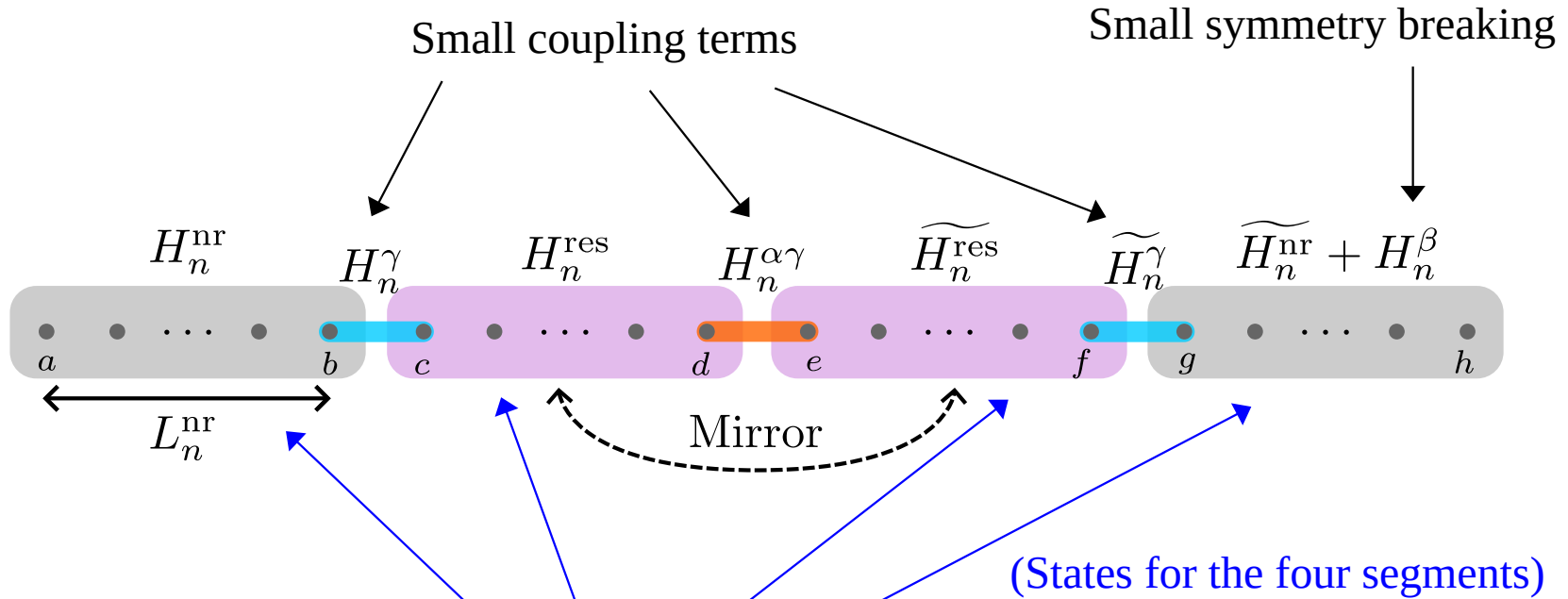


Construction of the solvable model of MBC



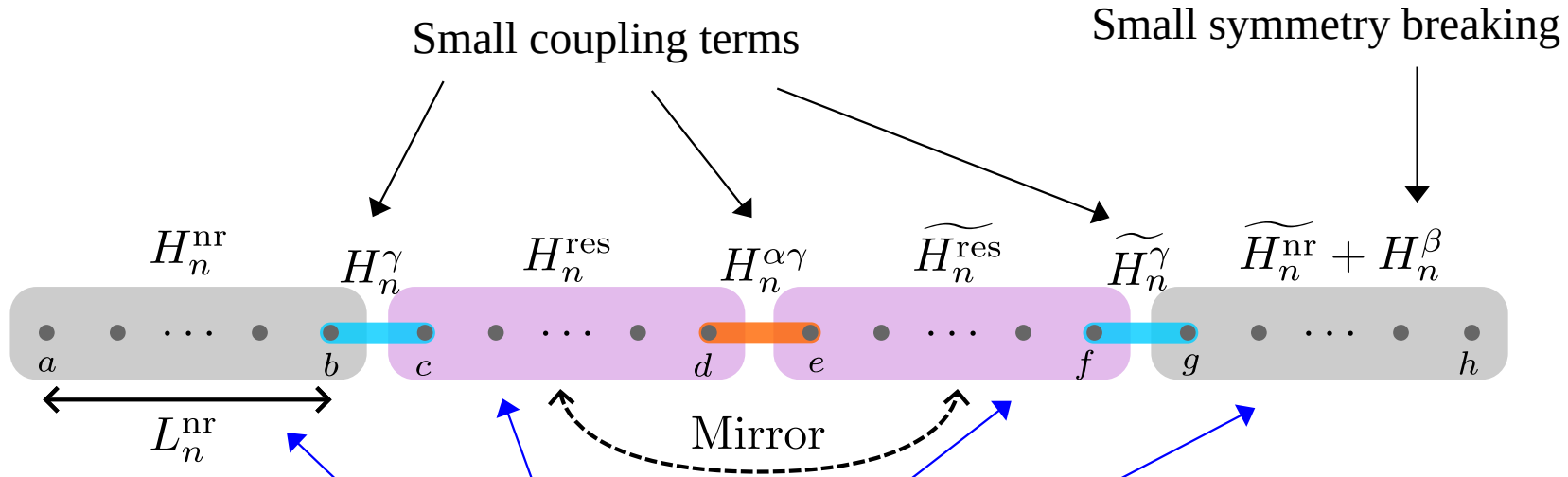
$$\text{perturbed eigenstate} \approx \begin{cases} |ABCD\rangle_n, & \text{if } A \neq \widetilde{D} \text{ or } B = C, \\ \frac{1}{\sqrt{2}} (|ABCD\rangle_n + |A\widetilde{C}\widetilde{B}D\rangle_n), & \text{if } A = \widetilde{D}, B > C, \\ \frac{1}{\sqrt{2}} (|ABCD\rangle_n - |A\widetilde{C}\widetilde{B}D\rangle_n), & \text{if } A = \widetilde{D}, B < C. \end{cases}$$

Construction of the solvable model of MBC



perturbed eigenstate $\approx$
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Construction of the solvable model of MBC

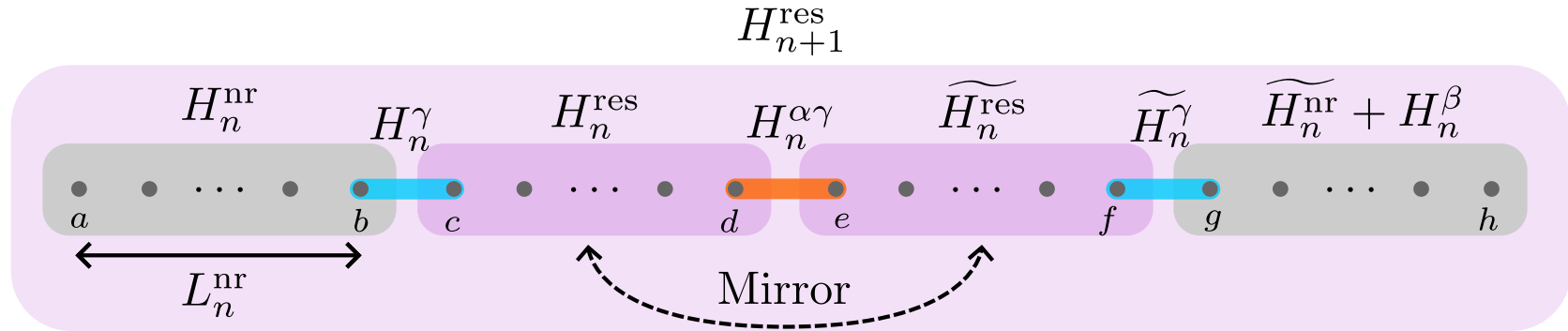


(States for the four segments)

$\widetilde{X}$: mirror reflection of X

$$\text{perturbed eigenstate} \approx \begin{cases} |ABCD\rangle_n, & \text{if } A \neq \widetilde{D} \text{ or } B = C, \\ \frac{1}{\sqrt{2}} (|ABCD\rangle_n + |A\widetilde{C}\widetilde{B}D\rangle_n), & \text{if } A = \widetilde{D}, B > C, \\ \frac{1}{\sqrt{2}} (|ABCD\rangle_n - |A\widetilde{C}\widetilde{B}D\rangle_n), & \text{if } A = \widetilde{D}, B < C. \end{cases}$$

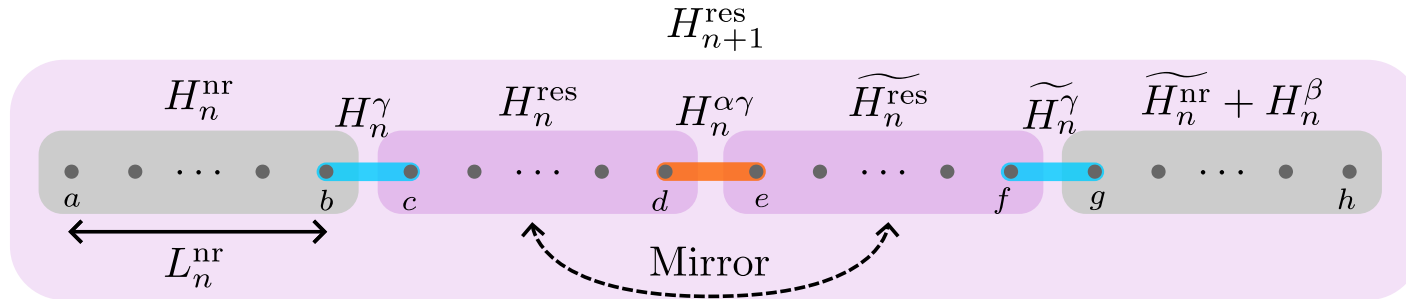
Construction of the solvable model of MBC



$\widetilde{X}$: mirror reflection of X

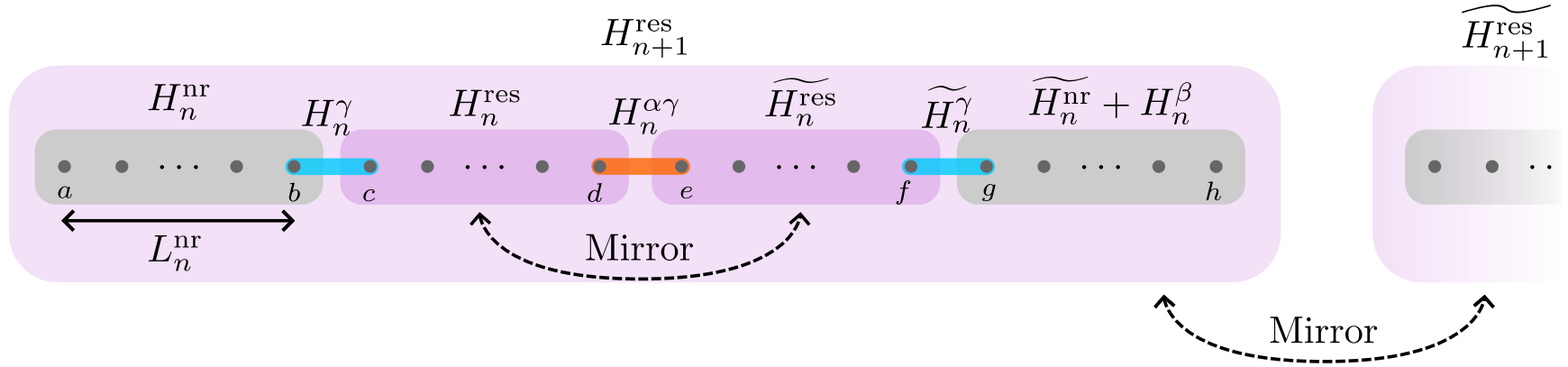
$$|ABCD\rangle_{n+1} \approx \begin{cases} |ABCD\rangle_n, & \text{if } A \neq \widetilde{D} \text{ or } B = C, \\ \frac{1}{\sqrt{2}} (|ABCD\rangle_n + |A\widetilde{C}\widetilde{B}D\rangle_n), & \text{if } A = \widetilde{D}, B > C, \\ \frac{1}{\sqrt{2}} (|ABCD\rangle_n - |A\widetilde{C}\widetilde{B}D\rangle_n), & \text{if } A = \widetilde{D}, B < C. \end{cases}$$

Construction of the solvable model of MBC



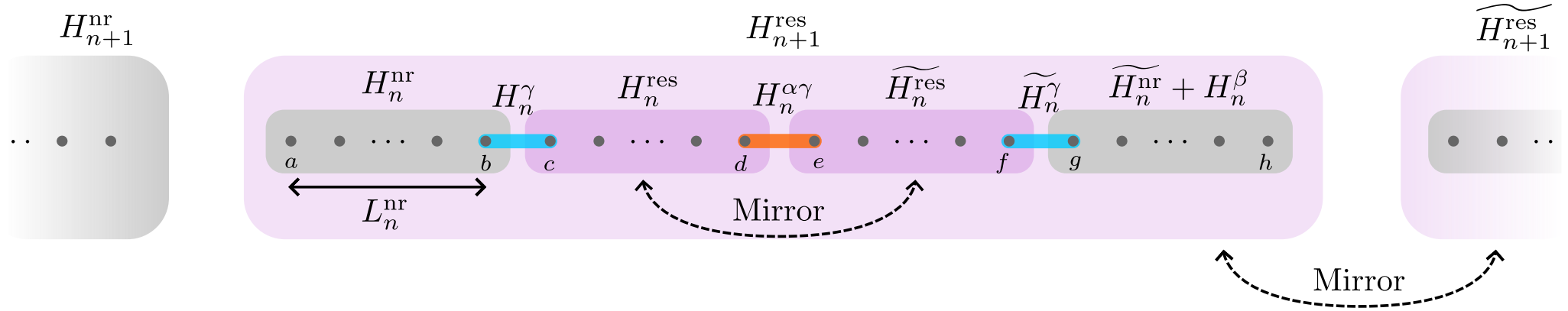
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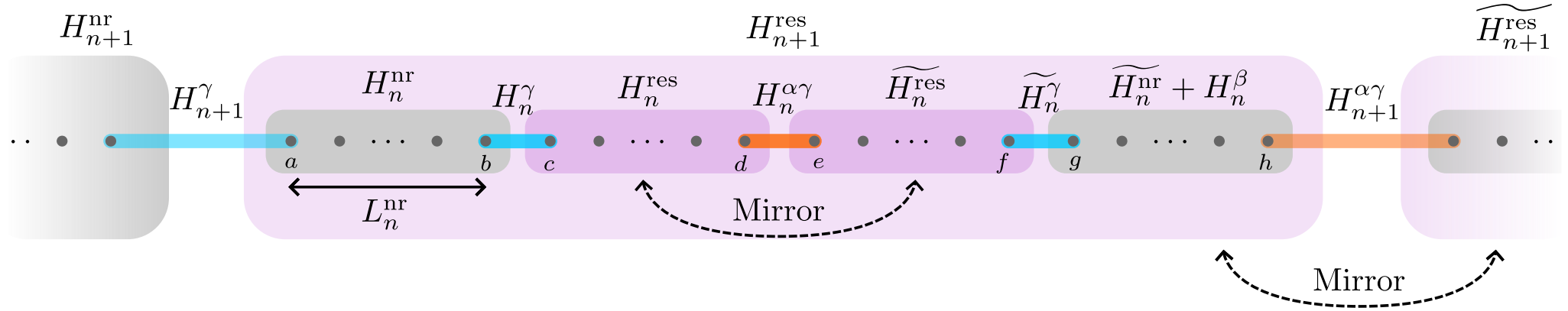
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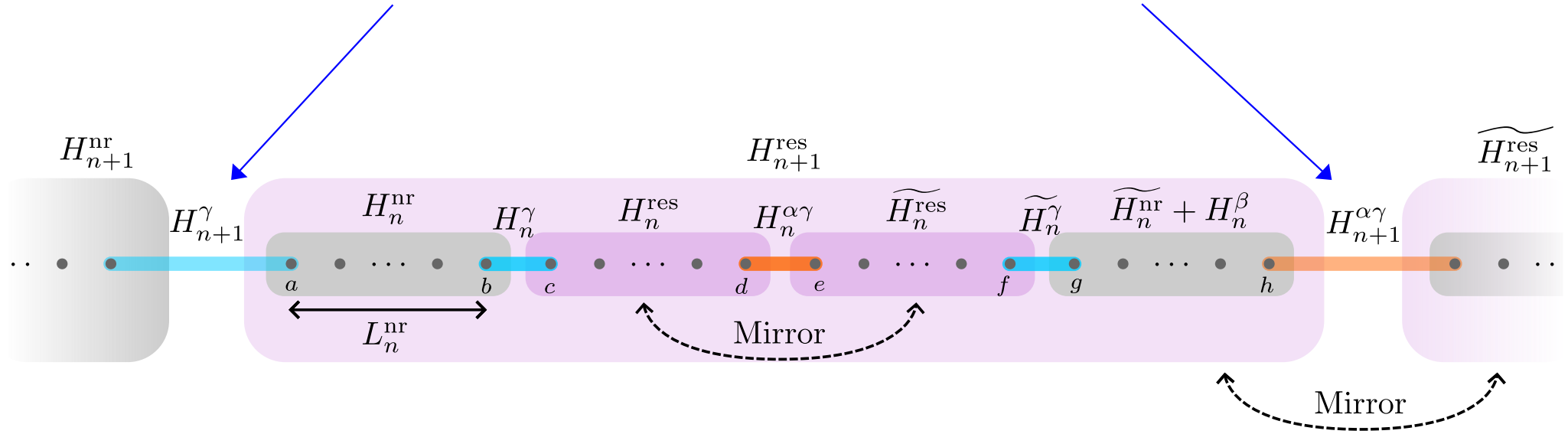
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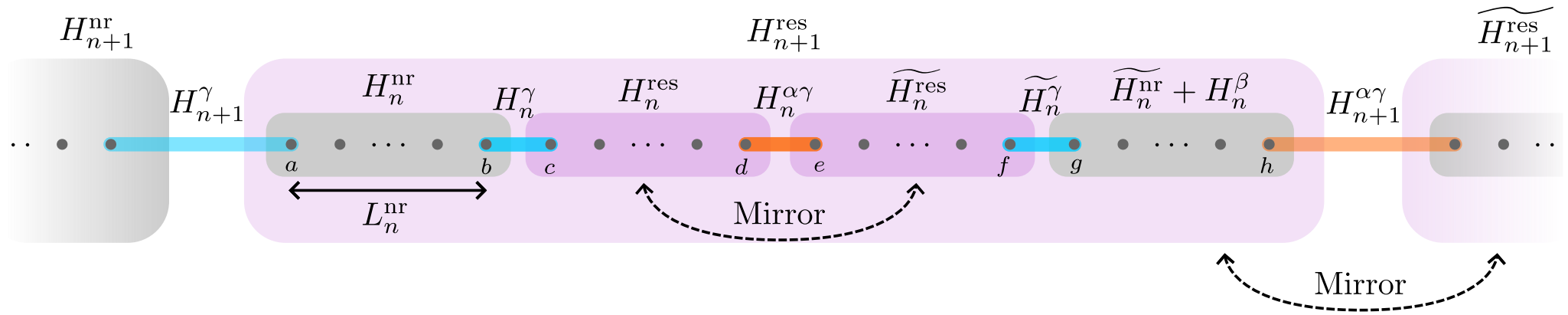
Construction of the solvable model of MBC

Weaker for larger n (avoid accidental resonances)

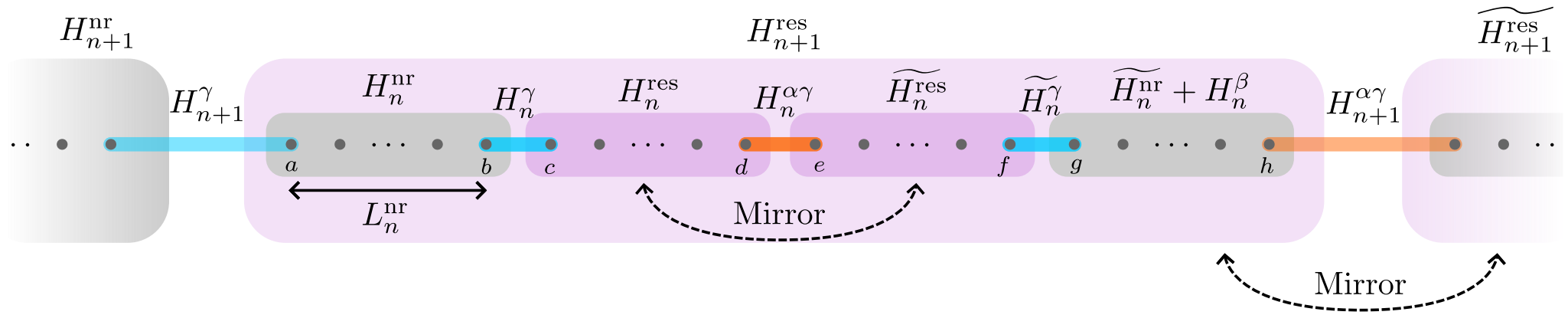


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Probability of mirror resonances



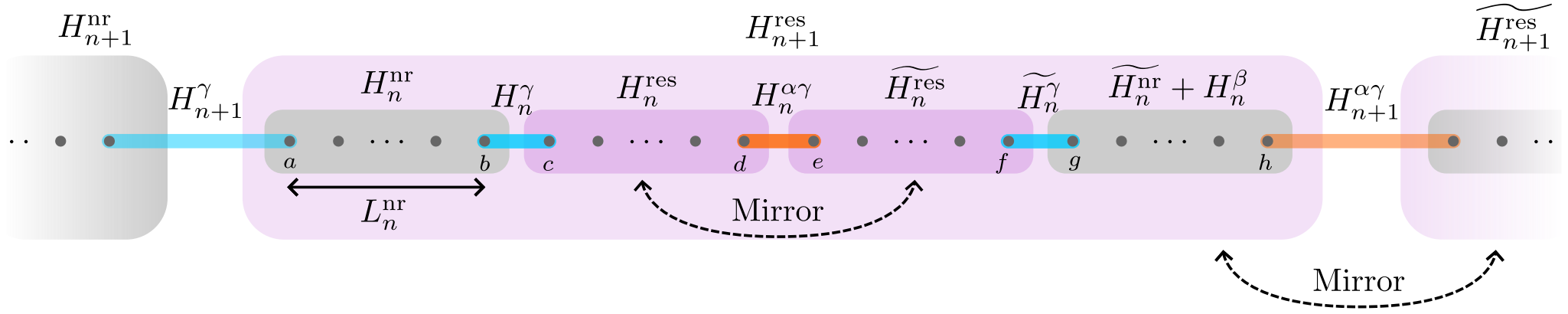
Probability of mirror resonances



Randomly choose a state from the product canonical ensemble of the blocks at level n :

$$|ABCD\rangle_n \sim \rho_n^{\text{nr}} \otimes \rho_n^{\text{res}} \otimes \widetilde{\rho}_n^{\text{res}} \otimes \widetilde{\rho}_n^{\text{nr}}$$

Probability of mirror resonances



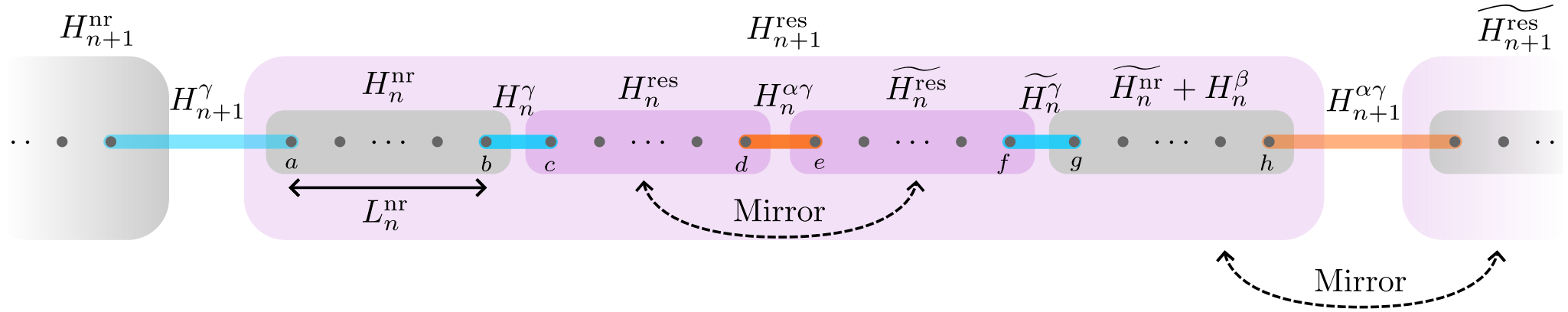
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The corresponding state at the next level becomes:

$$|ABCD\rangle_{n+1} \approx \begin{cases} |ABCD\rangle_n, & \text{with probability } 1 - p_n \\ \frac{1}{\sqrt{2}} (|ABCD\rangle_n \pm |A\widetilde{C}\widetilde{B}D\rangle_n), & \text{with probability } p_n \end{cases}$$

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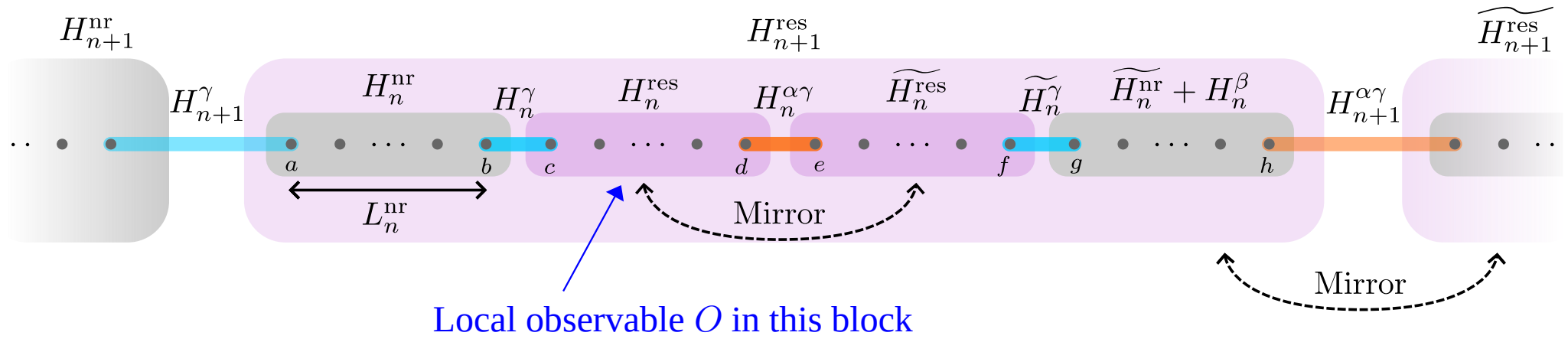
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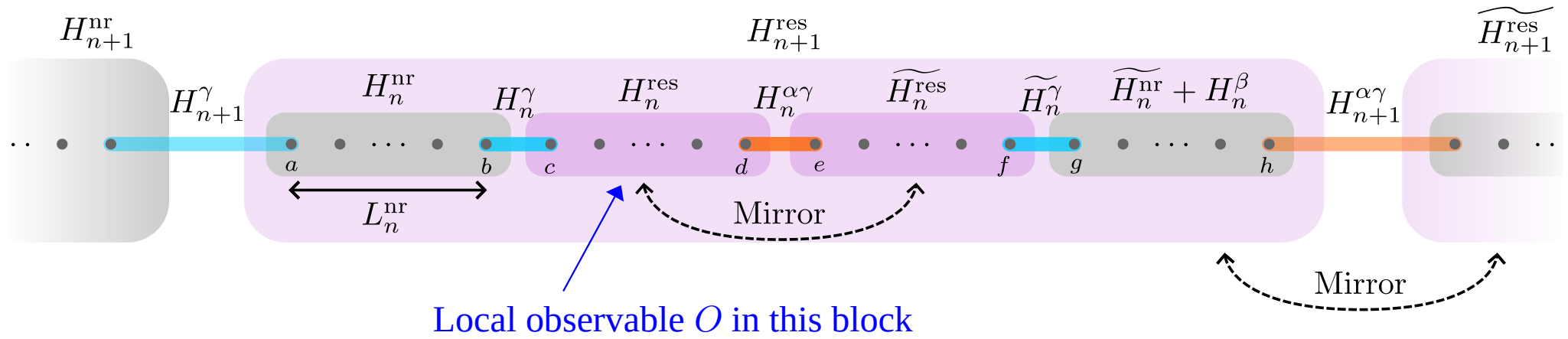
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$$p_n = \text{P}(A = \tilde{D}) = \exp(-S_{2,n}^{\text{nr}}), \quad S_2: \text{thermodynamic Rényi-2 entropy}$$

Variance of local observable expectations

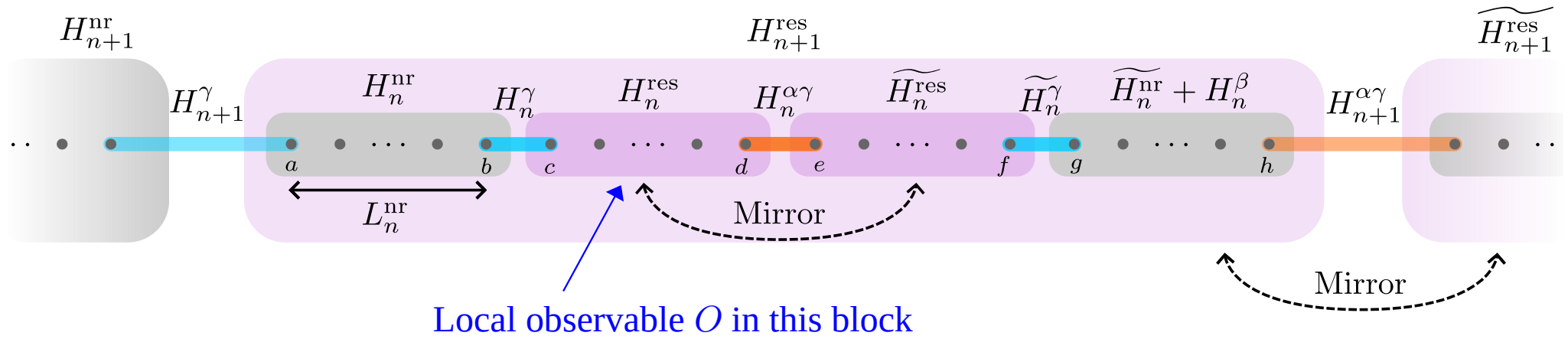


Variance of local observable expectations



$$|X\rangle_{n+1} \approx \begin{cases} |X\rangle_n, & \text{with probability } 1 - p_n \\ \frac{1}{\sqrt{2}} (|X\rangle_n \pm |\widetilde{X}\rangle_n), & \text{with probability } p_n \end{cases}$$

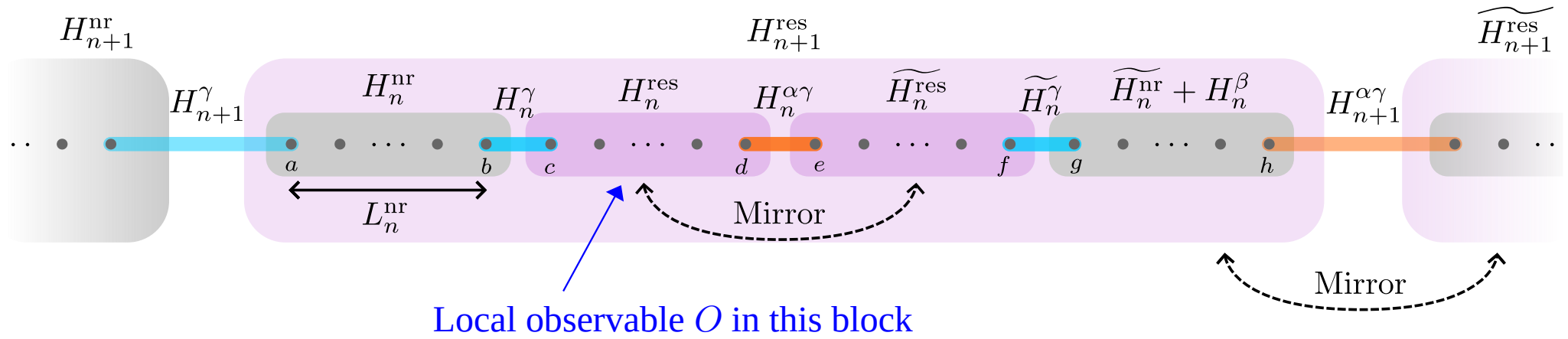
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$$\text{Var}\langle X|O|X\rangle_{n+1} \approx (1 - p_n) \text{Var}\langle X|O|X\rangle_n + p_n \frac{\text{Var}\langle X|O|X\rangle_n + \text{Var}\langle \tilde{X}|O|\tilde{X}\rangle_n}{4}$$

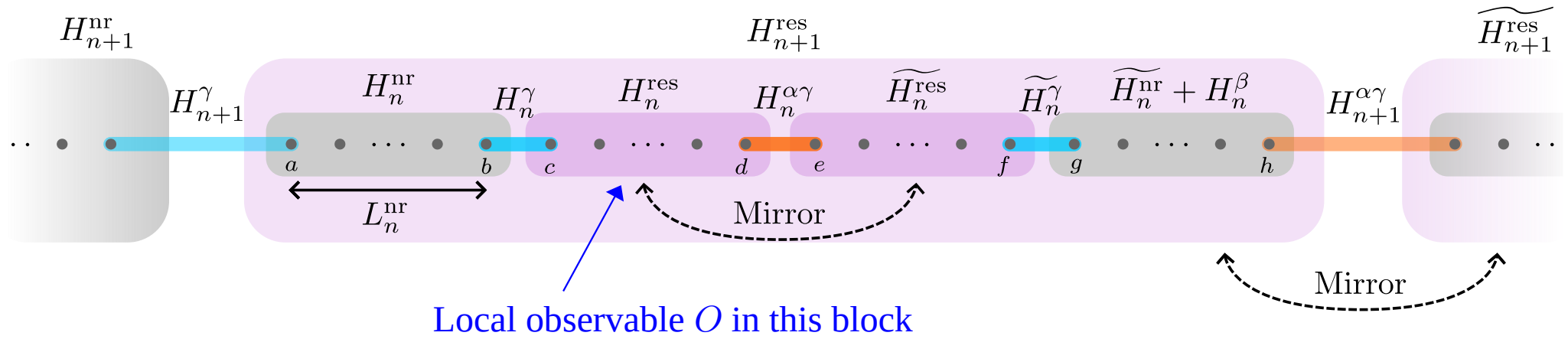
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$$\text{Var}\langle X|O|X\rangle_n \rightarrow 0 \text{ (weak ETH)} \iff \sum_n p_n = \infty$$

Many-body critical phases

YTT, Li, Das Sarma (in preparation)

Single
particle

Critical

(at least for our
solvable model)

E

$L_n^{\text{nr}} \ll \log n$

$L_n^{\text{nr}} \gg \log n$

Interacting
Many-body

E

E

— ETH-like states

— MBL-like states

$$\langle X|O|X \rangle_n \rightarrow \langle O \rangle_{\text{thermal}}$$

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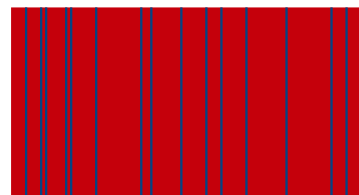
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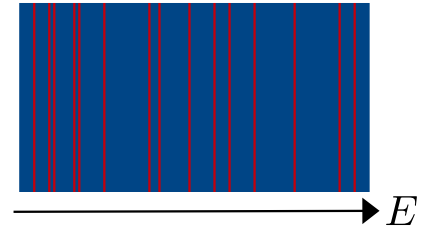
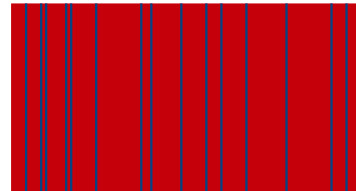
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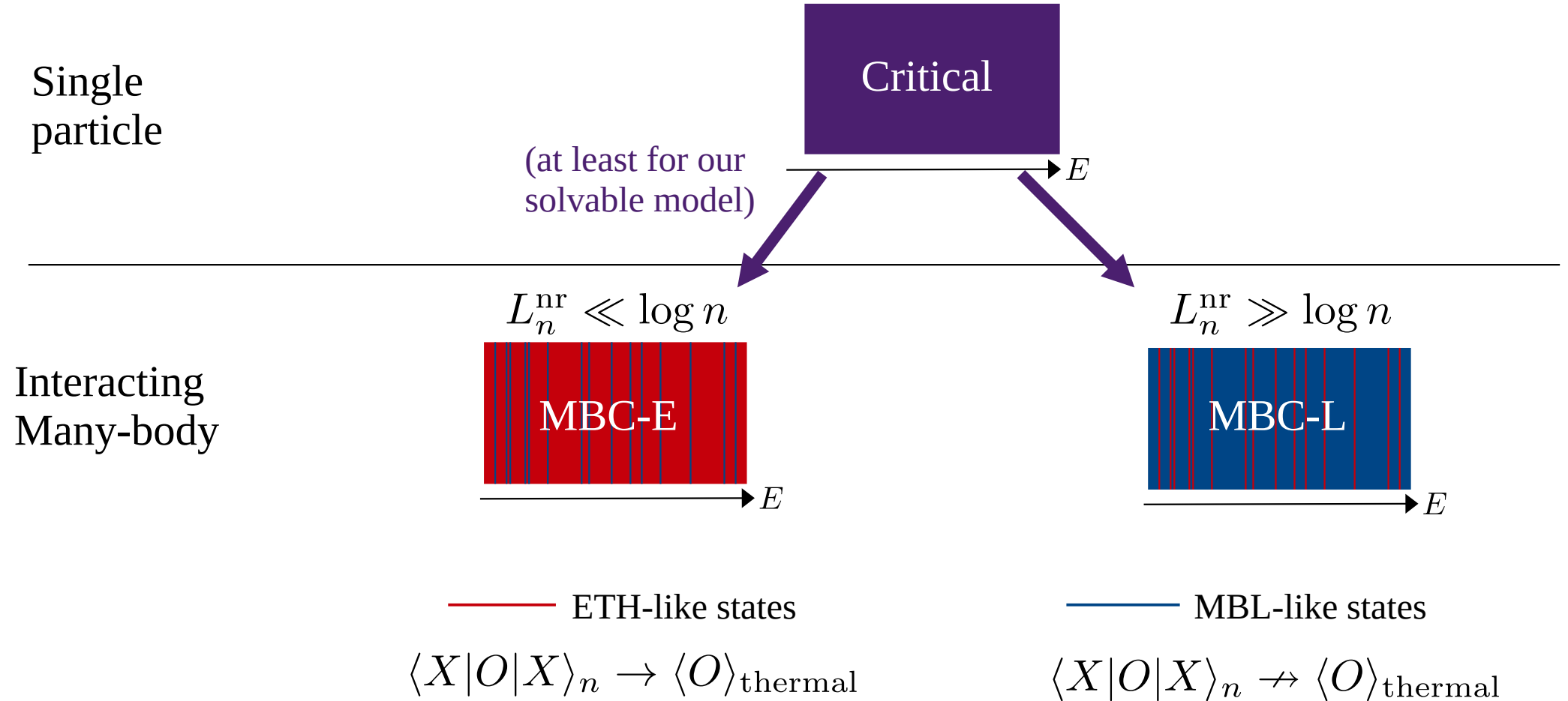
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Many-body critical phases

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E

$$L_n^{\text{nr}} \ll \log n$$

$$L_n^{\text{nr}} \sim \log n$$

$$L_n^{\text{nr}} \gg \log n$$

MBC-E

MBC-L

E

E

E

— ETH-like states

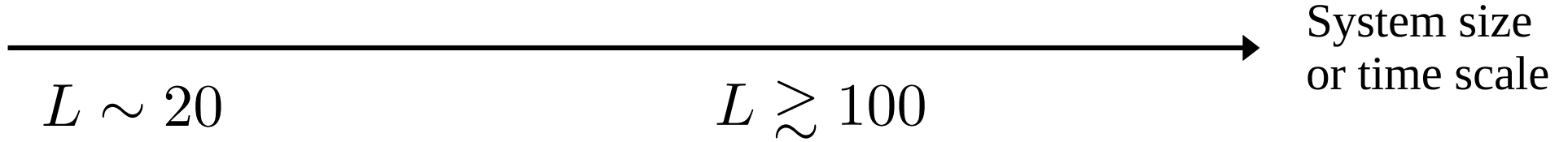
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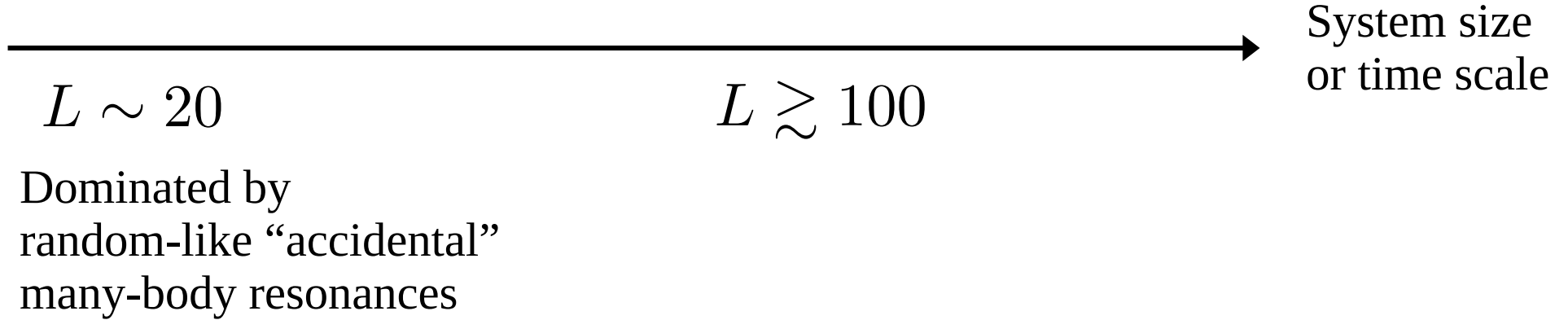
Conclusions & general picture

For a quasiperiodic spin chain, we expect:



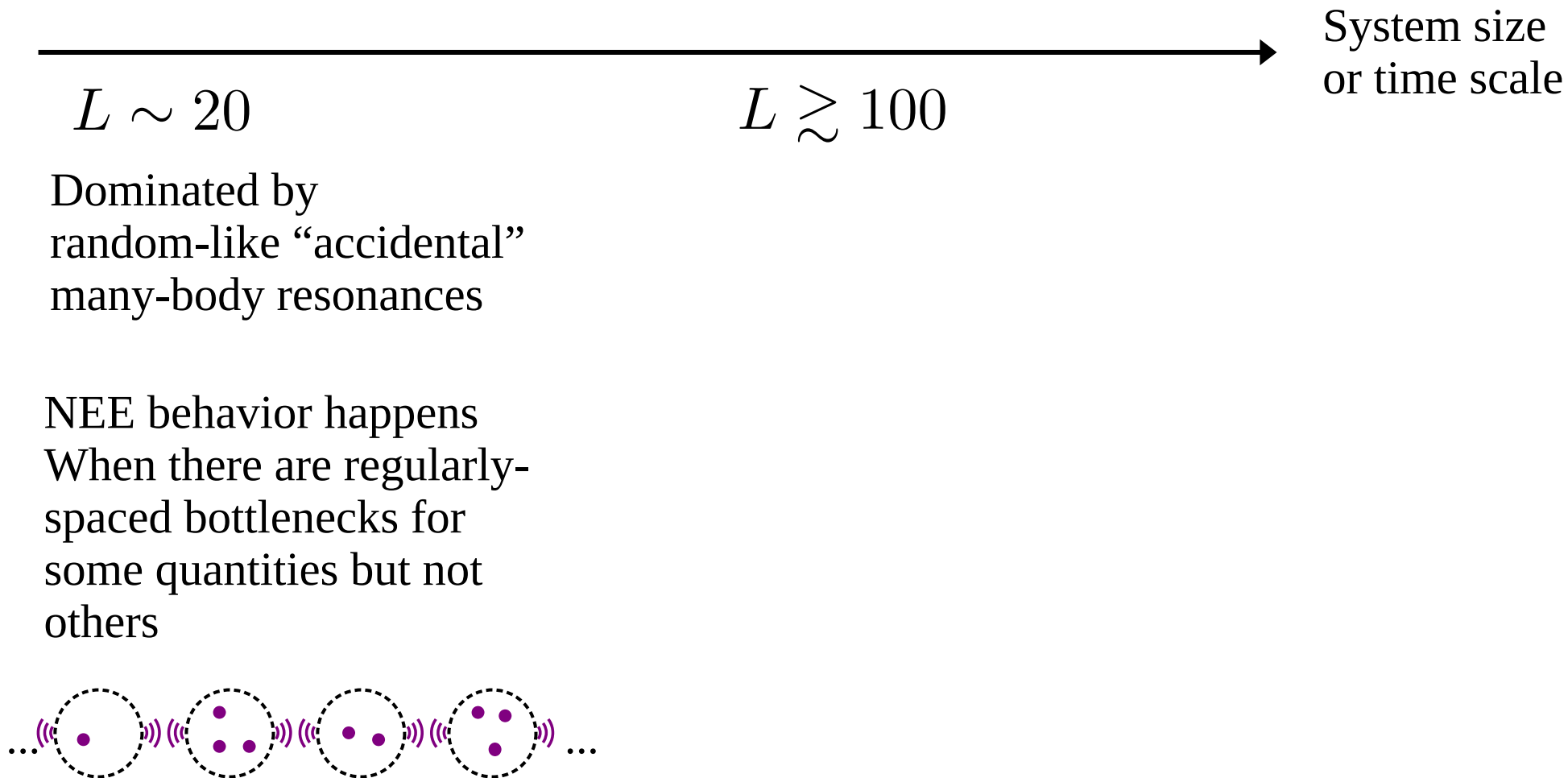
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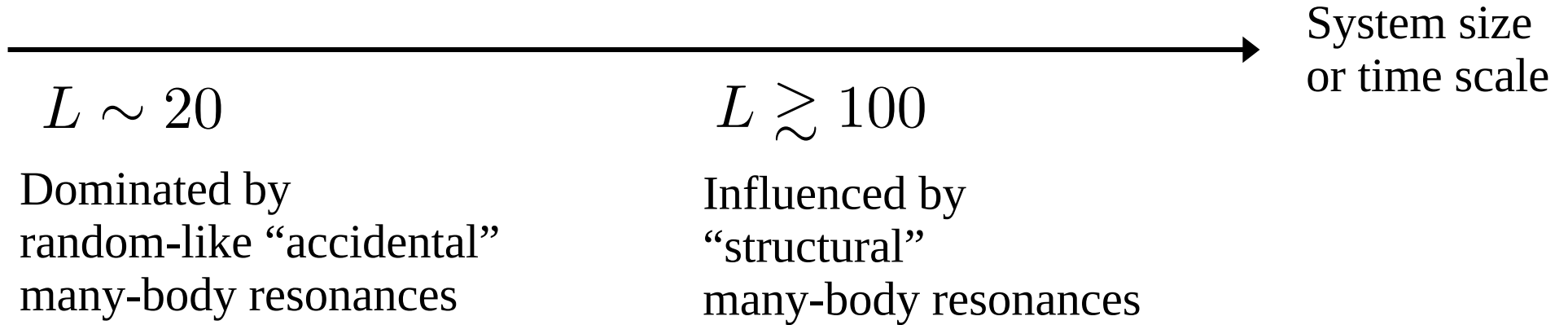
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Conclusions & general picture

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NEE behavior happens
When there are regularly-spaced bottlenecks for some quantities but not others



Conclusions & general picture

For a quasiperiodic spin chain, we expect:

